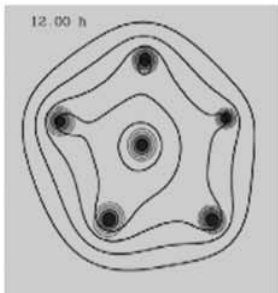
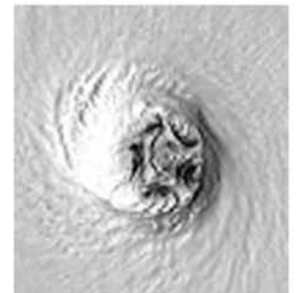
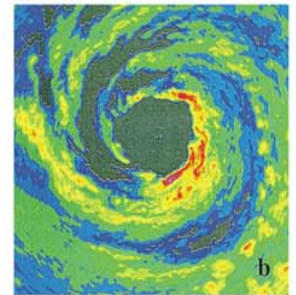
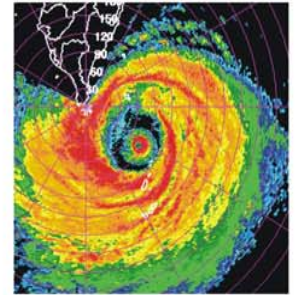


Vortex Interactions and the Barotropic Aspects of Concentric Eyewall Formation



Hung-Chi Kuo
Department of Atmospheric Sciences
National Taiwan University
Taipei, Taiwan

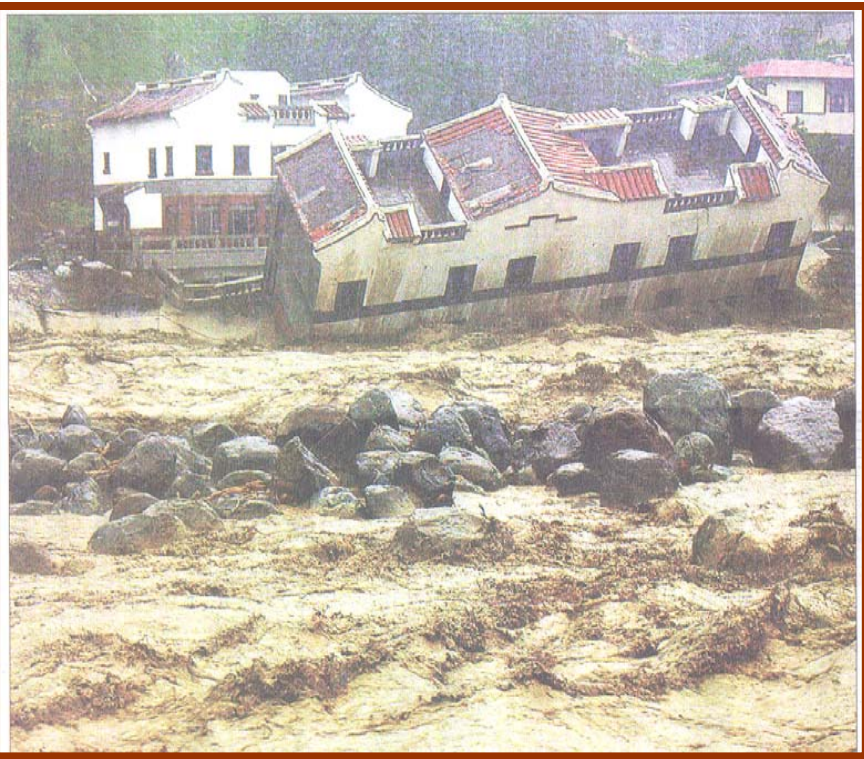
Institute of Mathematical Sciences
NUS
November 14 2006



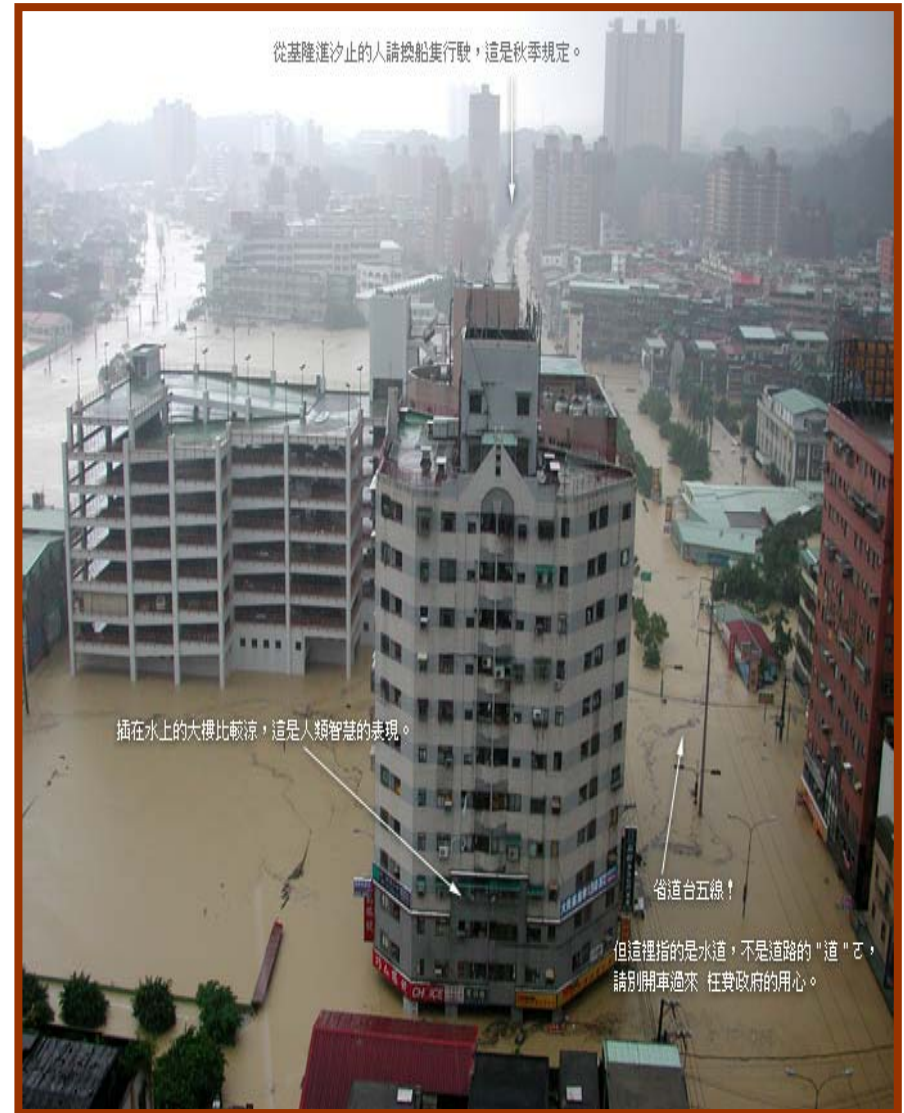
象神颱風侵襲後的台北市



納莉颱風 (2001)



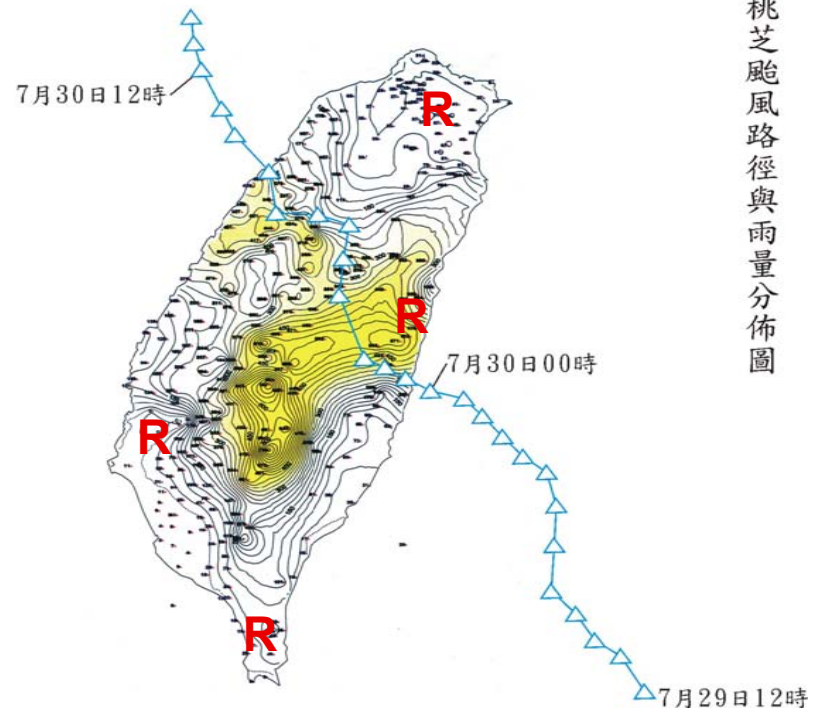
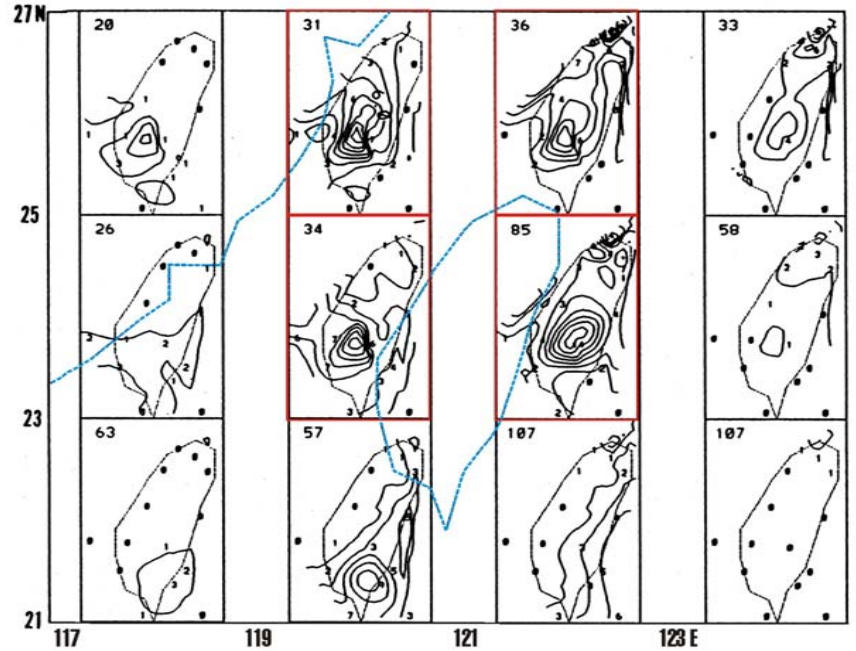
桃芝颱風 (2001)



**CWB is capable of 24 hr
and 100km scale ppn
(phase locked with
topography)**

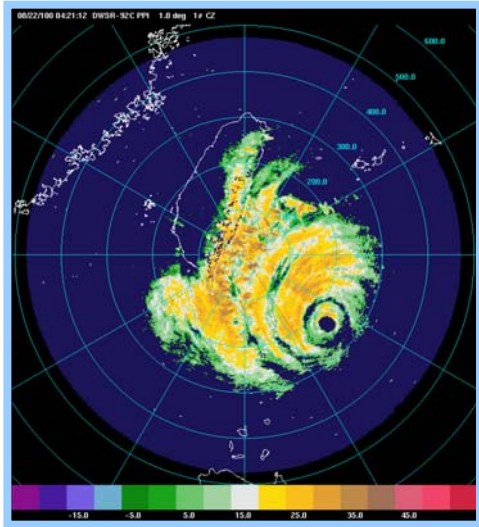
**0 to 12 hr and 10 km ppn
remain biggest challenges**

**355mm in 5 hr in the city of
KaoShung (5 pm to 10 pm
at the beginning of the rush
hour)**

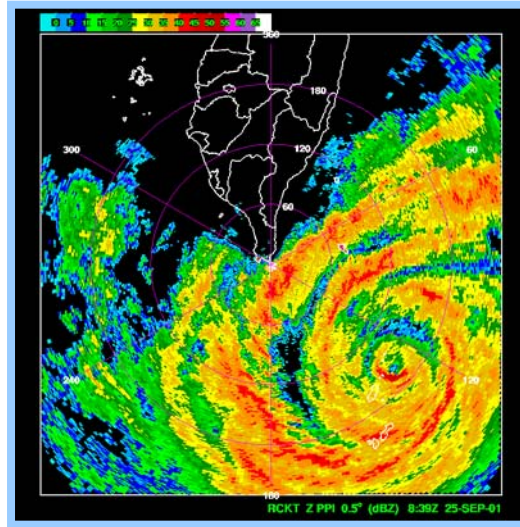


桃芝颱風路徑與雨量分佈圖

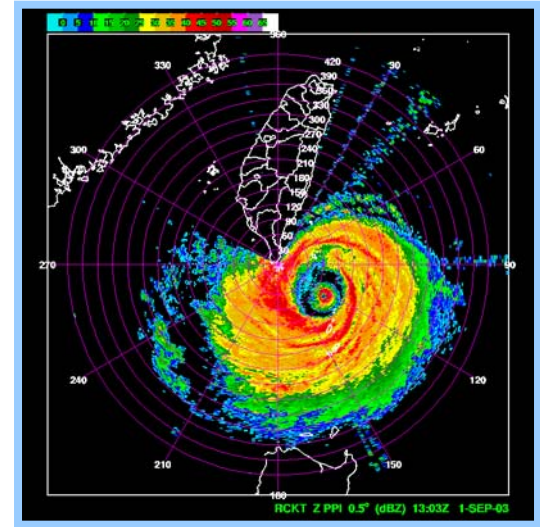
Concentric eyewalls near Taiwan



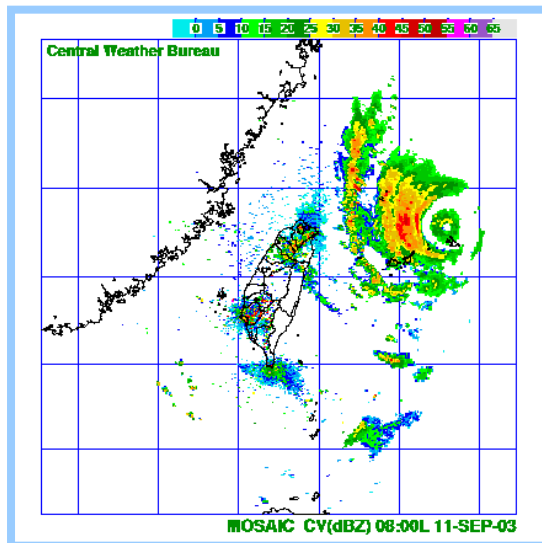
Bilis(2000)



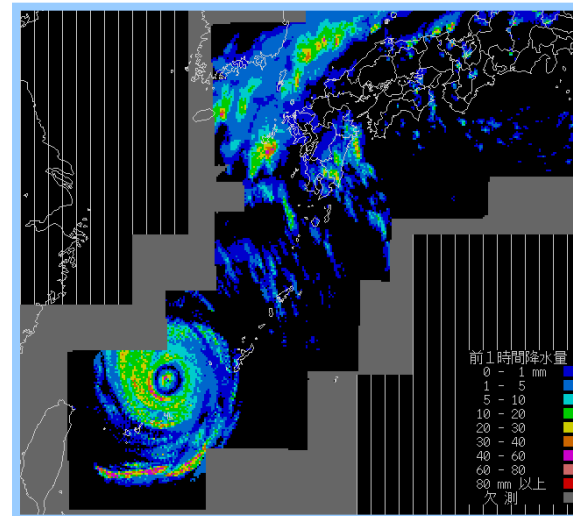
Lekima(2001)



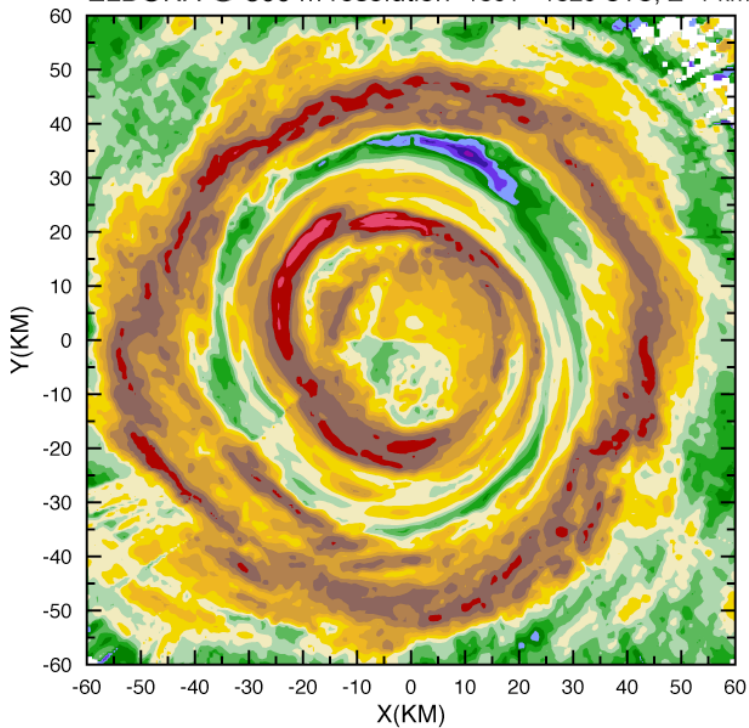
Dujan(2003)



Maemi(2003)

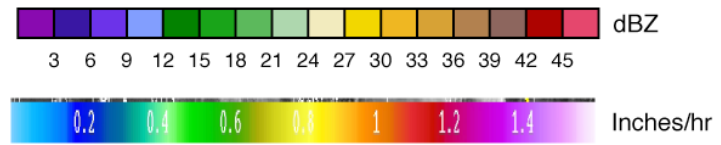
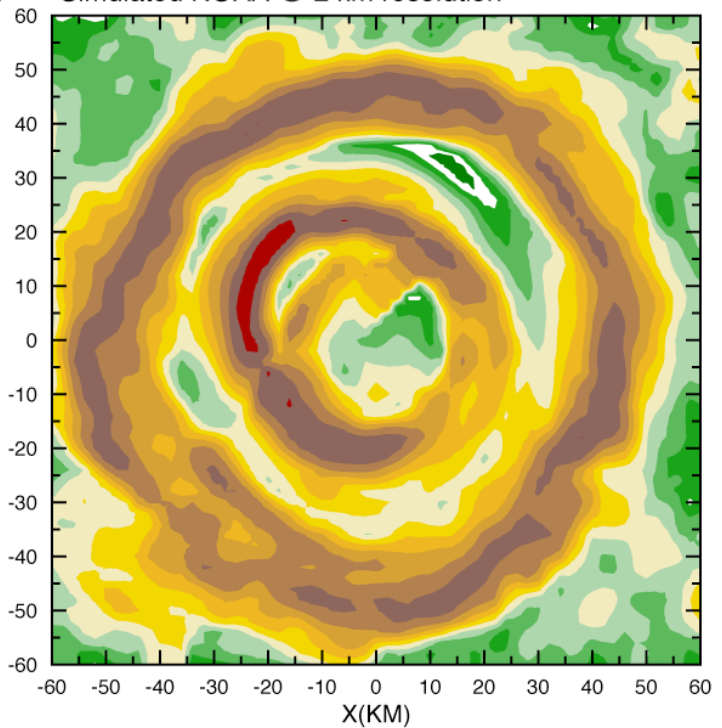


ELDORA @ 500 m resolution 1801 - 1820 UTC, Z=1 km

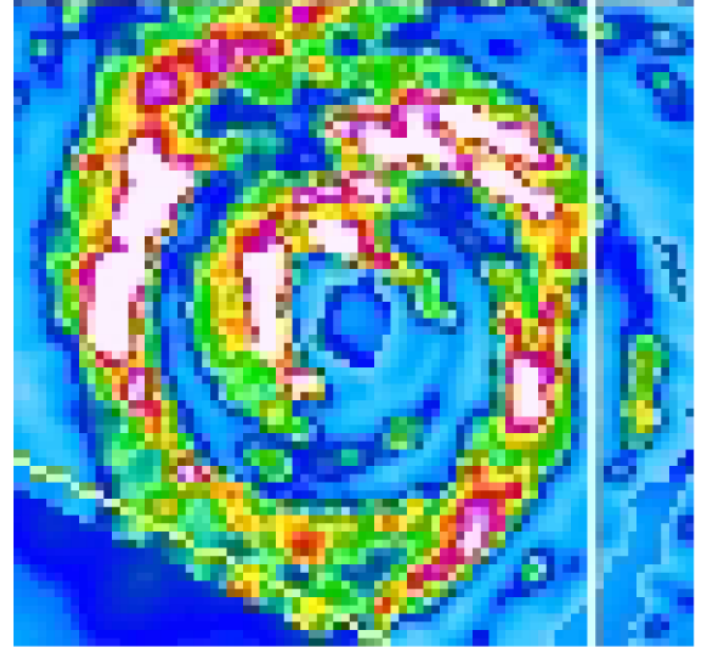


ELDORA

Simulated NOAA @ 2 km resolution



TRMM Precipitation Radar @ 4 km resolution (1443 UTC)

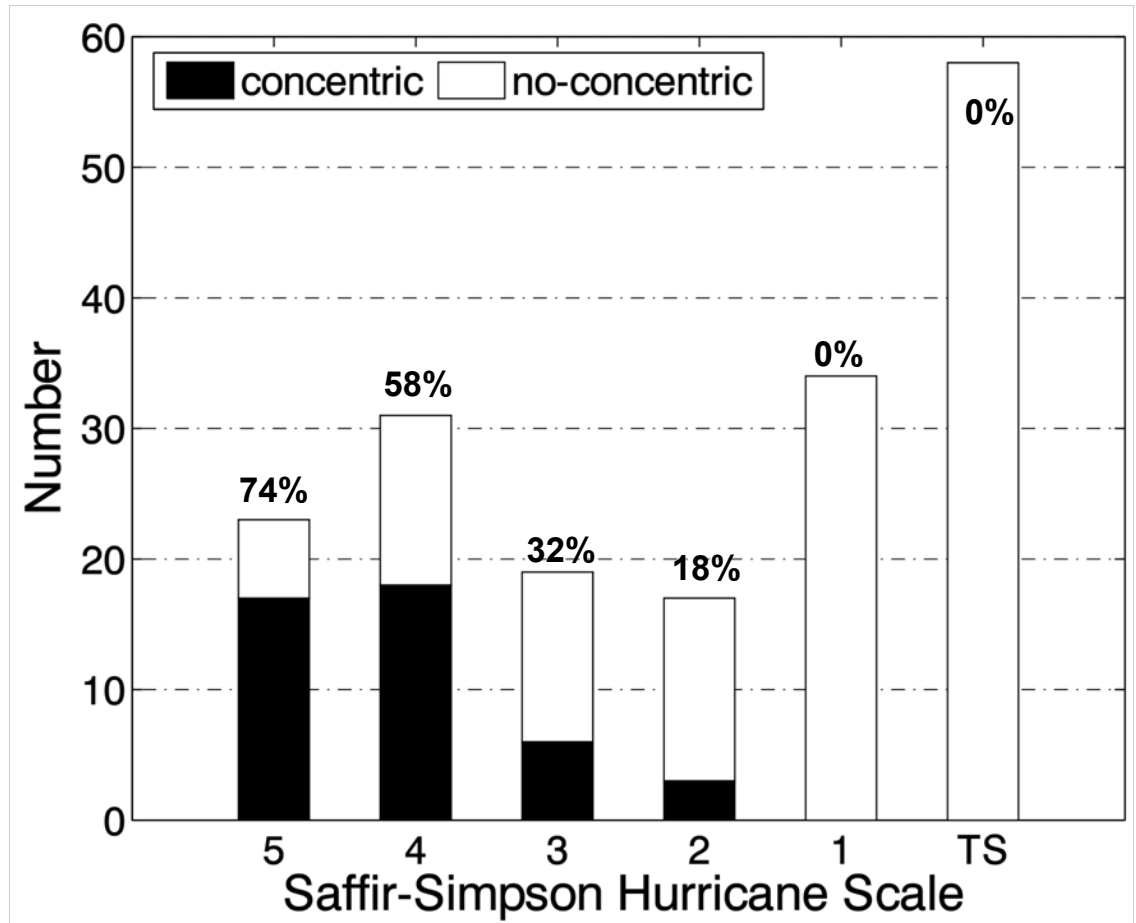


RAINEX

TRMM

WNPAC

- 1997~2004
- 197 tropical cyclones
- 182 tropical cyclones examined
- 44 concentric tropical cyclones (25%)
- higher intensity, higher percentage of concentric eyewall



category 5 (135+ kts)

category 4 (114-135 kts)

category 3 (96-113 kts)

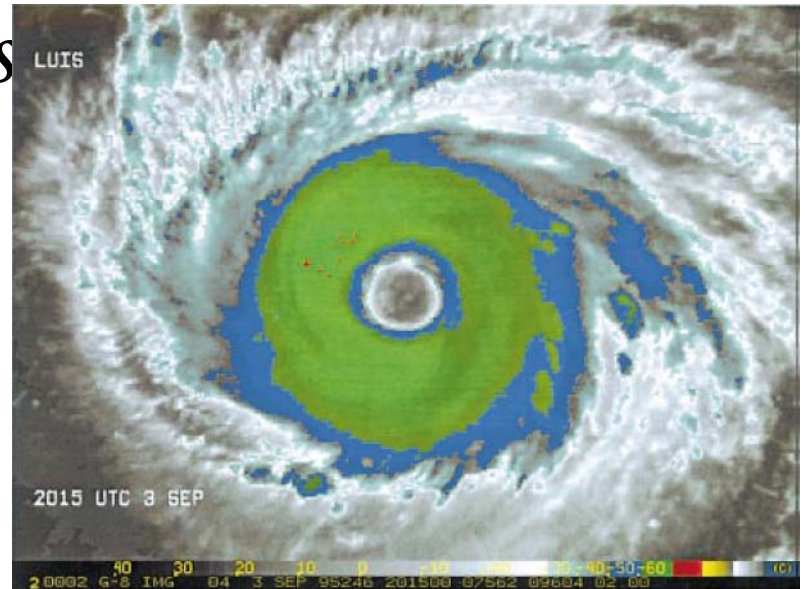
category 2 (83-95 kts)

category 1 (64-82 kts)

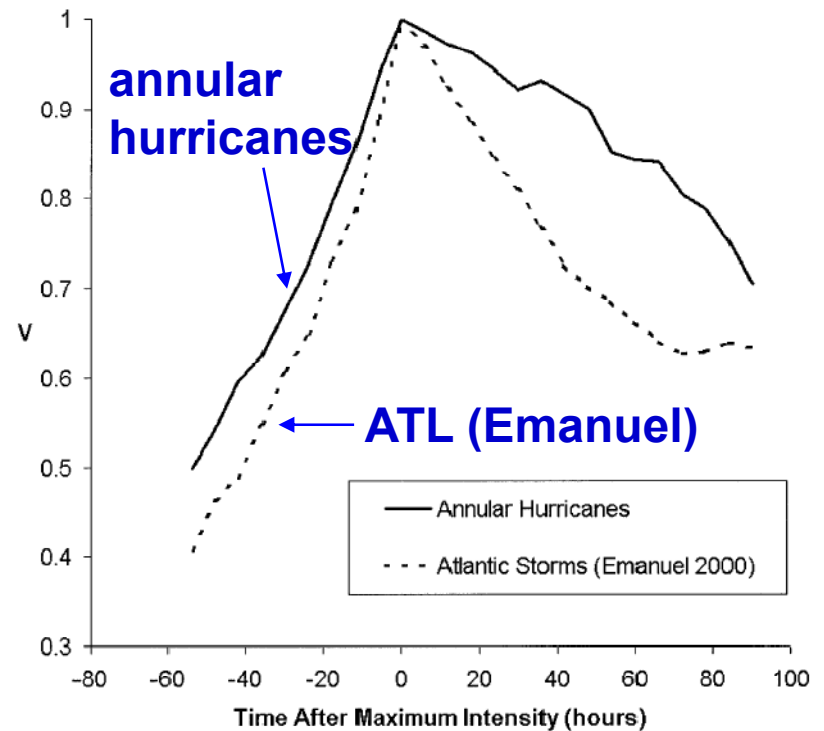
TS(63- kts)

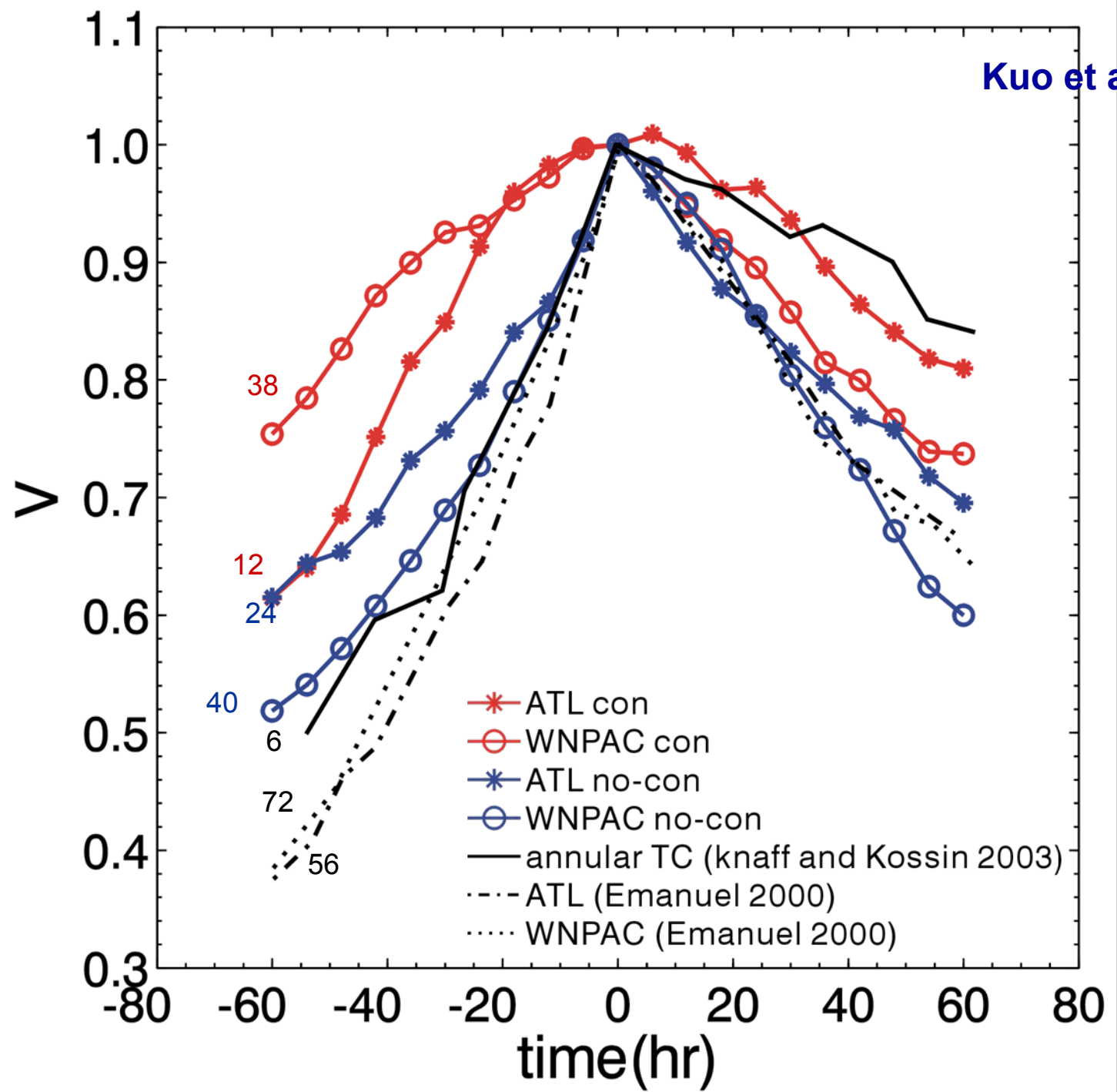
Knaff and Koss

- color-enhanced IR image of Hurricane Luis (1995) at 2015 UTC 3 Sep



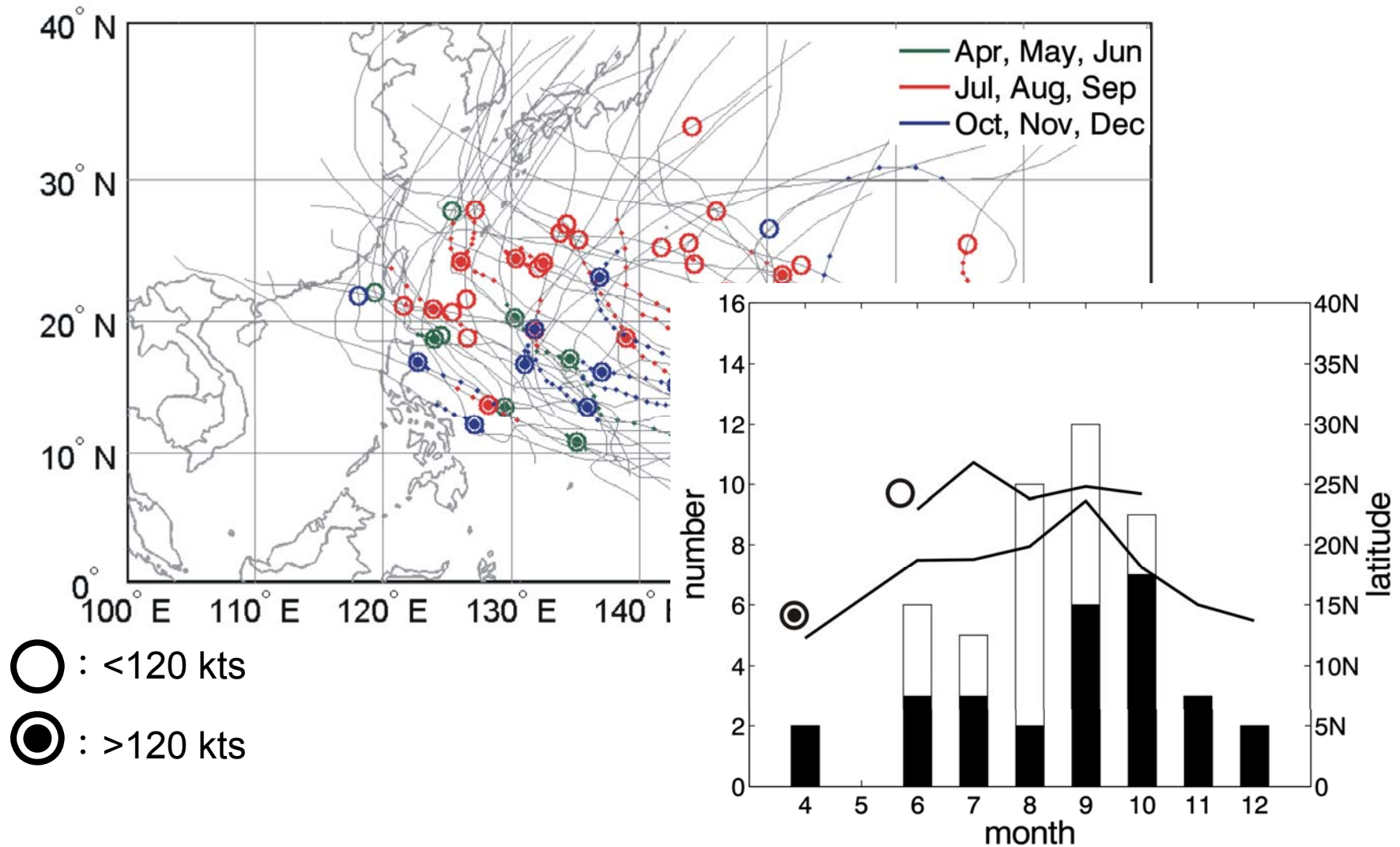
dimensionless	24-h weakening
ATL(56)	0.14
Annular hurricanes(6)	0.05





WNPAC Concentric eyewalls formation locations, intensity, and tracks

Kuo et al. (2006)



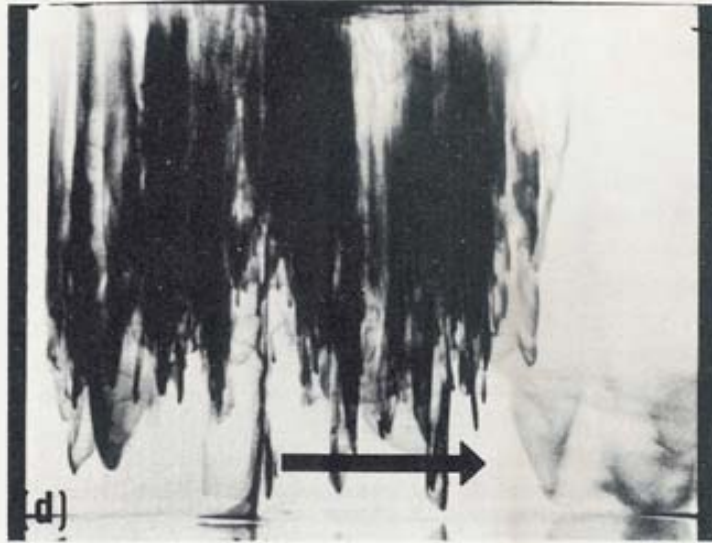
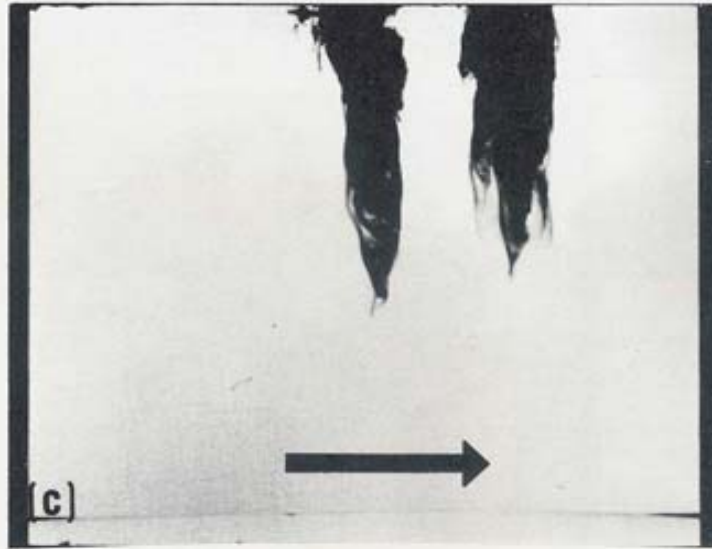
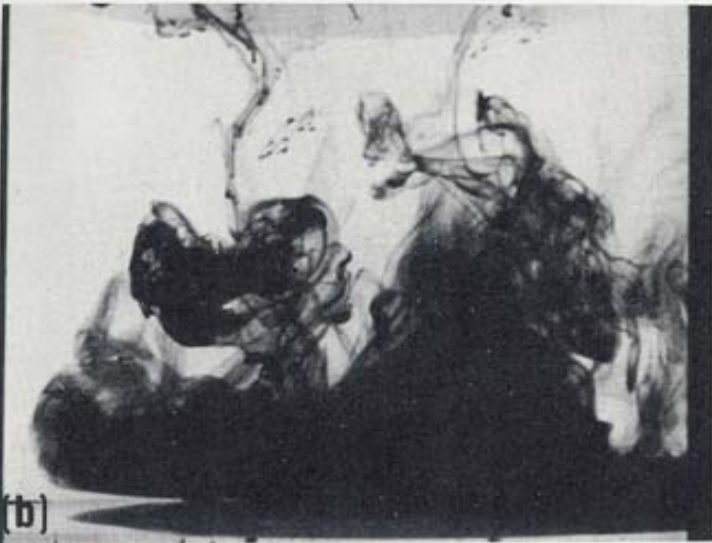
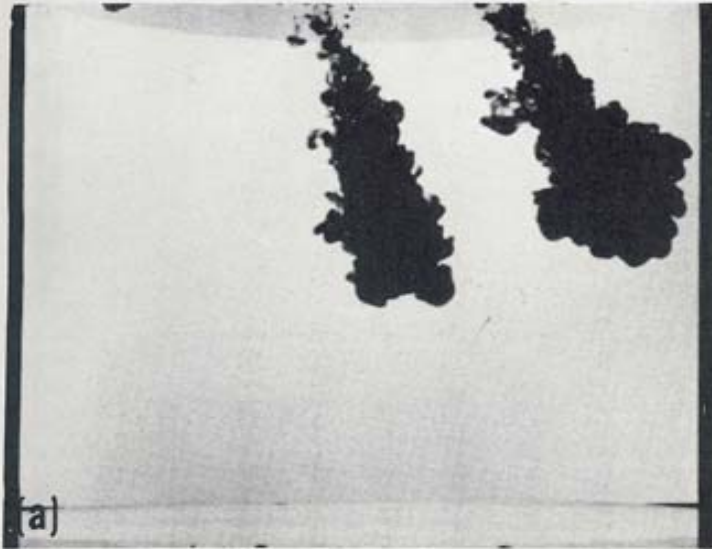
Coriolis Force



Non-inertial Frame

3-D

2-D (under rotation)

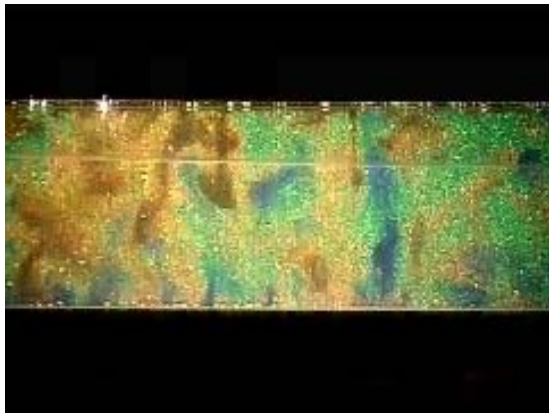


Taylor
column

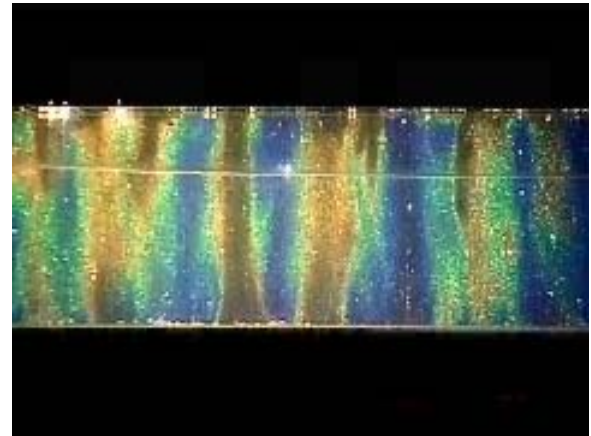
10. When dye is injected into non-rotating water (a), it spreads in a typical turbulent field (b). With identical ink injection into rotating water, the initial motions are

rapidly converted to nearly two-dimensional motions (c), and after a time, the ink is found distributed in vertical Taylor walls (d).

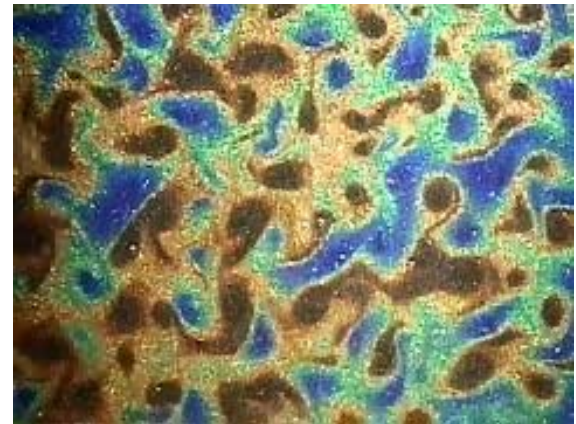
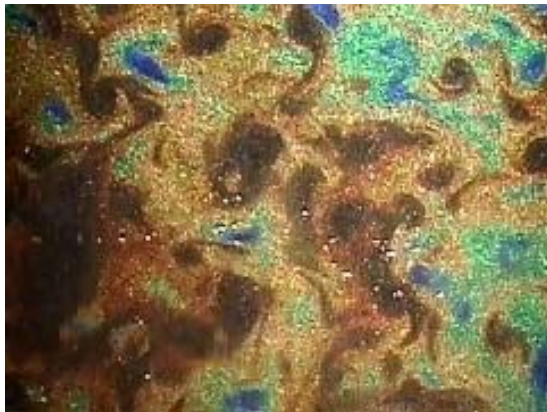
3D



2D (strong rotation)



Taylor columns



Vortices with sharp edge

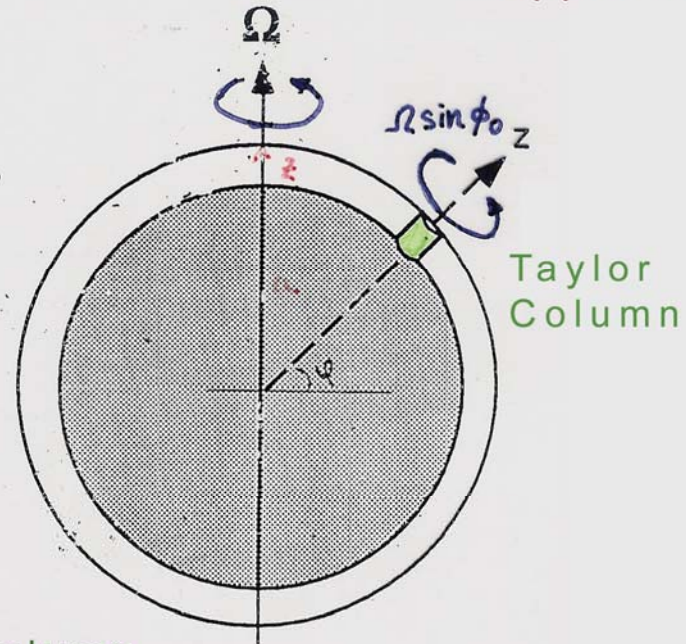
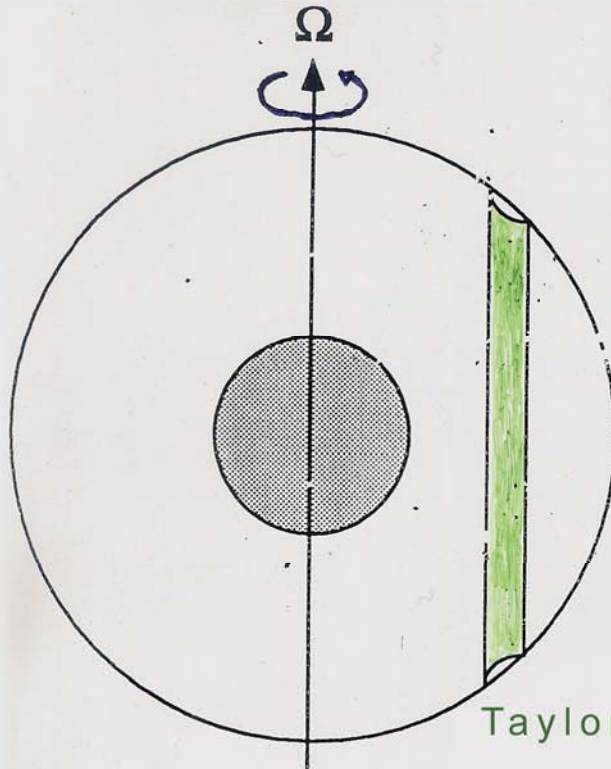
$$\vec{f} = 2\Omega \cos\phi \hat{j} + 2\Omega \sin\phi \hat{k}$$

$$\vec{f} \equiv 2\Omega \sin\phi \hat{k}$$

deep atmosphere
(Jupiter??)

shallow atmosphere
(Earth) $z \ll a$

traditional
approximation

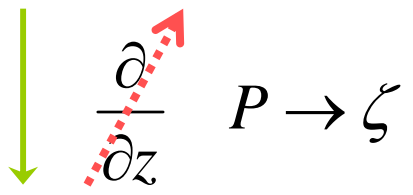


Taylor Column

Taylor
Column

$$\frac{DP}{Dt} = \frac{\partial P}{\partial t} + u \frac{\partial P}{\partial x} + v \frac{\partial P}{\partial y} + w \frac{\partial P}{\partial z} \quad P = \frac{\vec{\zeta} \cdot \nabla \theta}{\rho}$$

Potential Vorticity



$$\frac{D\zeta}{Dt} = \frac{\partial \zeta}{\partial t} + u \frac{\partial \zeta}{\partial x} + v \frac{\partial \zeta}{\partial y}$$

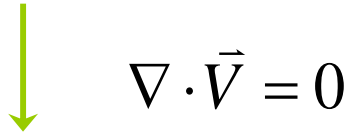
Barotropic Vorticity

$$\frac{D\zeta_g}{Dt} = \frac{\partial \zeta_g}{\partial t} + u_g \frac{\partial \zeta_g}{\partial x} + v_g \frac{\partial \zeta_g}{\partial y}$$

$$\zeta_g = \nabla^2 \psi \quad u_g = -\frac{\partial \psi}{\partial y} \quad v_g = \frac{\partial \psi}{\partial x}$$

Quasi-Geostrophic Vorticity

(Strong Rotation)



$$\frac{\partial \zeta}{\partial t} + \frac{\partial(\psi, \zeta)}{\partial(x, y)} = \nu \nabla^2 \zeta$$

Nondivergent Barotropic Equation

Non-divergent barotropic model (Nearly Inviscid Fluid)

$$\frac{\partial}{\partial t} \zeta + \mathbf{J}(\psi, \zeta) = \nu \nabla^2 \zeta \quad \boxed{\nabla^2 \psi = \zeta}$$

The energy and enstrophy relations

$$\frac{d\mathcal{E}}{dt} = -2\nu \mathcal{Z}$$

$$\mathcal{E} = \iint \frac{1}{2} (u^2 + v^2) dx dy \quad \text{kinetic energy}$$

$$\mathcal{Z} = \iint \frac{1}{2} \zeta^2 dx dy \quad \text{enstrophy}$$

$$\frac{d\mathcal{P}}{dt} = -2\nu \mathcal{P}$$

$$\mathcal{P} = \iint \frac{1}{2} \nabla \zeta \cdot \nabla \zeta dx dy \quad \text{palinstrophy}$$

**Small viscosity led to
large palinstrophy and the
large enstrophy cascade**



Stirring



$E \sim p'^2 / L^2$ (KE) geostrophy

$Z \sim p'^2 / L^4$ (Enstrophy)

KE nearly conserved $L \sim p'$

Enstrophy cascade $L \uparrow$ (L increase Z decrease)

Selective Decay of 2D turbulence

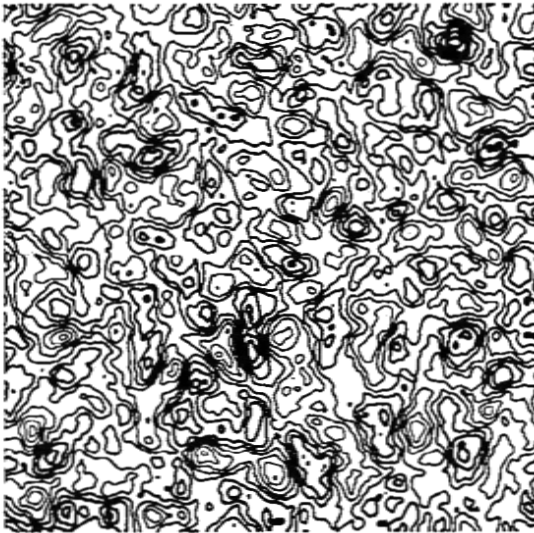
The vortices become, on the average,
larger, stronger, and fewer.

Merger and Axisymmetrization Dynamics

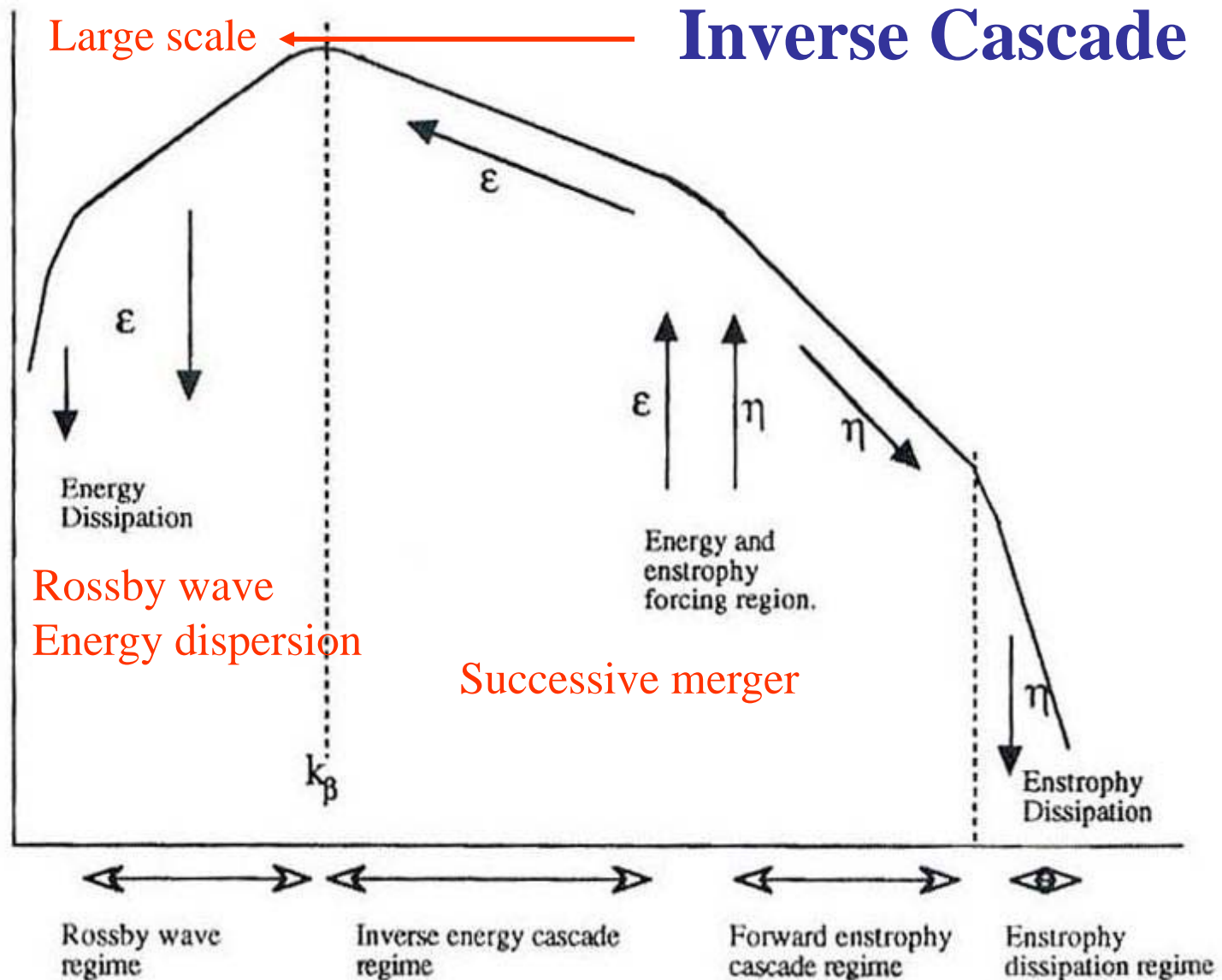
Emergence of coherent structure in 2-D turbulence

The tendency toward successive merger

Fewer and stronger vortices !!!



Waves, turbulence, and coherent vortex



Huang and Robinson
1998

Turbulence



Rhines curve



Waves

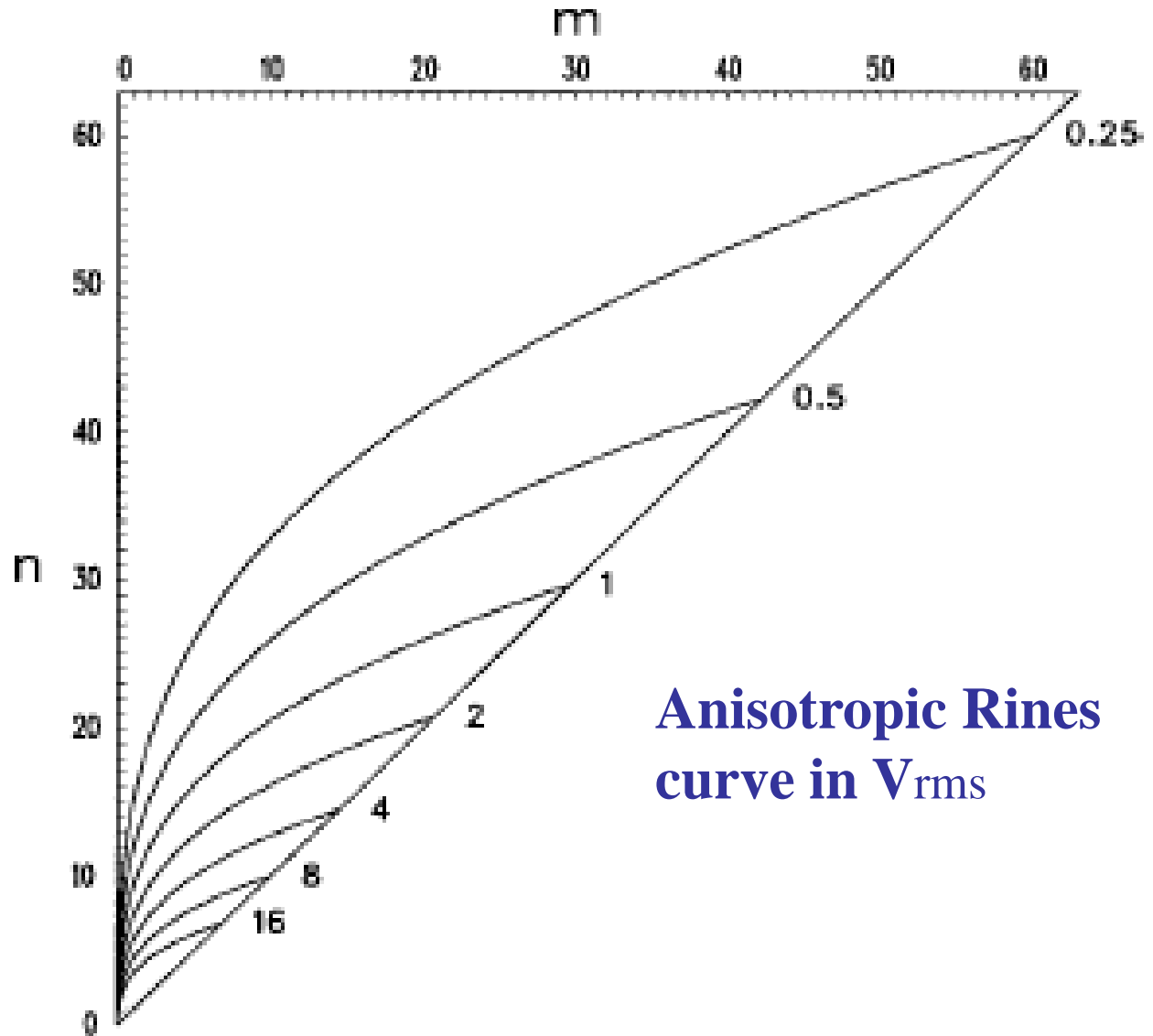
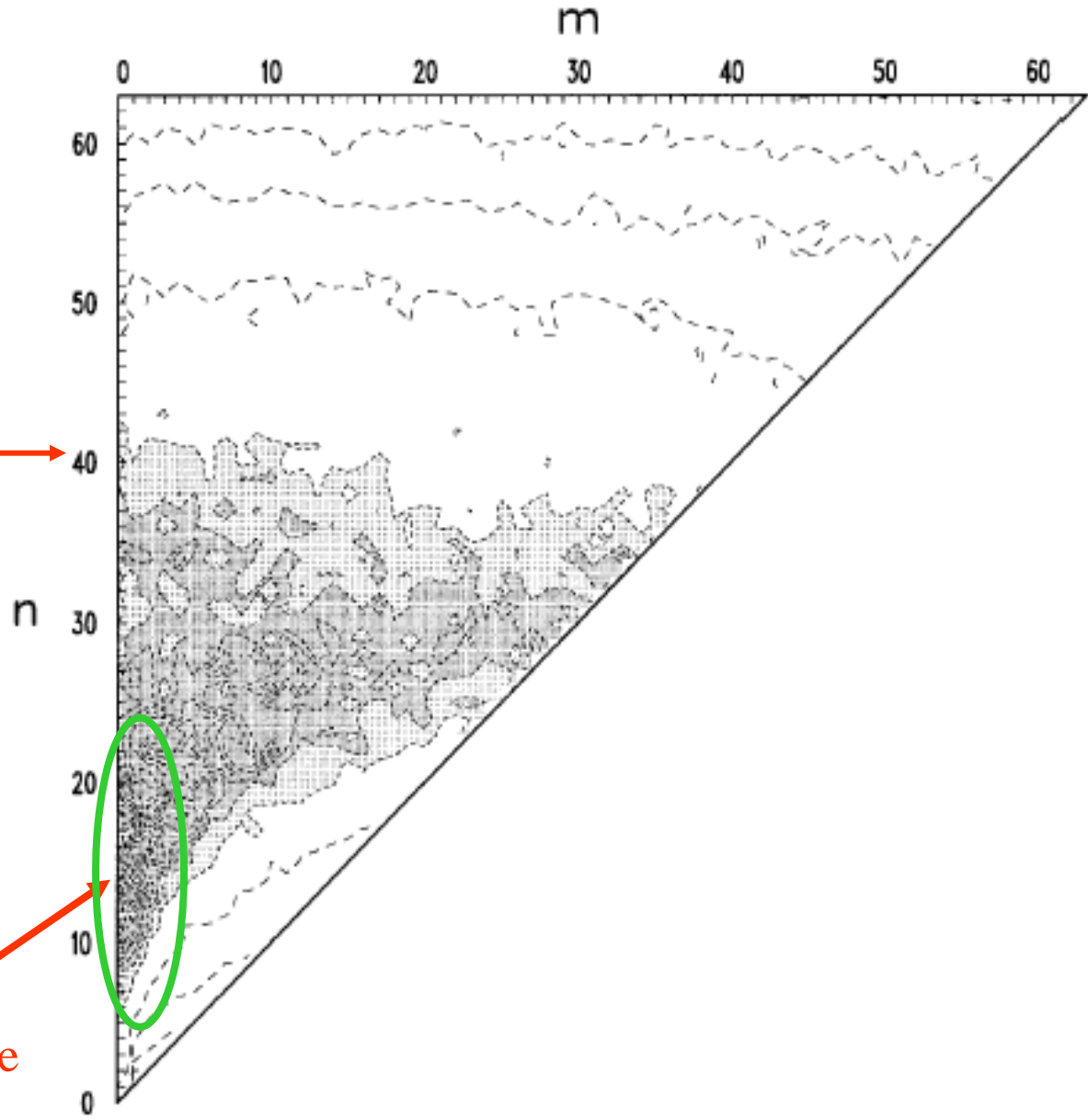


FIG. 1. Anisotropic Rhines curve on the wavenumber plane based on Eq. (3). Dimensional values of V_{rms} (in $m s^{-1}$) are labeled at right.

Huang and Robinson
1998

Forcing added



Inverse energy cascade
to zonal Harmonics

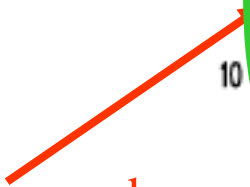
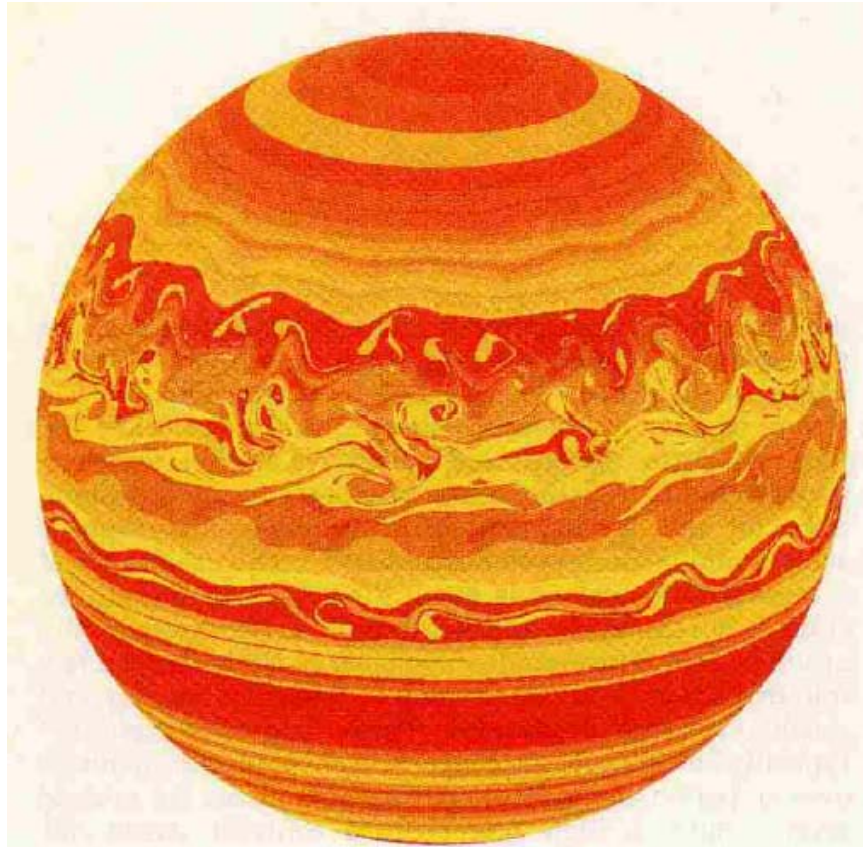


FIG. 2. Energy spectrum $E(m, n)$ of an ensemble mean at day 80 of 10 decaying turbulence experiments. The magnitude of the spectrum is normalized by the maximum value on the map. Contour levels are 0.0001, 0.001, 0.01, 0.1–0.9 with increment 0.1. Area with $E(m, n) > 0.1$ is lightly shaded, $E(m, n) > 0.2$ heavily shaded.



Simulated Jovian atmosphere calculated by contour surgery for a single-layer planetary atmosphere starting with the observed zonal winds of Jupiter.¹⁰ The overall strong potential-vorticity gradient from pole to pole (from positive to negative q) is characteristic of rapid, almost rigid rotation of the atmosphere. Superposed on this global gradient are numerous latitudinal striations indicating zonal gradient reversals, some of which give rise here to nonlinear instabilities. **Figure 4**

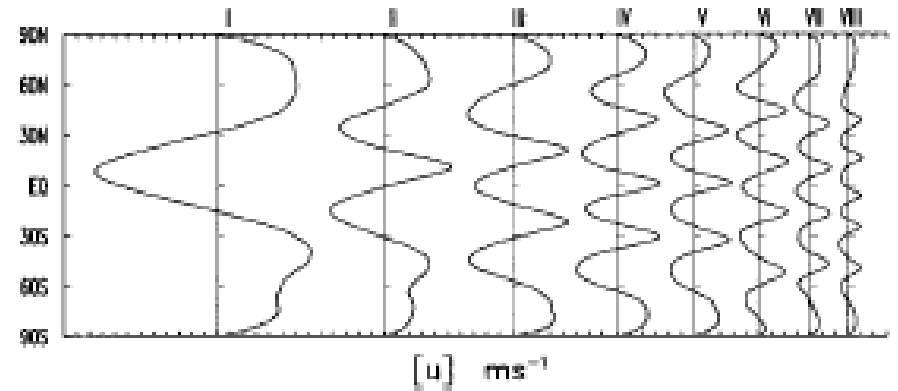
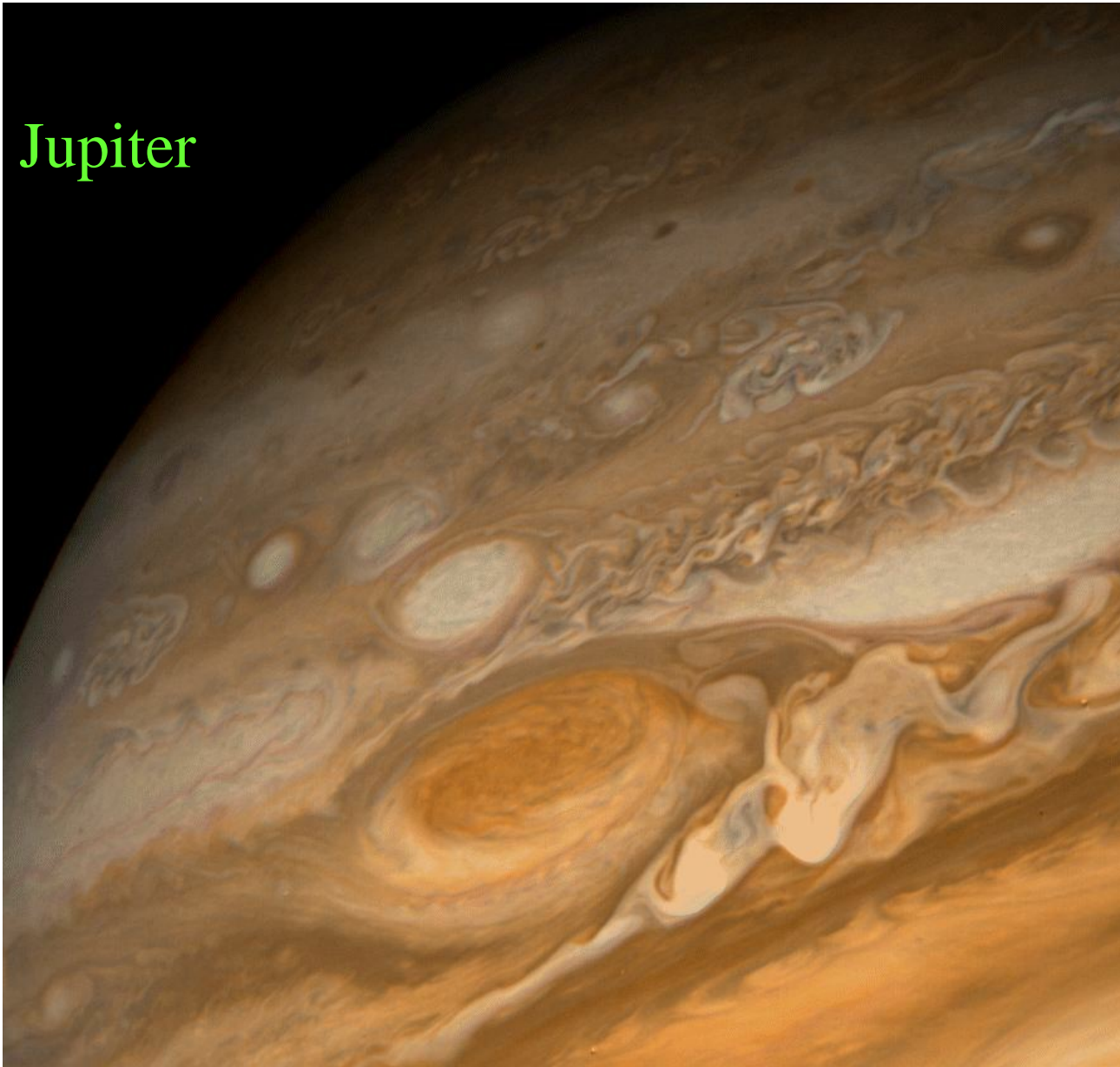


FIG. 4. Time-mean zonal-mean zonal wind profiles for cases I–VIII in Table 1 (the eight open circles in Fig. 3). Each grid on the abscissa represents 1 m s^{-1} .

**These alternating
Easterly and westerly
Jets are similar to
observed patterns
on Jupiter and
Saturn.**

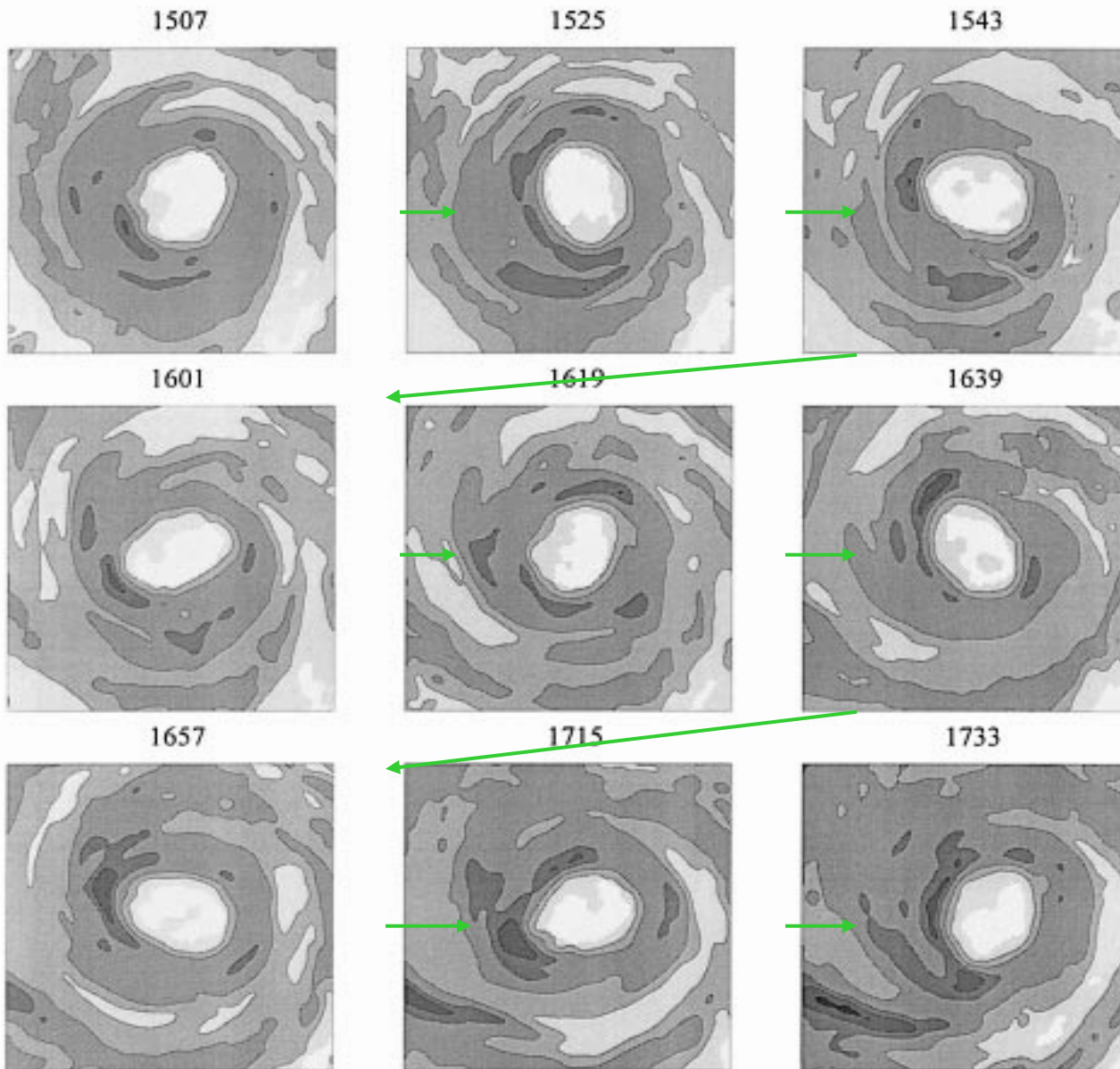
Rotational period 9.84hr

Jupiter



Alternating Zonal Structures

The Great Red Spot



144 min rotation period

**Coherent vortex
Dispersion resistant**

Vorticity Dynamics

Nonlinear glue

Nonlinear Dynamics

FIG. 1. Horizontal distribution of maximum reflectivity in the vertical column for Typhoon Herb from the Central Weather Bureau WSR-88D (10 cm) radar at Wu-Feng Mountain on 31 Jul 1996. The sequence of images is from left to right and from top to bottom. The time interval between each image is approximately 18 min. The local time of observation is indicated on top of each image. The major axis radius in the eye region is about 30 km and the minor axis radius is about 20 km. The nine images illustrate one eye rotation period of 144 min.

● Kirchhoff vortex (nonlinear) Lamb, 1932

$$\frac{x^2}{a^2} + \frac{y^2}{b^2} = 1$$

$$\zeta \frac{ab}{(a+b)^2} = \omega$$

$$\text{rotating period } \mathbf{P} = \frac{2\pi (a+b)^2}{\zeta ab}$$

$$a \sim 30 \text{ km} \quad b \sim 20 \text{ km}$$

$$P \sim (2\pi/\zeta)*4 \quad \zeta \sim 3*10^{-3} \text{ s}^{-1}$$

$$V_{\max} \sim 50 \text{ ms}^{-1}$$

● Kelvin PV wave (linear)

$$c = V_{\max} \left(1 - \frac{1}{m}\right) \quad m = 2$$

$$\text{angular velocity} = \frac{c}{r} = \frac{V_{\max}}{2r}$$

rotating period \mathbf{P}

$$\sim \frac{2\pi}{\omega} = \frac{2\pi}{2 \frac{V_{\max}}{r}} * 4 = \frac{2\pi}{\zeta} * 4$$

Same as Kirchhoff vortex !!

Eye rotation is nonlinear

$$J(\psi, \zeta) \gg$$

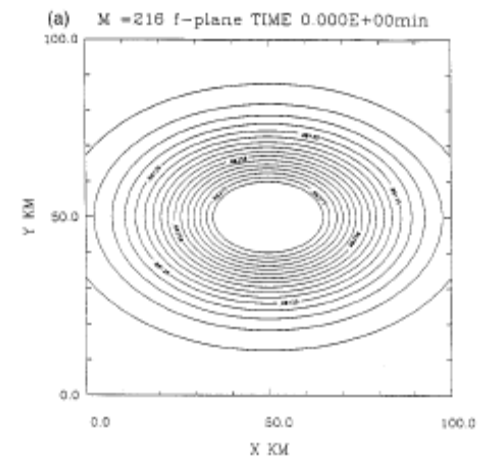
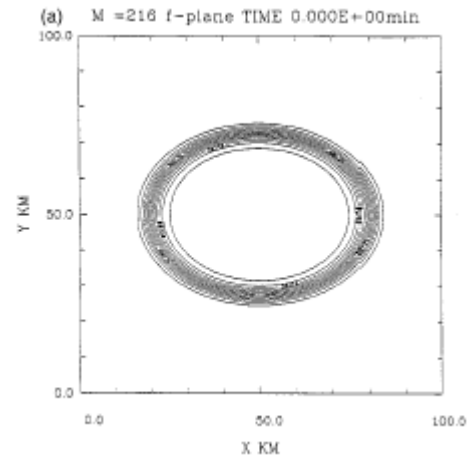
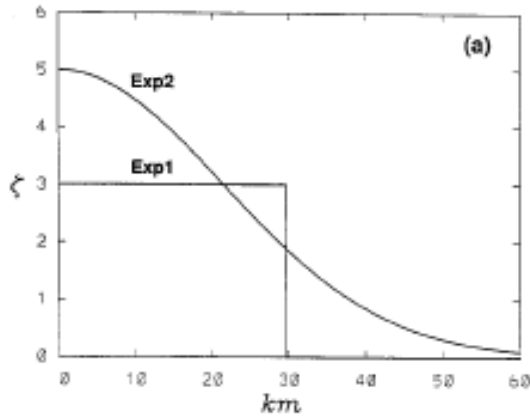
EXP 1

T = 0 min

$$J(\psi, \zeta)$$

EXP 2

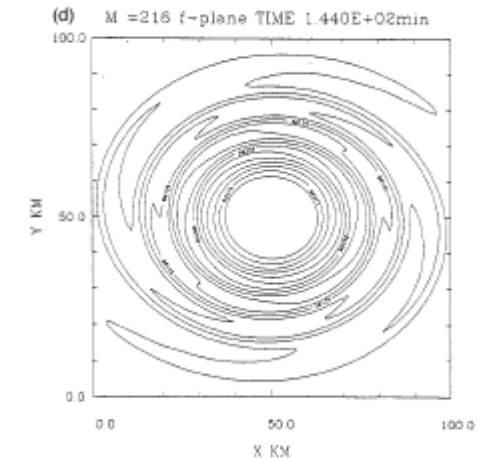
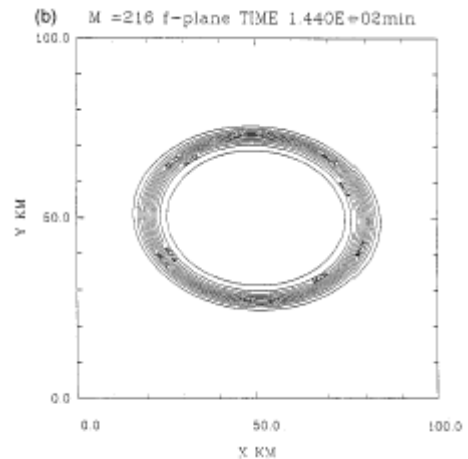
T = 0 min



T = 144 min

T = 144 min

$$\frac{\partial}{\partial t} \zeta + J(\psi, \zeta) = 0$$

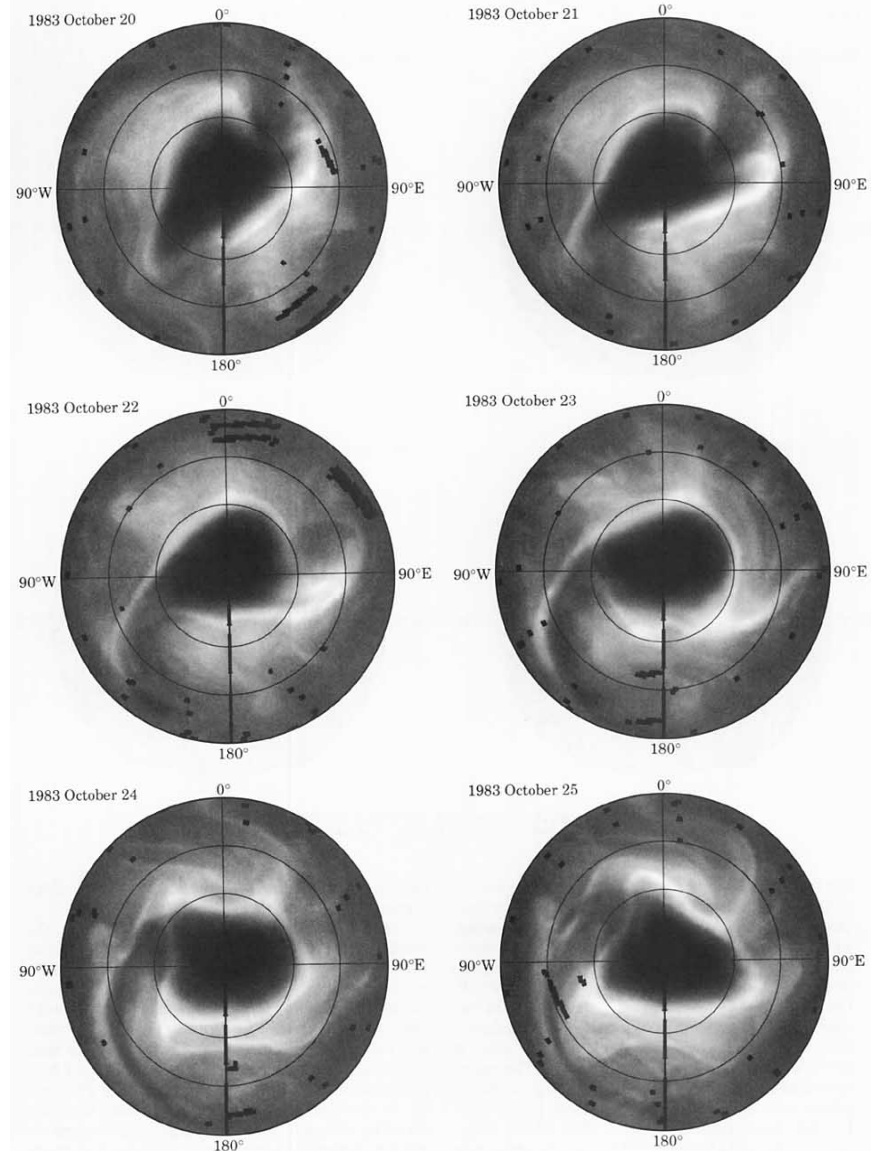


Bowman and Mangus (1993)

Observations of deformation and mixing of the total ozone field in the Antarctic polar vortex

Surf Zone Dynamics

Fig.1: Daily TOMS images of total ozone in the Southern Hemisphere for six consecutive days in October 1983. Latitude circles are drawn at 40° , 60° , and 80° S. The outermost latitude is 20° S.



Electron density redistribution in experimental plasma physics

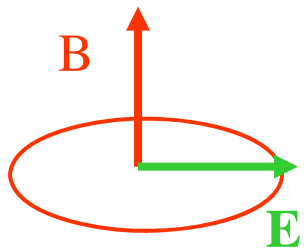
single sign charge

+

axial magnetic field
confinement

$$\mathbf{E} = -\nabla\psi$$

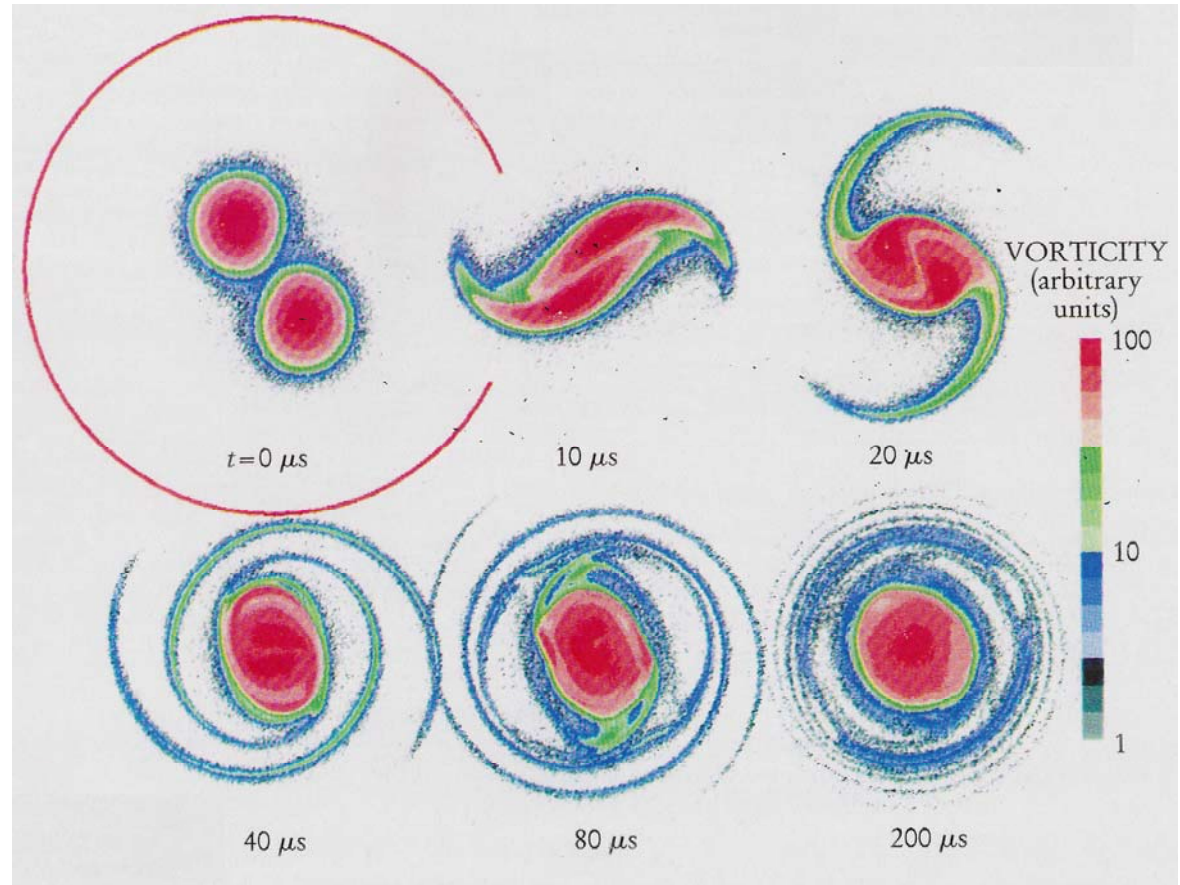
$$\nabla \cdot \mathbf{E} = -\nabla^2\psi = \frac{\rho}{\epsilon}$$



$\vec{E} \times \vec{B}$ drift

Coriolis force

Axisymmetrization 軸對稱化



Core is protected, thin filaments from edges

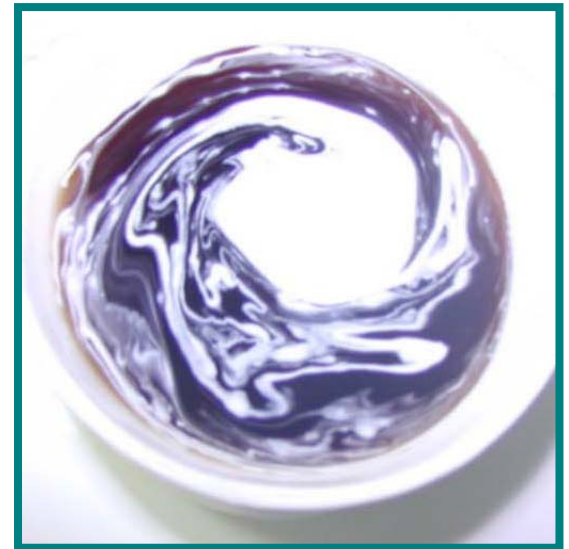
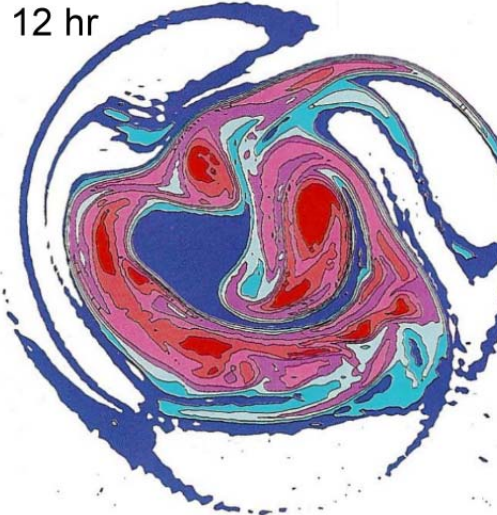
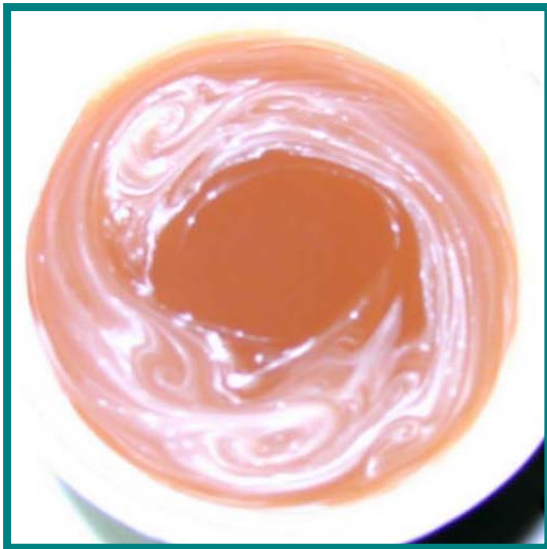
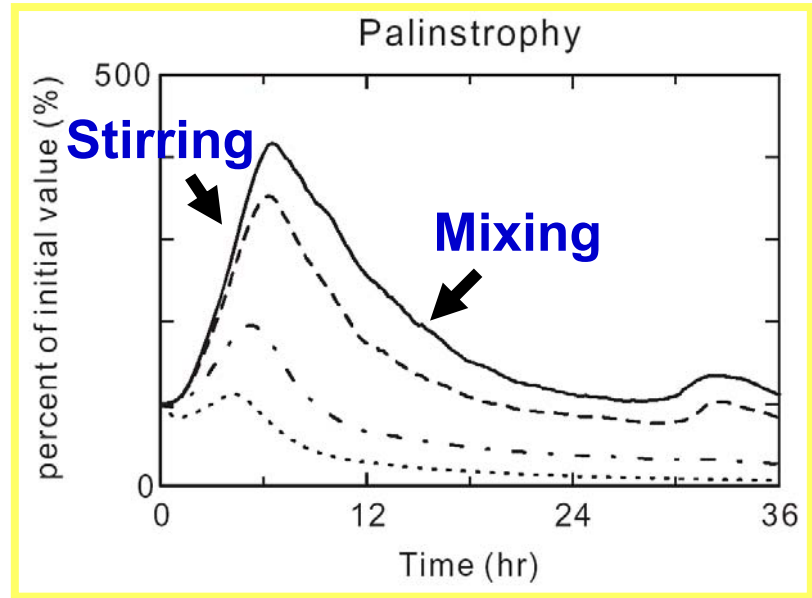
$$\frac{D\theta}{Dt} = \frac{\partial\theta}{\partial t} + \vec{V} \cdot \nabla\theta = \nu \nabla^2\theta$$

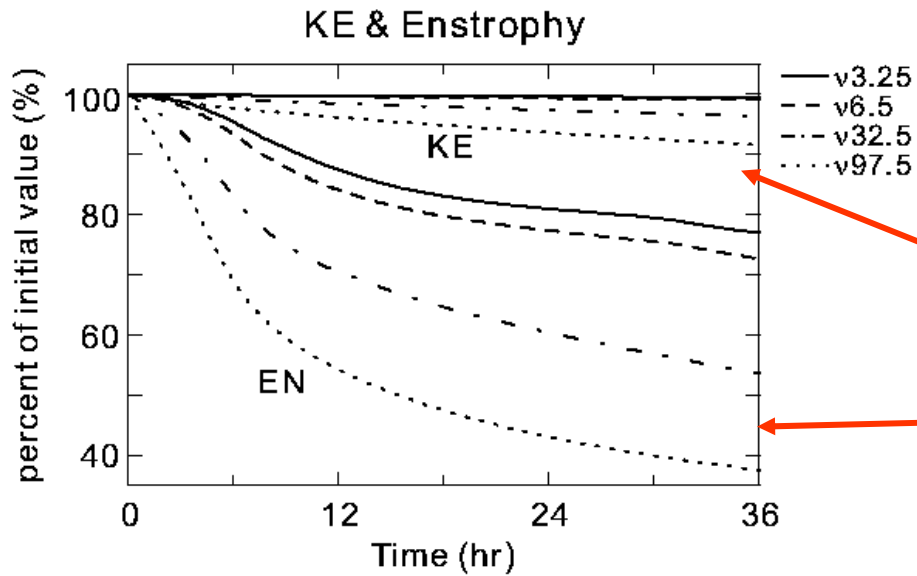
$$C = \frac{1}{2} \int \nabla\theta \cdot \nabla\theta dV$$

$$\frac{dC}{dt} = \int (\vec{V} \cdot \nabla\theta) \nabla^2\theta dV - \nu \int (\nabla^2\theta)^2 dV$$

Stirring

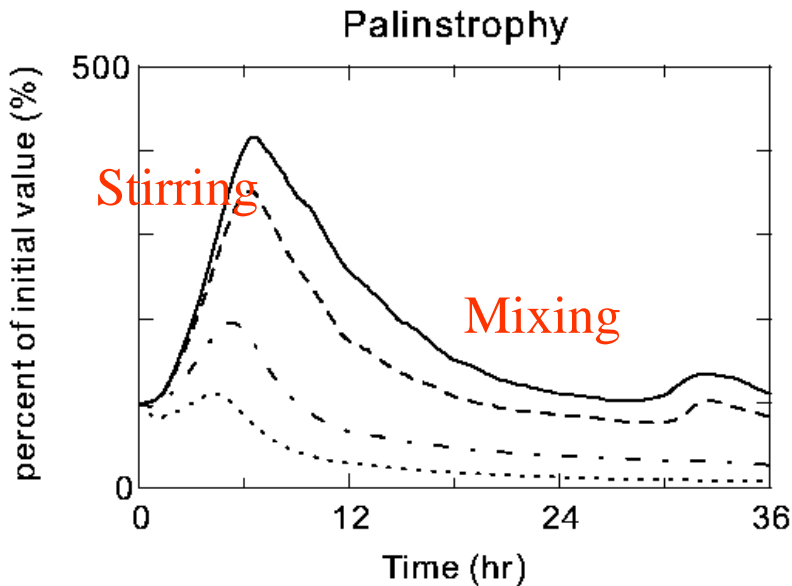
Mixing





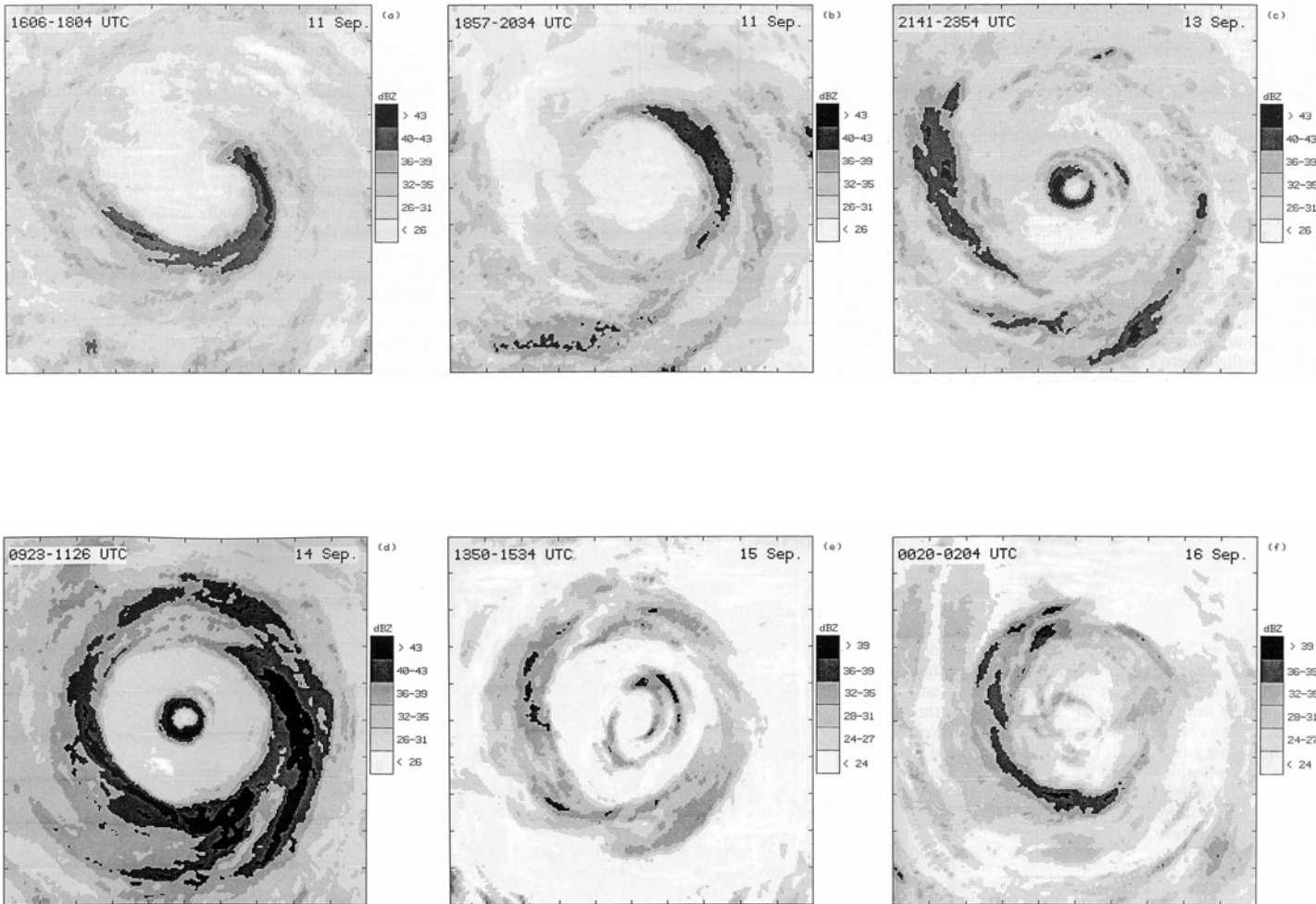
KE nearly conserved

Enstrophy cascade

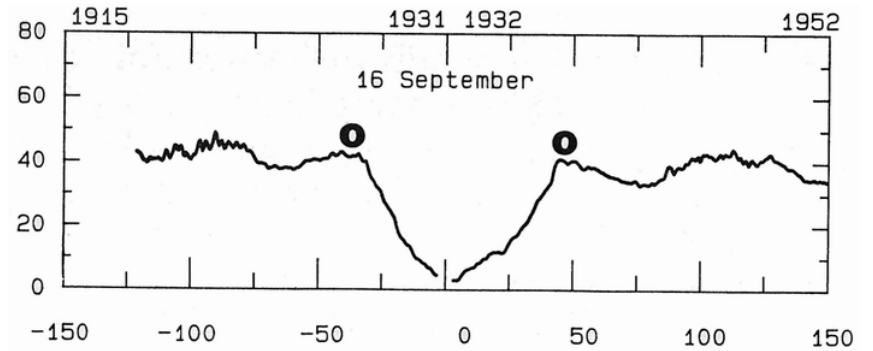
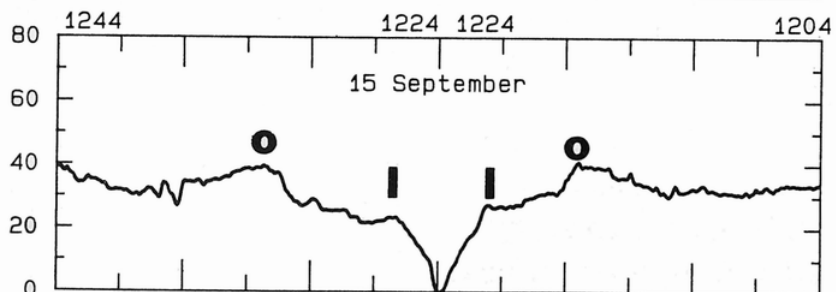
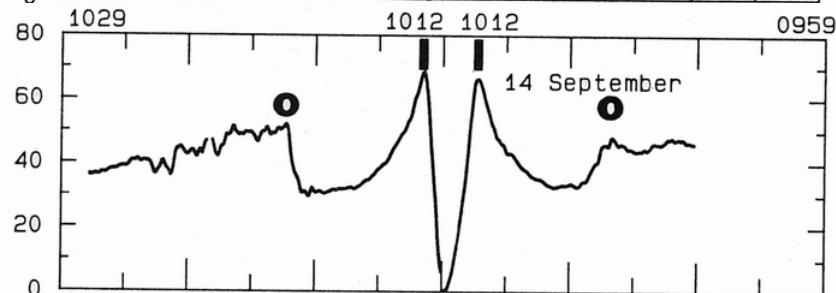
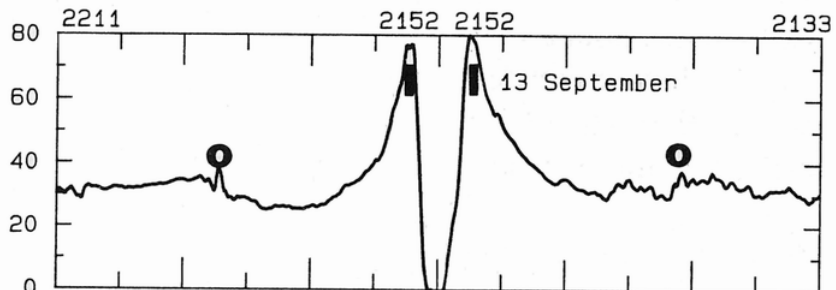
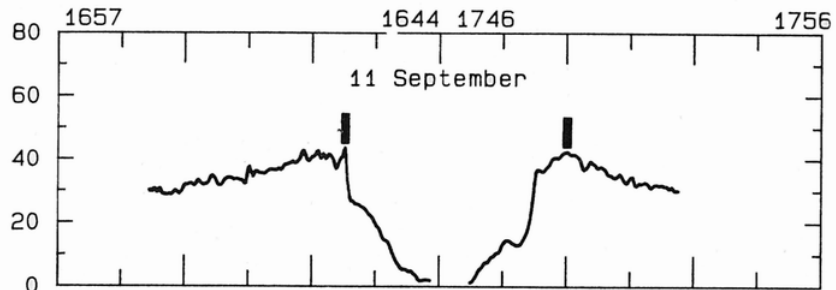


Selective Decay of 2D turbulence

Black and Willoughby (1992) Hurricane Gilbert (1988)



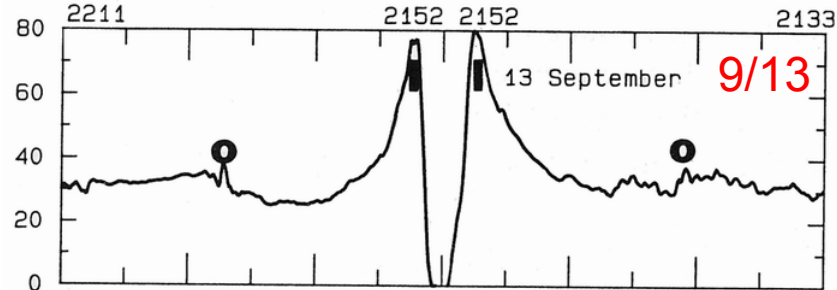
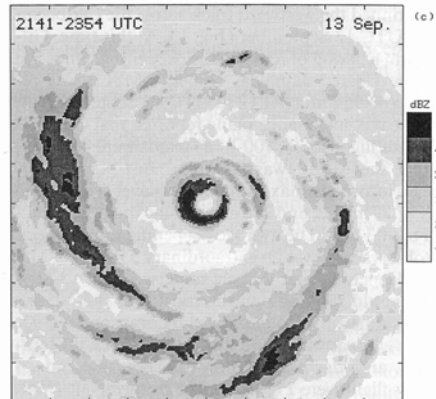
Black and Willoughby (1992)



Hurricane Gilbert (1988)

A major issue in understanding changes in typhoon intensity

Black and Willoughby (1992) Hurricane Gilbert (1988)



Development of symmetric structure from asymmetric convection in 12 hours

The contraction of the Outer tangential wind maximum

Core vortex intensity remains approximately the same during the contraction period

Inner core dissipate, TC weakens

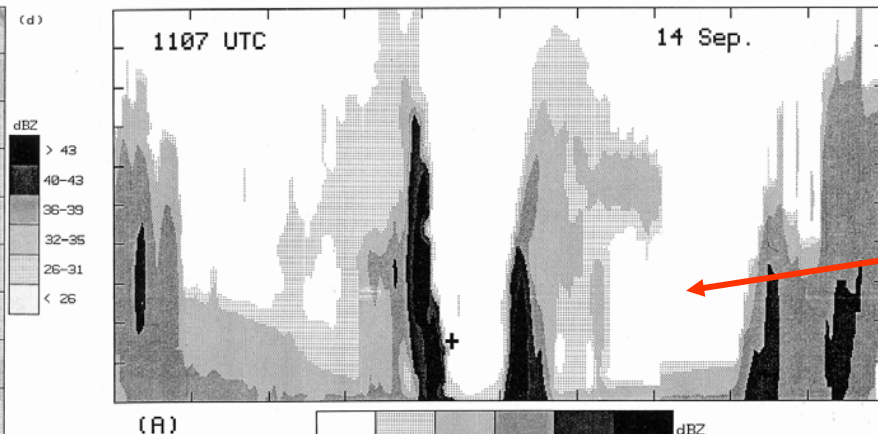
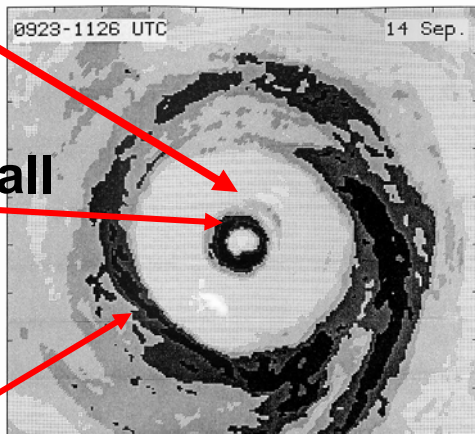
Black and Willoughby (1992)

Vertical cross sections of radar reflectivity of the concentric eyewall

Moat

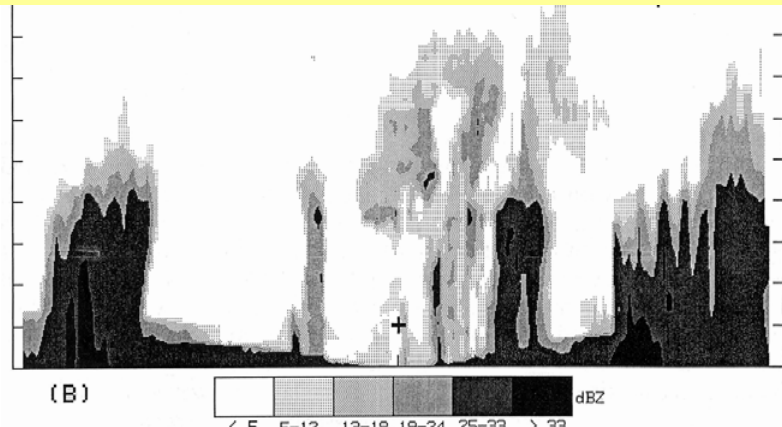
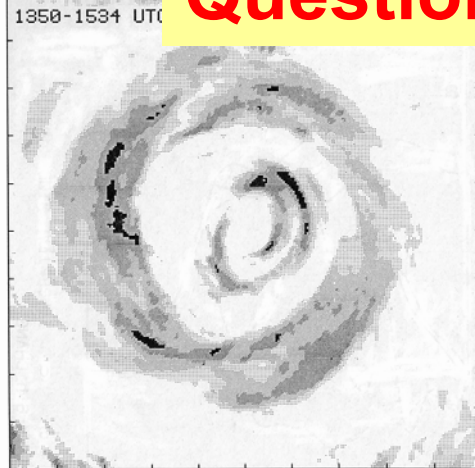
Inner eyewall

Outer eyewall

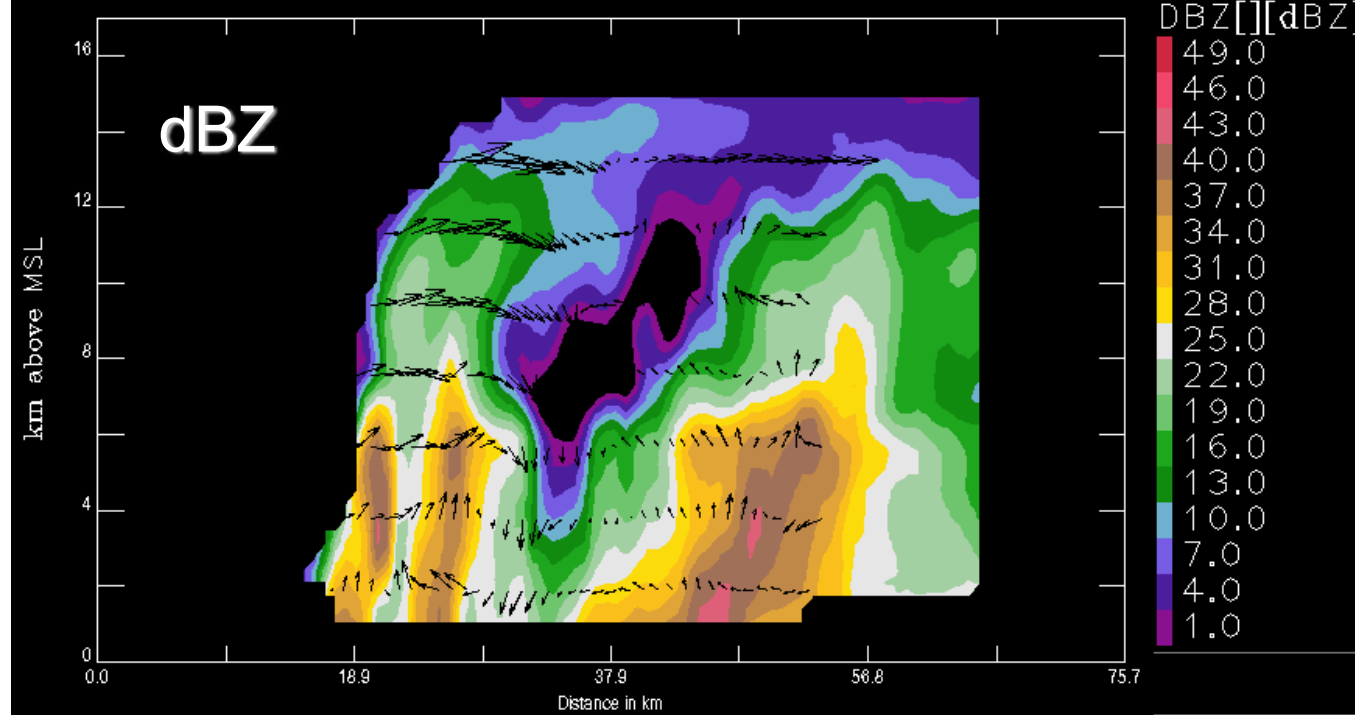


Moat

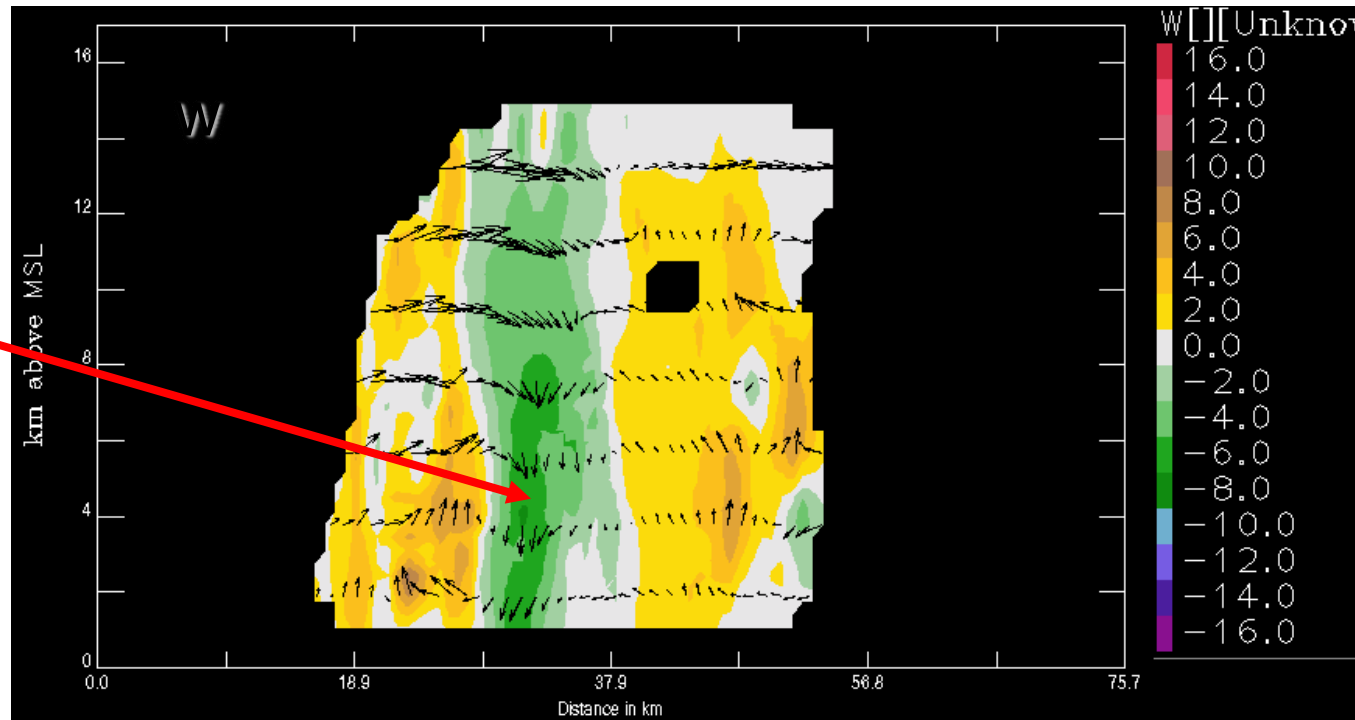
Question: How does the “moat form?”

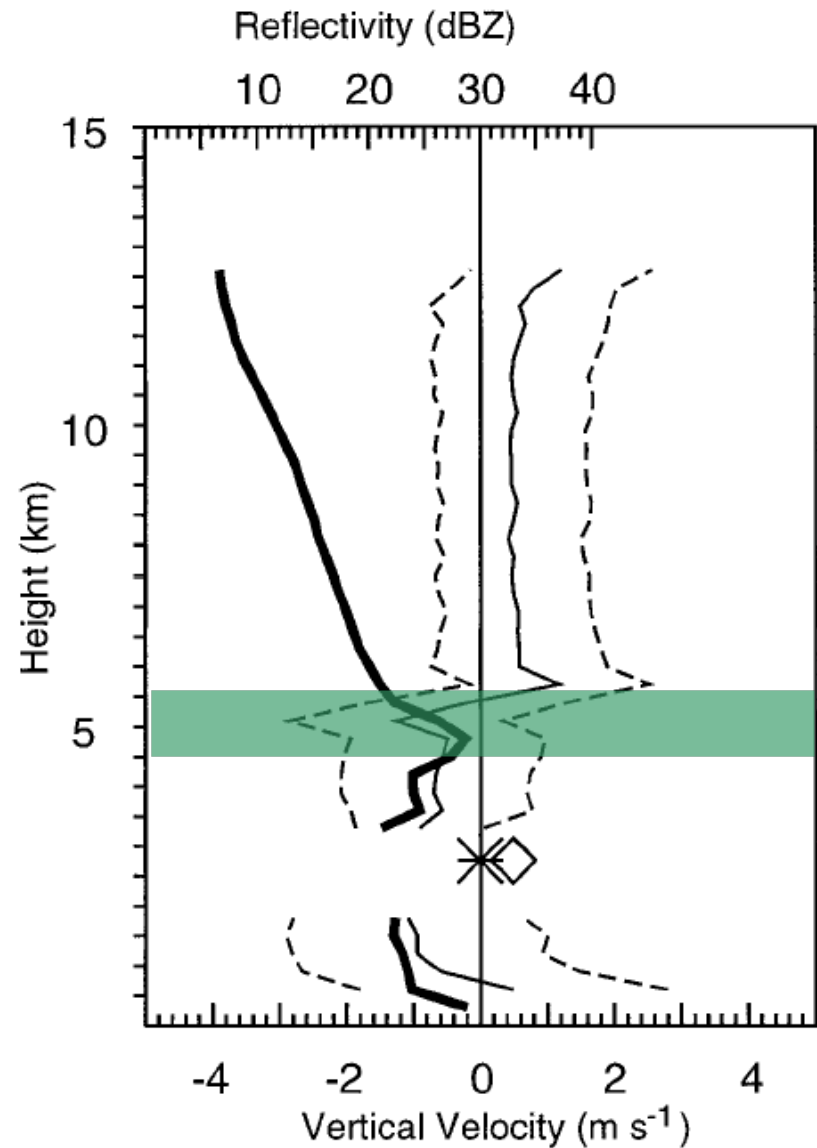


ELDORA data show downward motion between the two eyewalls



Subsidence





Hurricane Gilbert (1988)

The strong subsidence induced by the intensified eyewall convection may also contribute to the formation of the moat.

← bright band (4.5~5.5 km)

- reflectivity
- average vertical velocity
- - - one standard deviation on each side of the vertical velocity

FIG. 12. Vertical profiles of average vertical velocity (thinner line) and reflectivity (thicker line), for the stratiform region from 30 to 55 km from the center of circulation. The average flight-level vertical velocity is denoted by a star and the diamond denotes the average of the absolute value of the vertical velocity. The dashed lines are plotted one standard deviation on each side of the vertical velocity plot.

Dodge et al. (1999)

Thoughts from the 80's and 90's

Shapiro and Willoughby (1982) and Schubert and Hack (1982) proposed that heating-vorticity interaction can lead to convective-ring contraction.

$$d \zeta / dt \sim \zeta \nabla \cdot \mathbf{V}$$

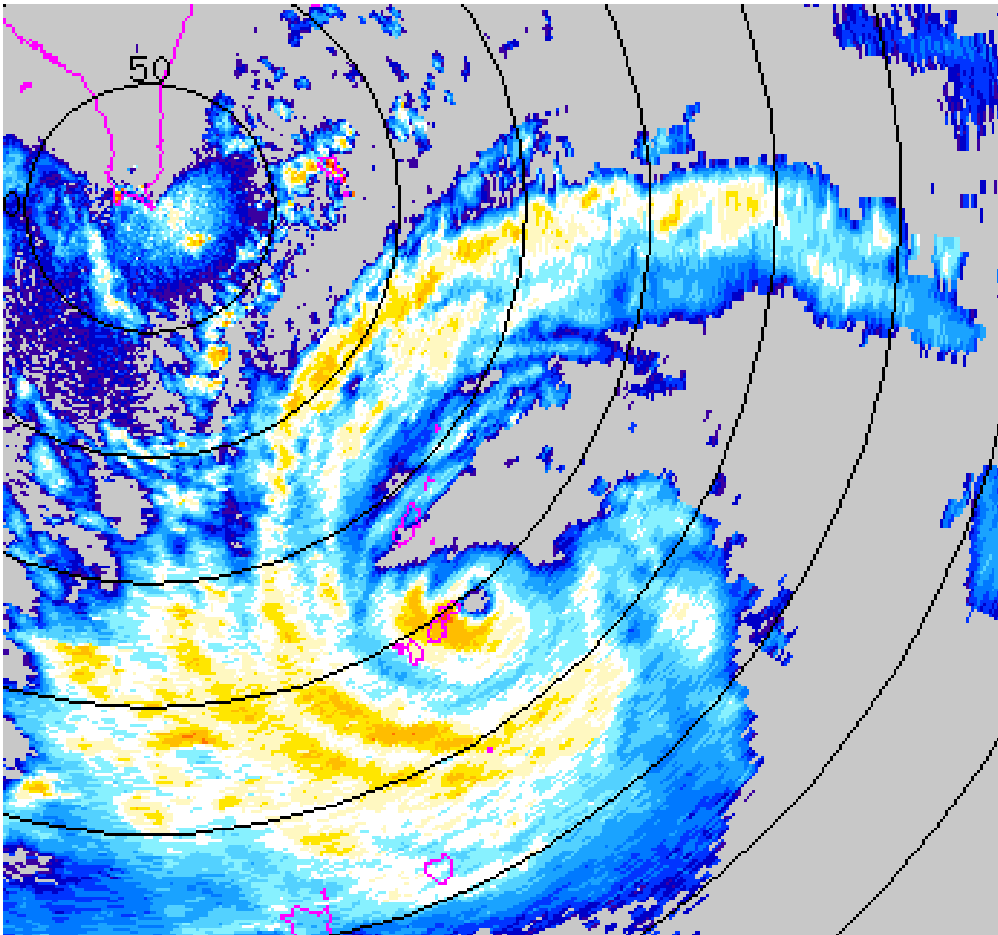
Stronger ζ near the TC core favors the inward response

Symmetrical Model

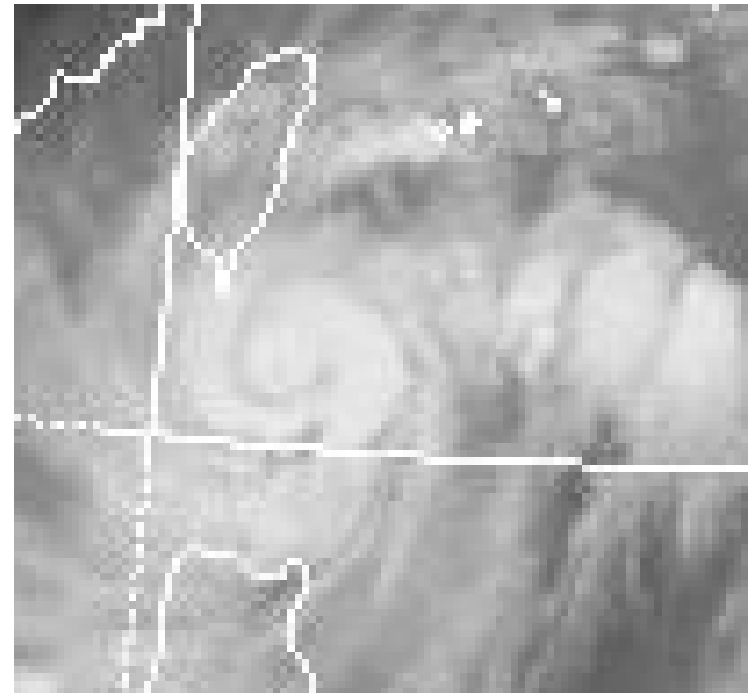
Moat formation and eyewall replacement are related to the subsidence and the moisture cut-off.

Typhoon Lekima (2001)

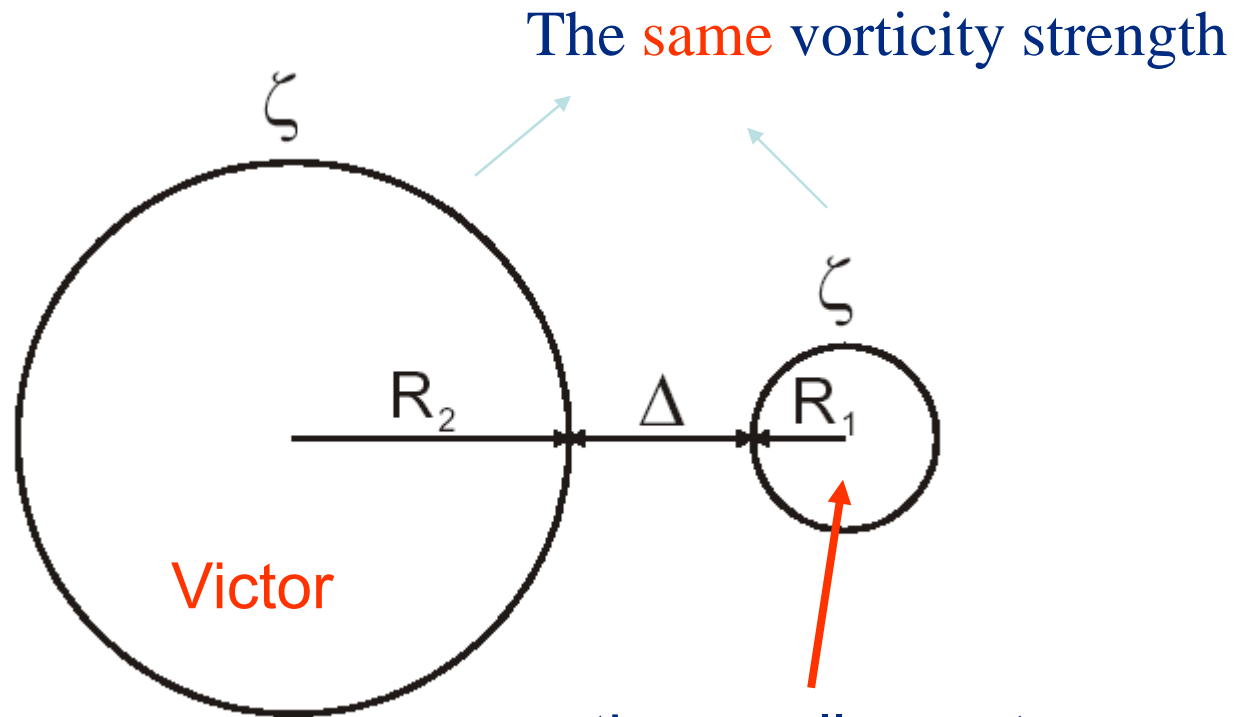
0935-1935 LST



0925 1900LST



Dritschel and Waugh (1992)



$$r = \frac{R_1}{R_2}$$

Vortex radius ratio

$$\frac{\Delta}{R_1}$$

Dimensionless gap

the smaller vortex was the one often being partially or totally destroyed

Elastic interaction regime



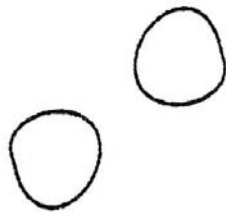
$t=0.0$



$t=2.0$



$t=3.0$



$t=4.0$

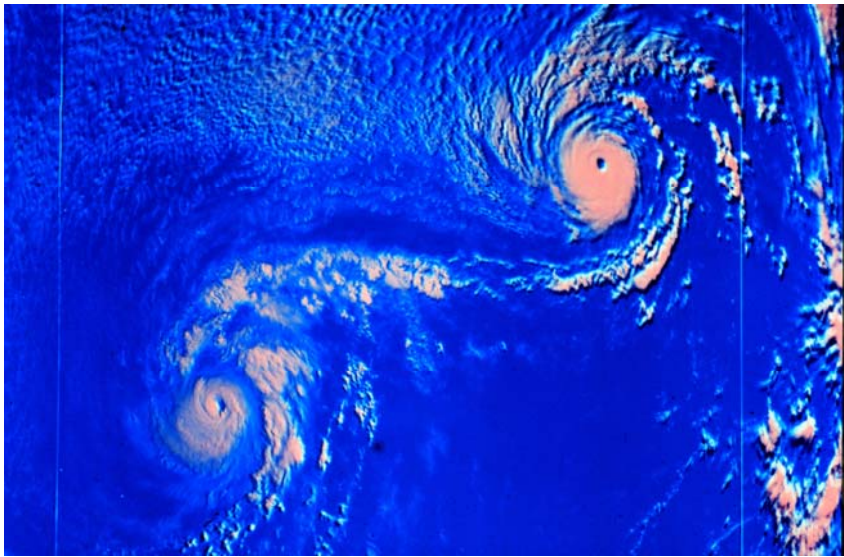


$t=6.0$

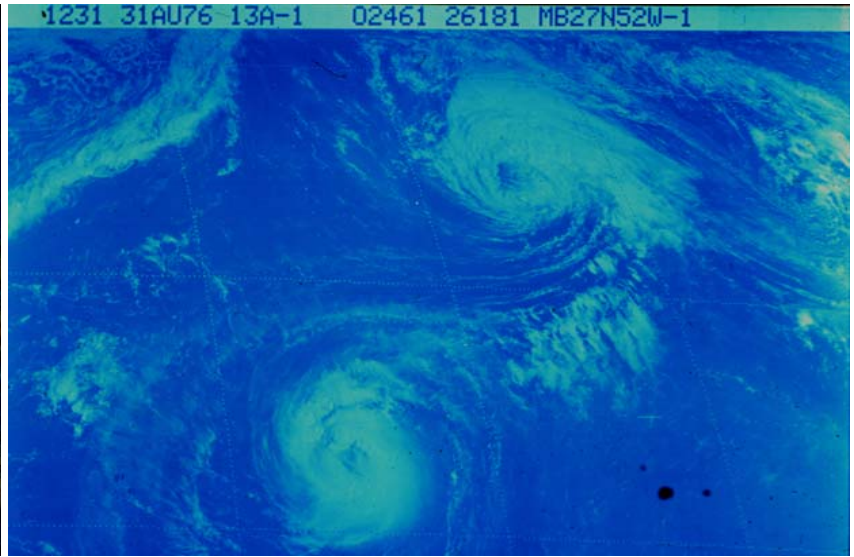


$t=20.0$

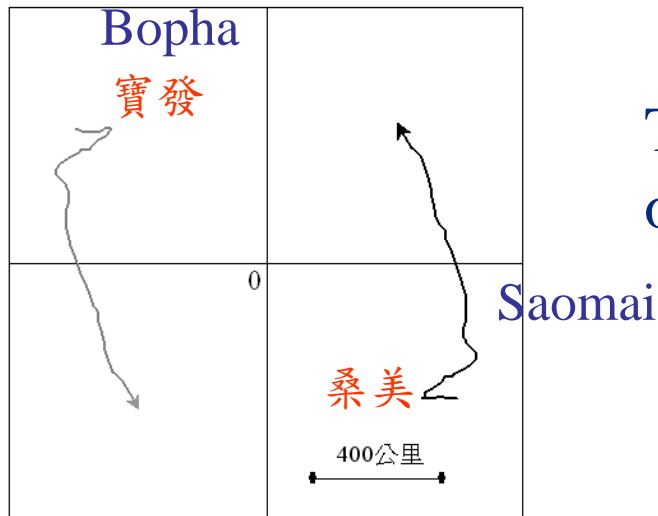
雙颱風的互繞 ---- 藤原效應



颱風 Ione 與 Kristen



颱風 Emmy 與 Frances

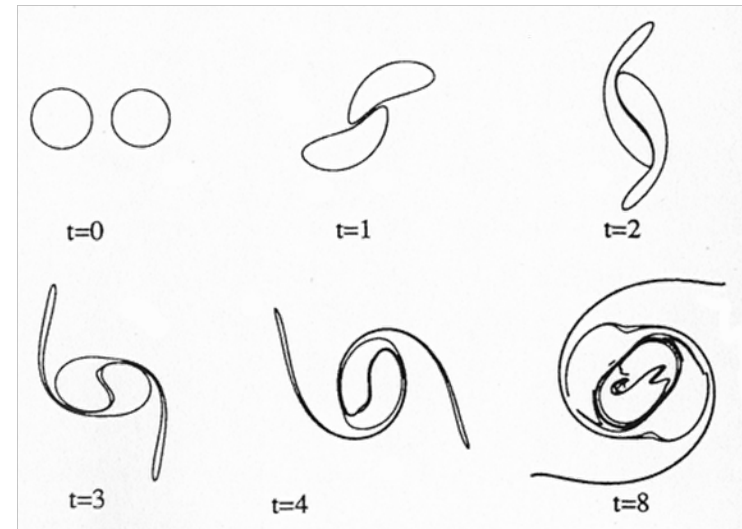
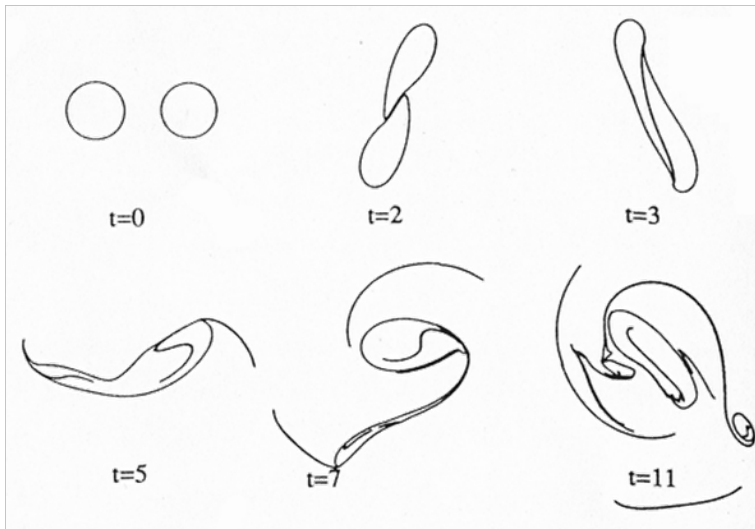


The unusual south movement of Typhoon Bopha

Merger regime

partial merger (PM)

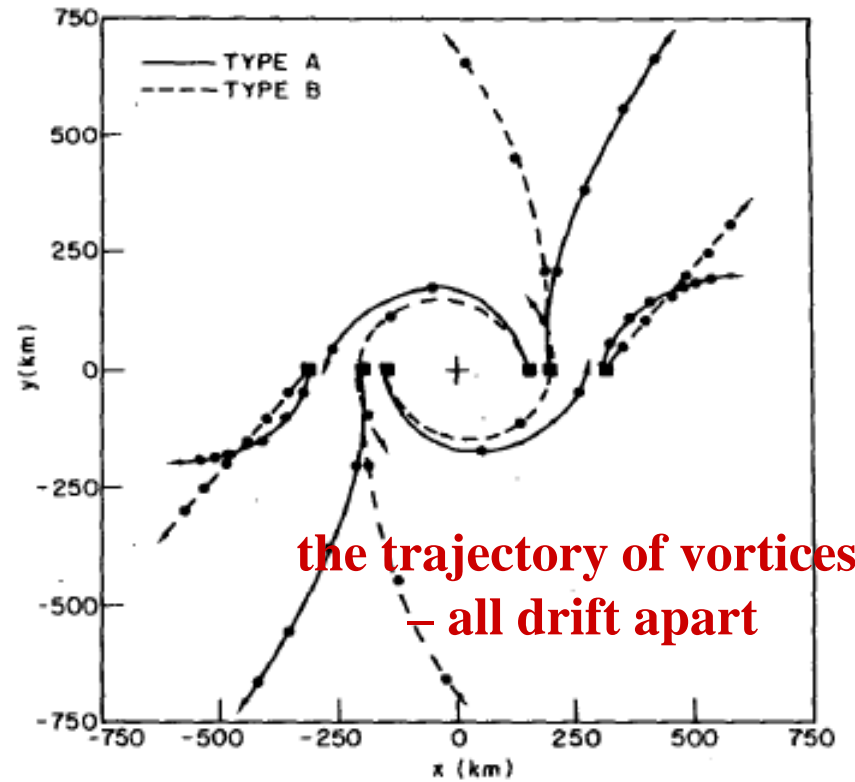
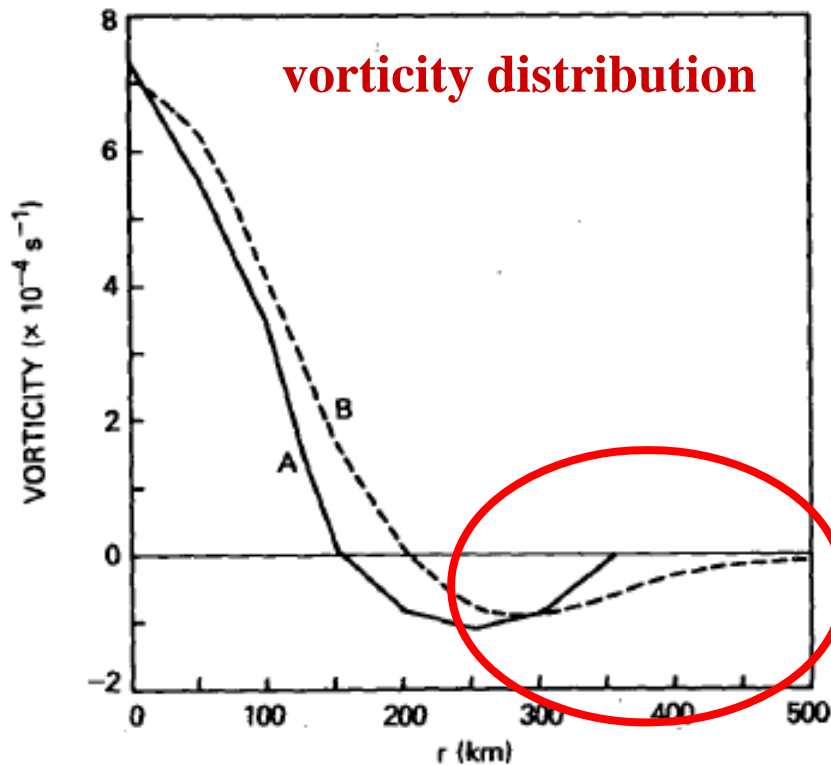
complete merger (CM)



Why 'merger' ?

- Chang (1983) – diabatic heating
- DeMaria & Chan (1984) – vortex vorticity gradient
- Dritschel and Waugh (1992)
 - advection + selective decay of 2D turbulence

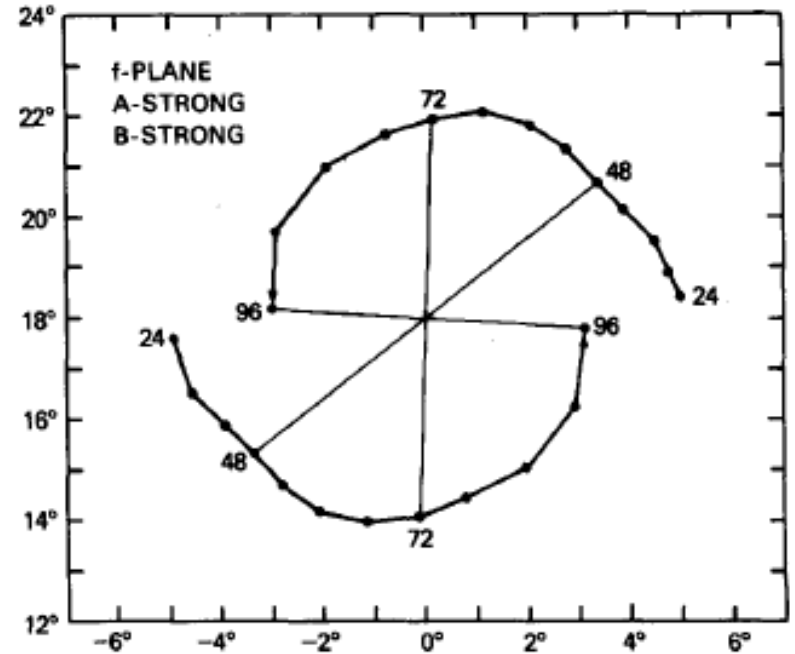
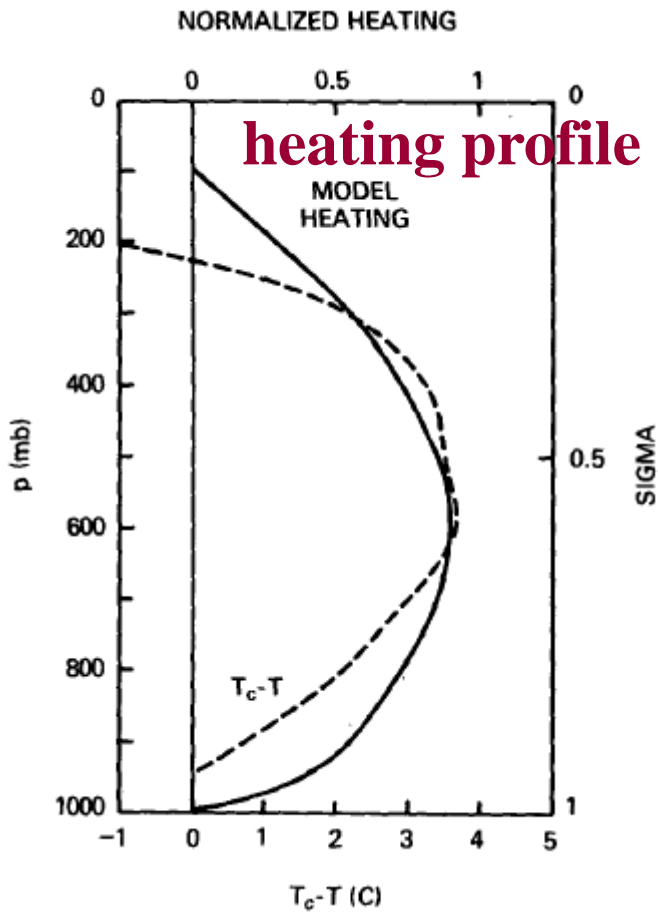
Chang (1983)



non-diabatic heating

Chang (1983)

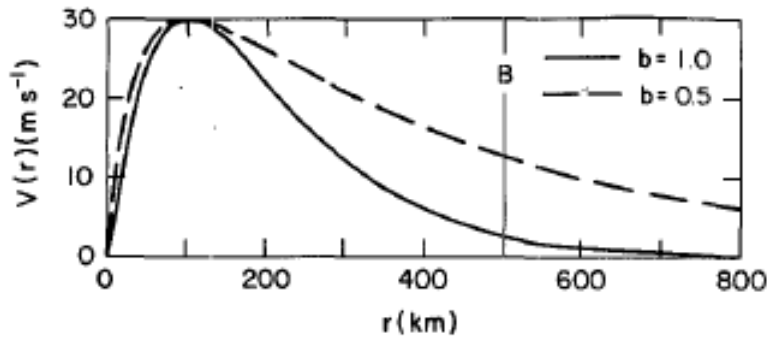
diabatic heating



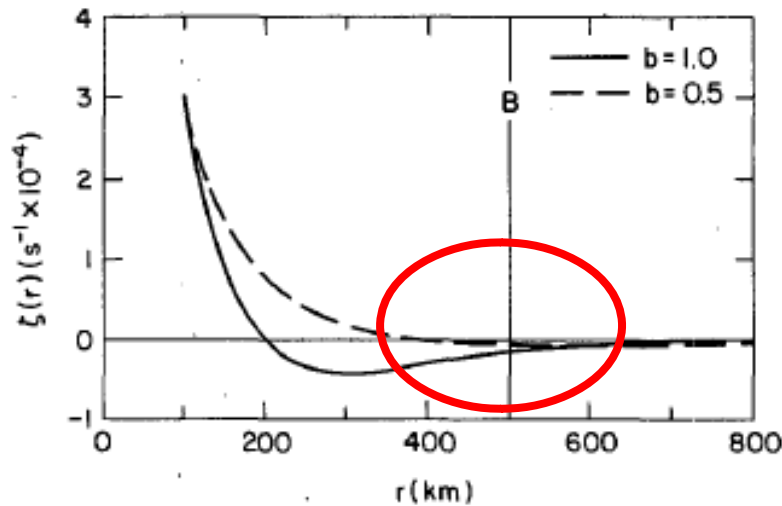
$$\frac{dQ}{dt} = 200 \text{ (K day}^{-1}\text{)}$$

The distance between two vortices decreases with time

DeMaria & Chan (1984)



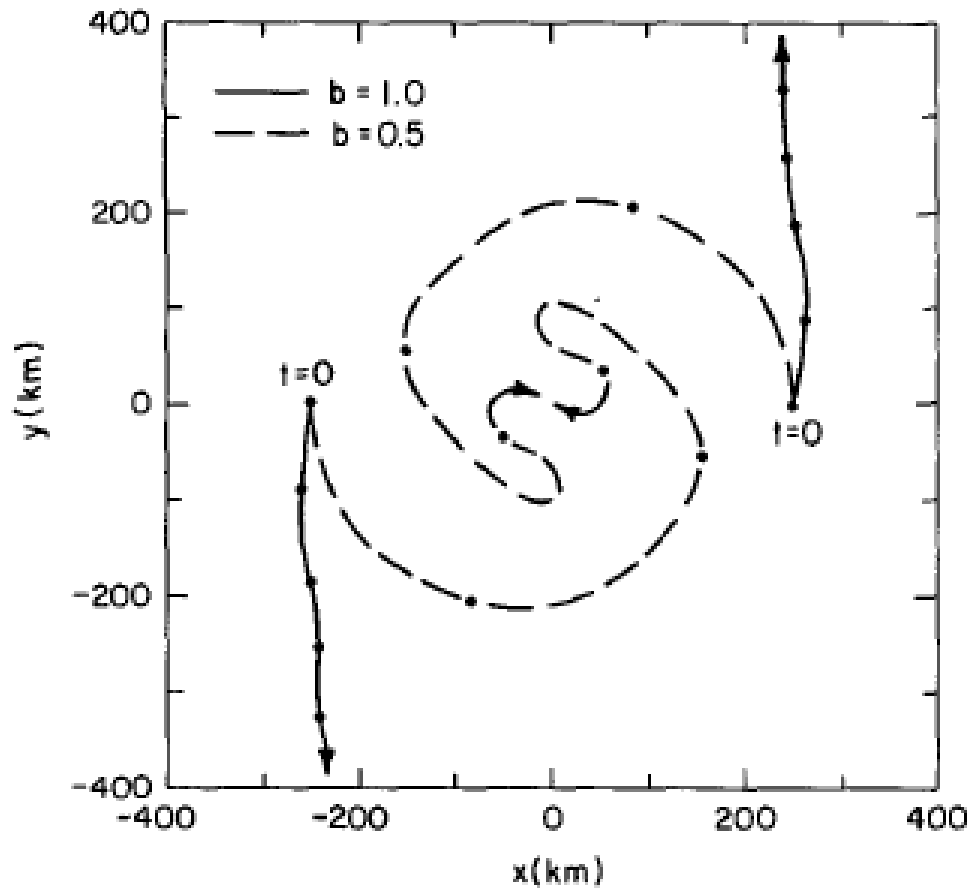
$$V(r) = V_m \left(\frac{r}{r_m} \right) \exp \left\{ \frac{1}{b} \left[1 - \left(\frac{r}{r_m} \right)^b \right] \right\},$$



$$\zeta(r) = \frac{2V_m}{r_m} \left[1 - \frac{1}{2} \left(\frac{r}{r_m} \right)^b \right] \exp \left\{ \frac{1}{b} \left[1 - \left(\frac{r}{r_m} \right)^b \right] \right\}$$

b : the factor determines the rate of tangential wind decays

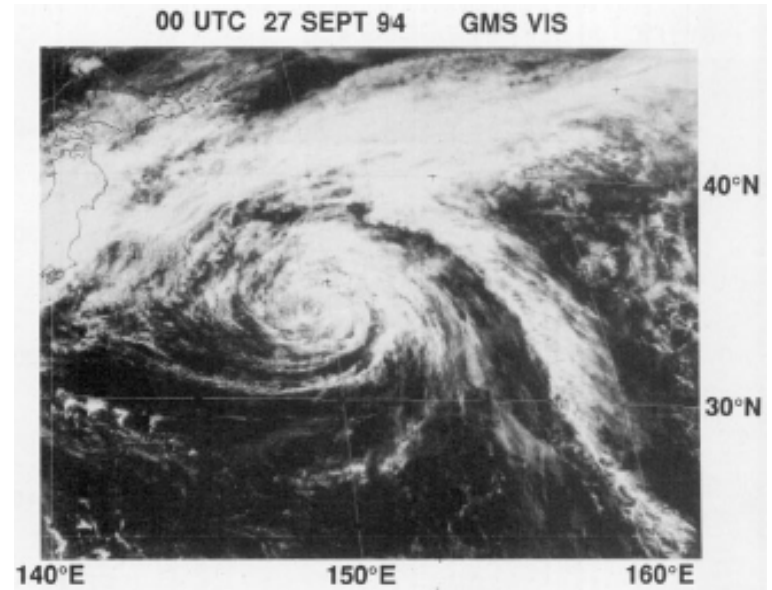
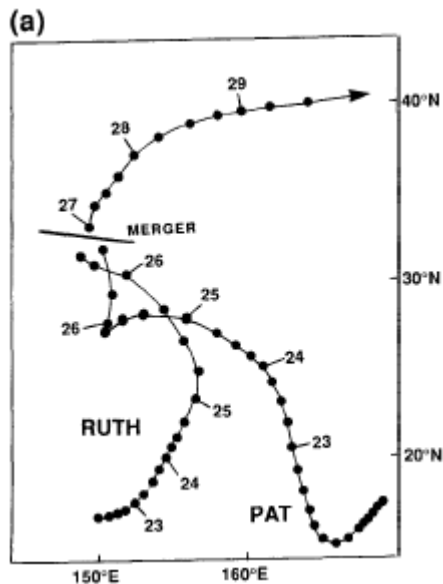
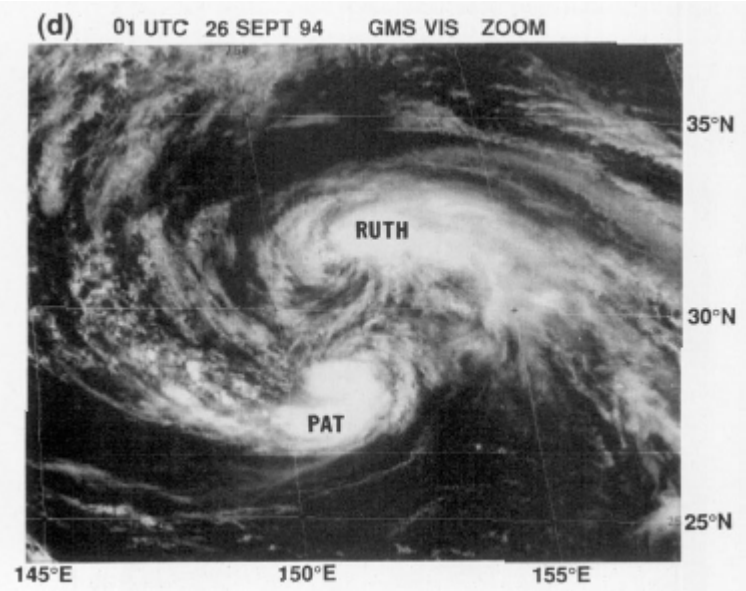
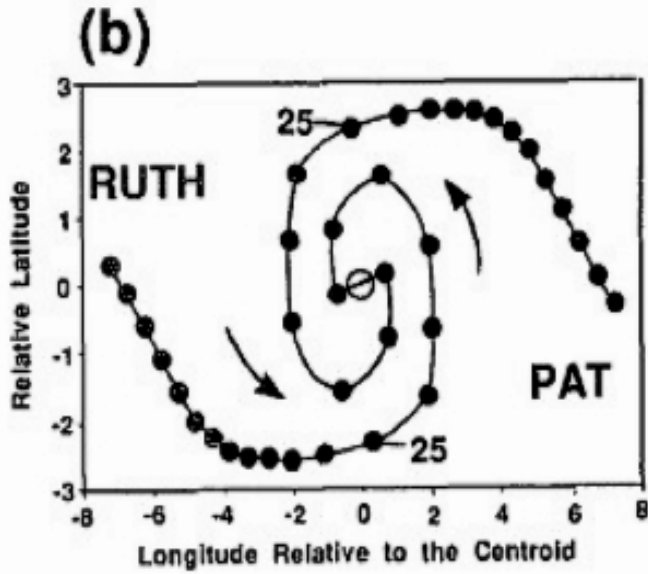
DeMaria & Chan (1984)



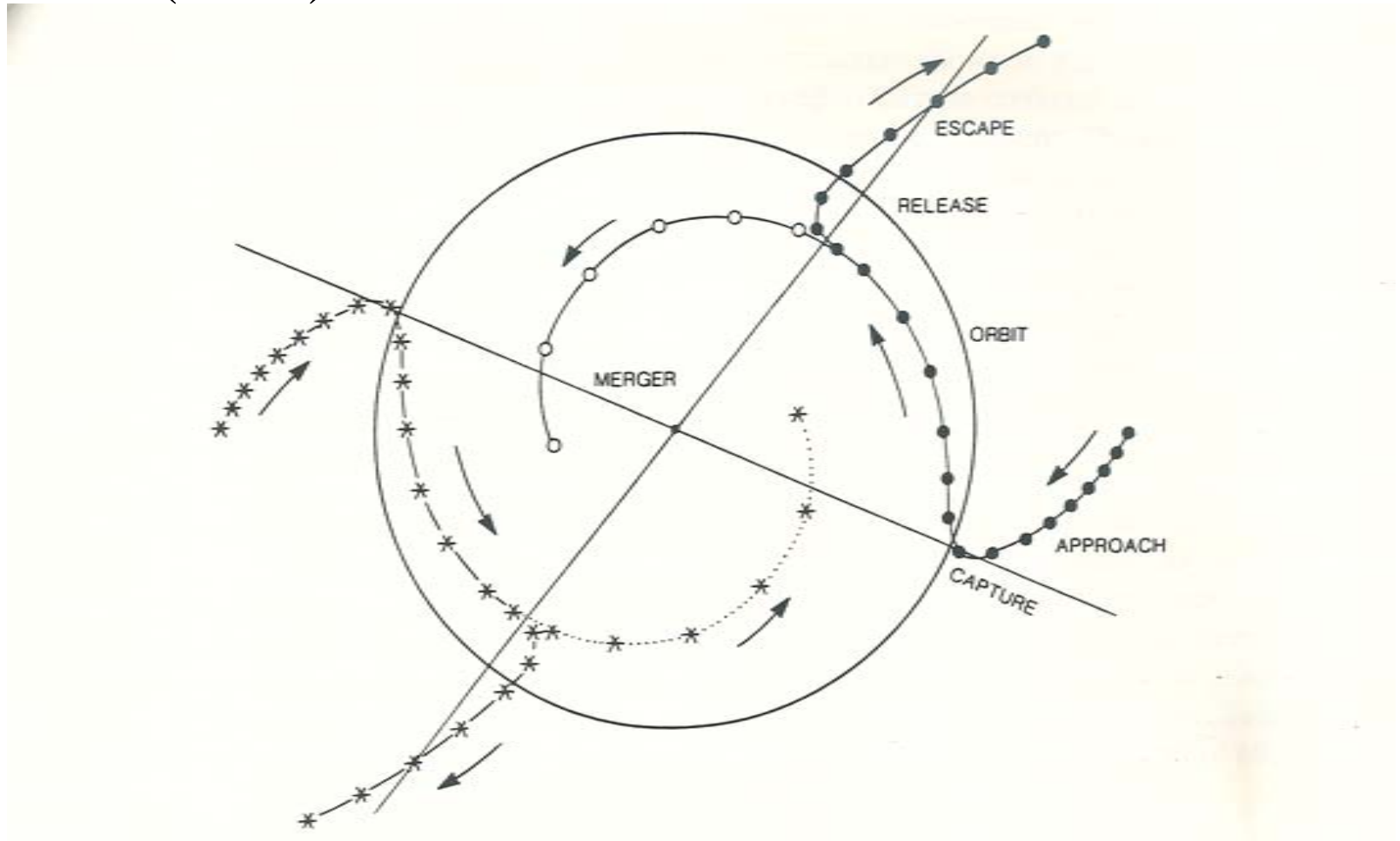
$b = 1.0$ drift apart

$b = 0.5$ merge

Merger ---- 颱風 PAT 與 RUTH (1994)



Lander and Holland (1993)

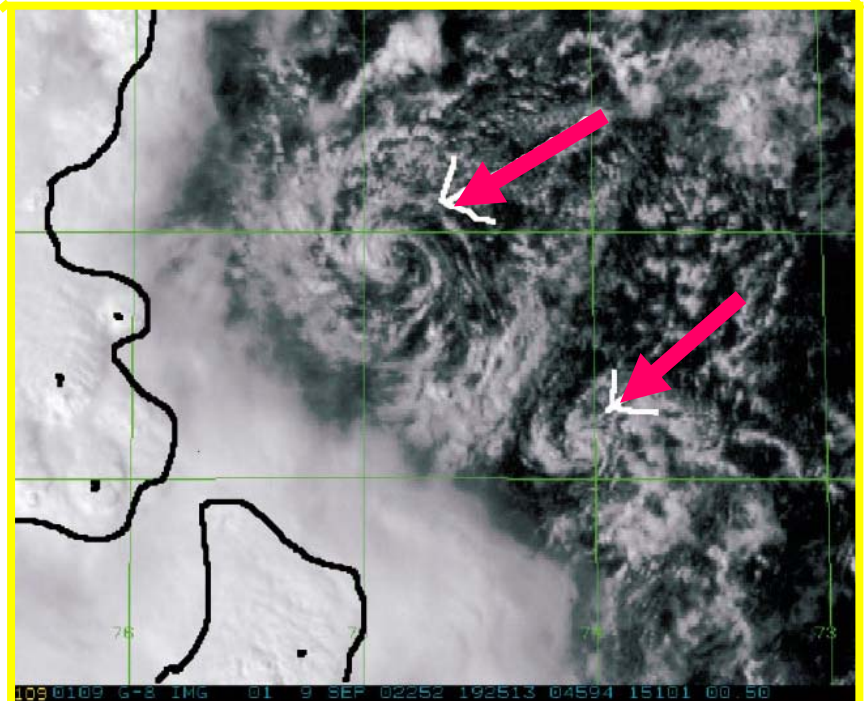
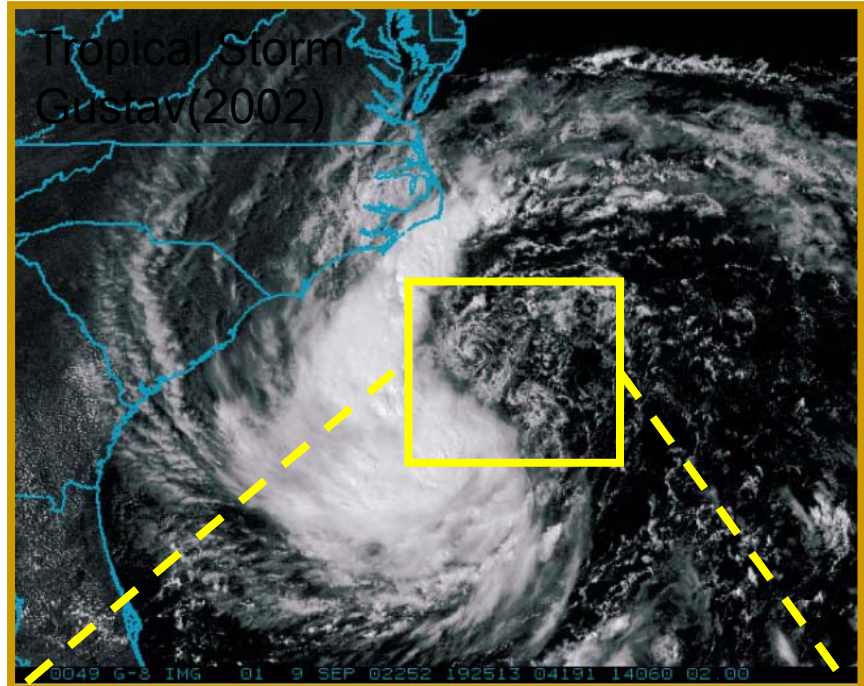


Merger and Elastic Interaction

Chaotic Behavior

“Vortical” Hot Towers

- The net effect of the hot towers is to produce strong small-scale (10km in diameter on average) lower-tropospheric (below $z \approx 5\text{km}$) cyclonic PV towers.
- The strong updrafts in the hot towers converge and stretch existing low-level vertical vorticity into **intense small-scale vortex tubes**.
- Multiple **mergers / axisymmetrization** of these tubes in the lower troposphere.



“Vortical” Hot Towers

t=26.16 hr

$$\frac{D(PV)}{Dt} \approx \frac{f + \zeta}{\rho} \frac{\partial \dot{\theta}}{\partial z}$$

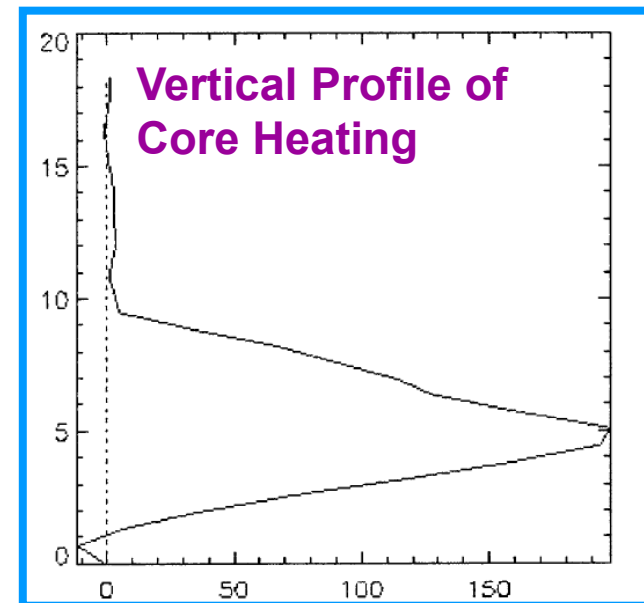
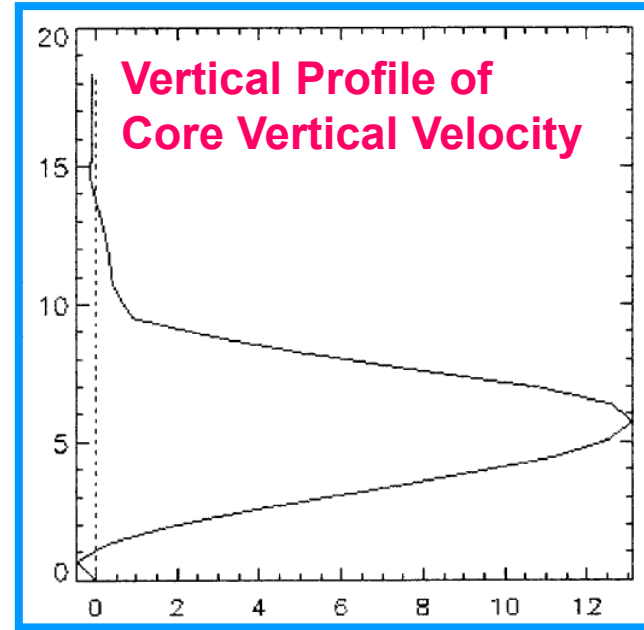
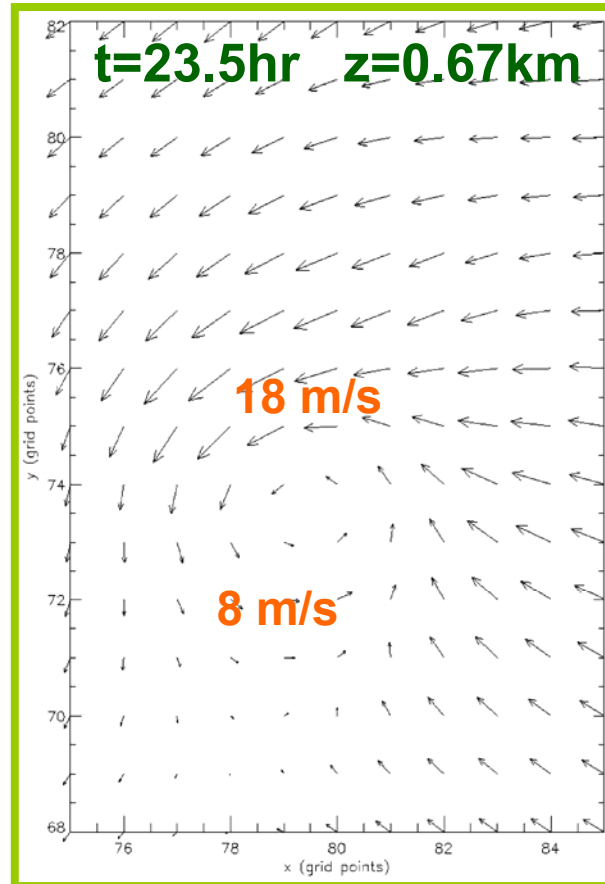


$$\approx \frac{10^{-3} s^{-1}}{0.7 kg m^{-3}} (1.25 \times 10^{-5} K s^{-1} m^{-1})$$

$$\approx 64 PVU h^{-1}$$

$$\left(\begin{array}{l} 1 PVU \\ = 1 \times 10^{-6} m^2 s^{-1} K kg^{-1} \end{array} \right)$$

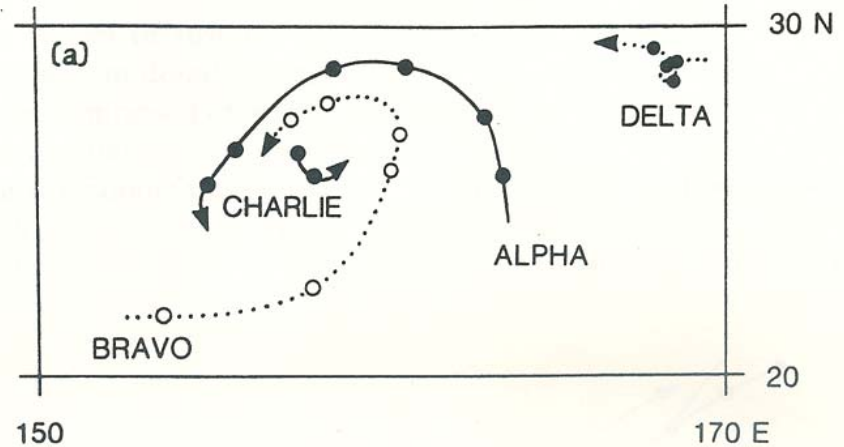
$$\left| \frac{\partial \dot{\theta}}{\partial z} \right| \approx 45 K h^{-1} km^{-1}$$



Hendricks et al.
(2004)

Lander and Holland (1993)

Tracks of mesoscale vortices



Centroid-relative motion of Alpha and Bravo

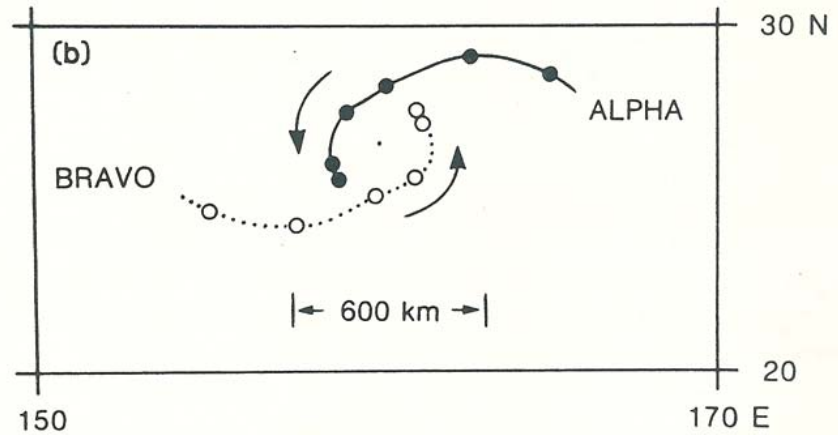
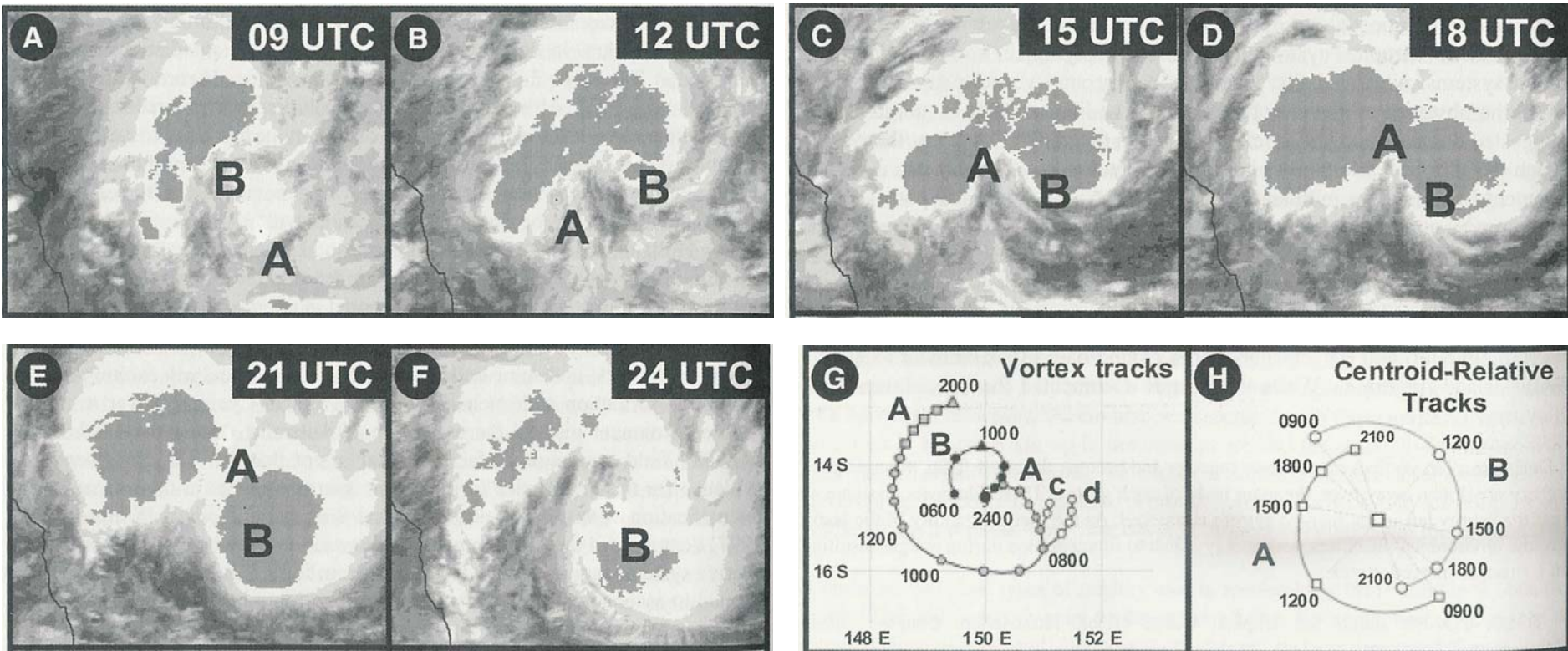


Figure 6. (a) Tracks of mesoscale vortices Alpha, Bravo, Charlie and Delta and (b) centroid-relative motion of Alpha and Bravo. Dots on the tracks are at irregular time intervals and show fixes obtained from visible satellite imagery.



(A)-(F) The locations of the two mesovortices A and B during the development of Tropical Cyclone Oliver superposed on three-hourly satellite imagery for the period 0900 UTC February 4 to 0000 UTC February 5, 1993; (G) Tracks of four of the vortices obtained from radar data. The positions are not evenly spaced and so times (in UTC) of some of the vortex positions are marked; (H) three-hourly centroid-relative tracks of mesovortices A and B from 0900 UTC to 2100 UTC February 4 [Simpson et al., 1997].

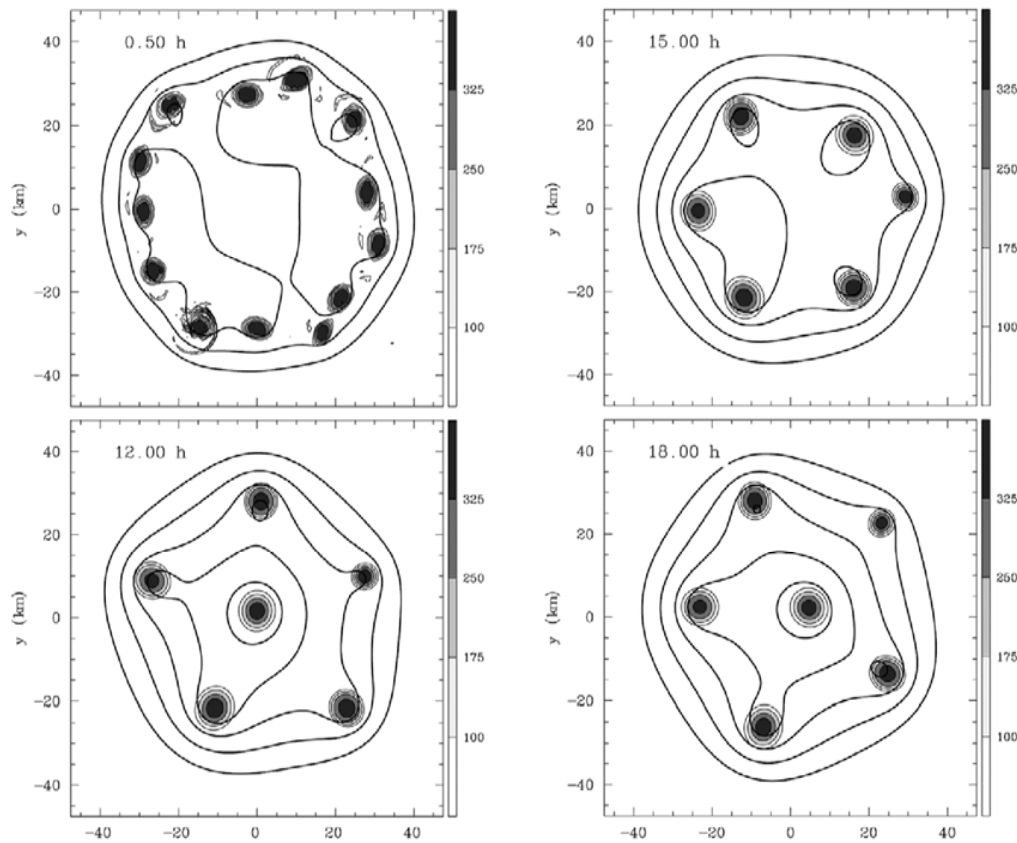


FIG. 2. Evolution of vorticity (shaded) and streamfunction contours (bold) for the numerical experiment of Kossin and Schubert (2001). Values along the label bar are in units of 10^{-4} s^{-1} . The shape of the streamlines transitions from a pentagon to a hexagon and back to a pentagon over 6 h.

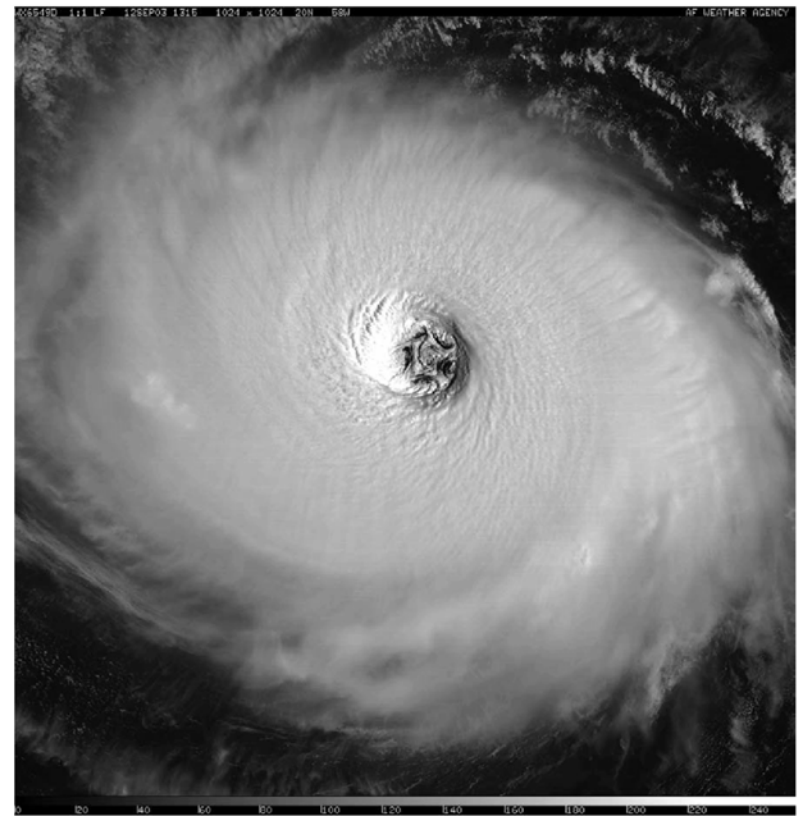


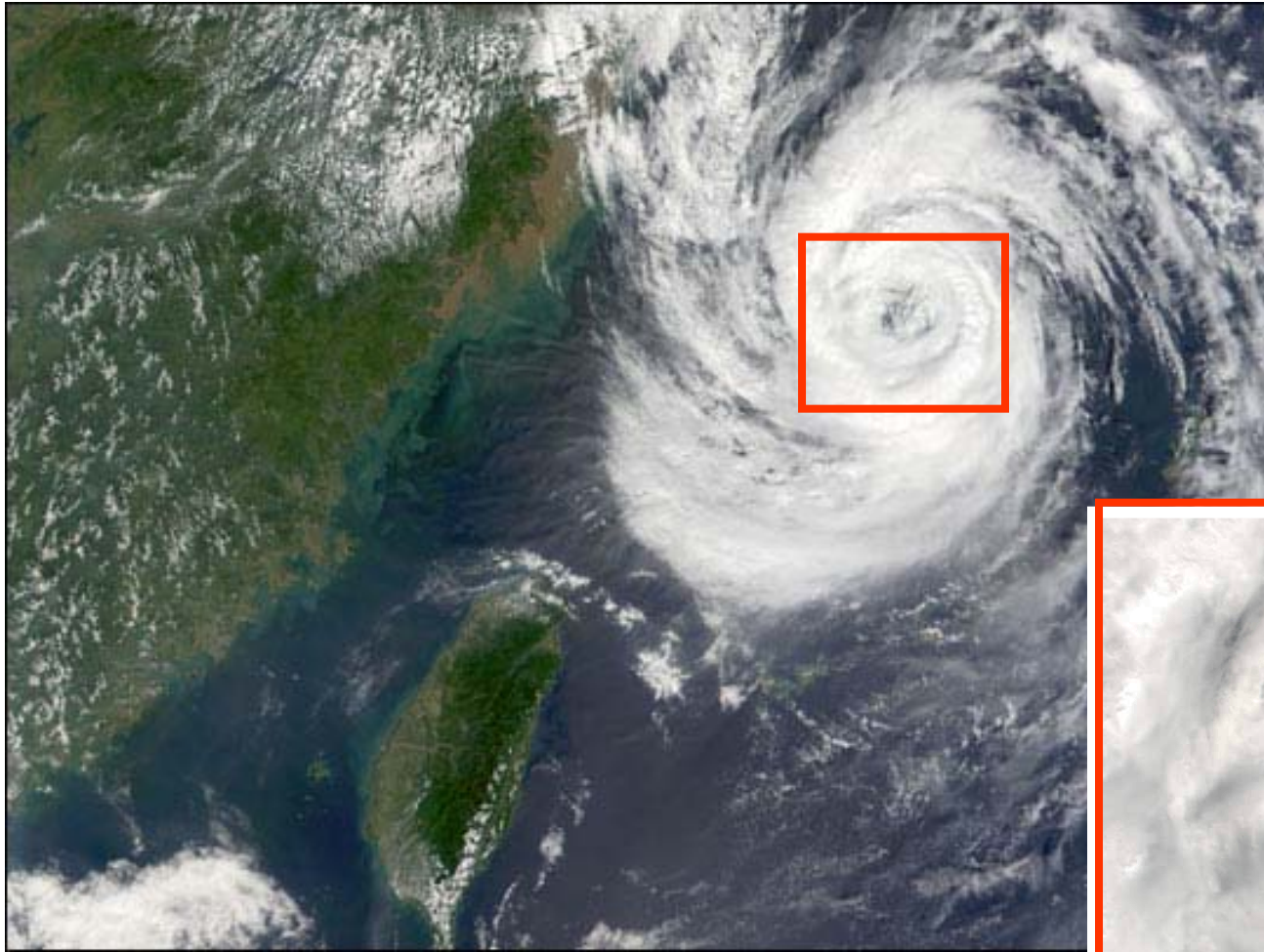
FIG. 1. Defense Meteorological Satellite Program (DMSP) image of Hurricane Isabel at 1315 UTC 12 Sep 2003. The starfish pattern is caused by the presence of six mesovortices in the eye—one at the eye center and five surrounding it.

MESOVORTICES IN HURRICANE ISABEL

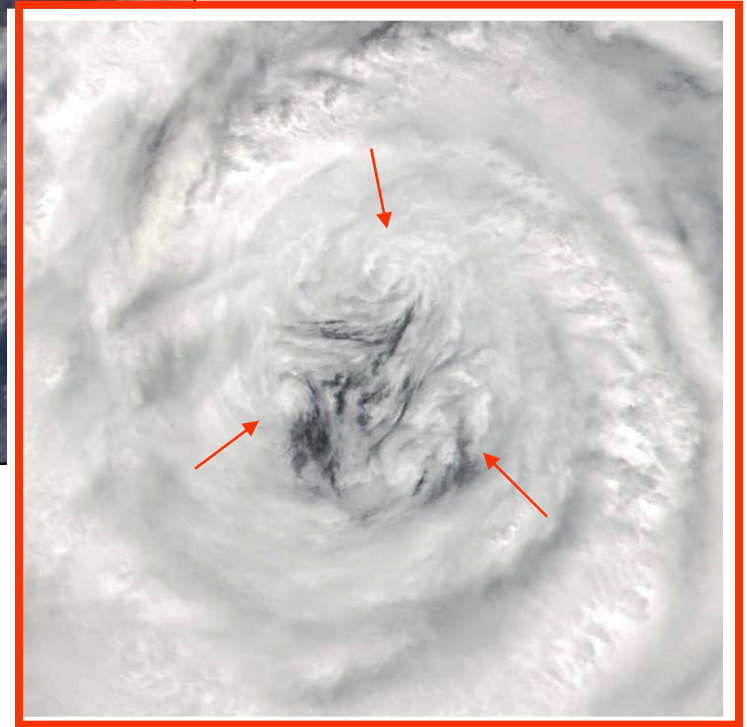
BY JAMES P. KOSSIN AND WAYNE H. SCHUBERT

Importance of Asymmetric Vorticity Dynamics

納莉颱風眼附近的中尺度渦旋

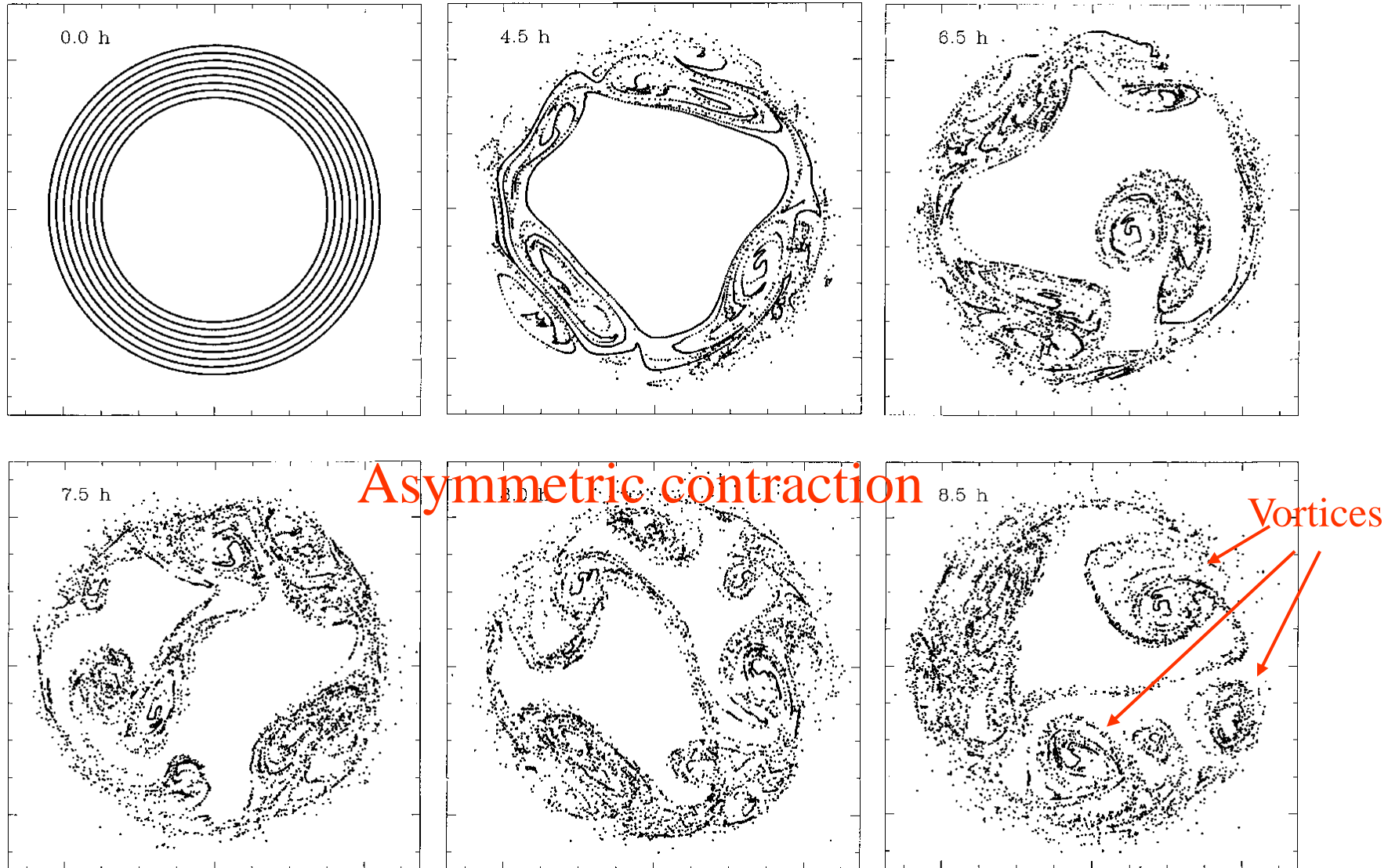


侵台納莉颱風登陸前
颱風眼附近觀測到3個
中尺度渦旋



Kossin and Schubert (2002) Mixing due to Barotropic Instability

20X20 km

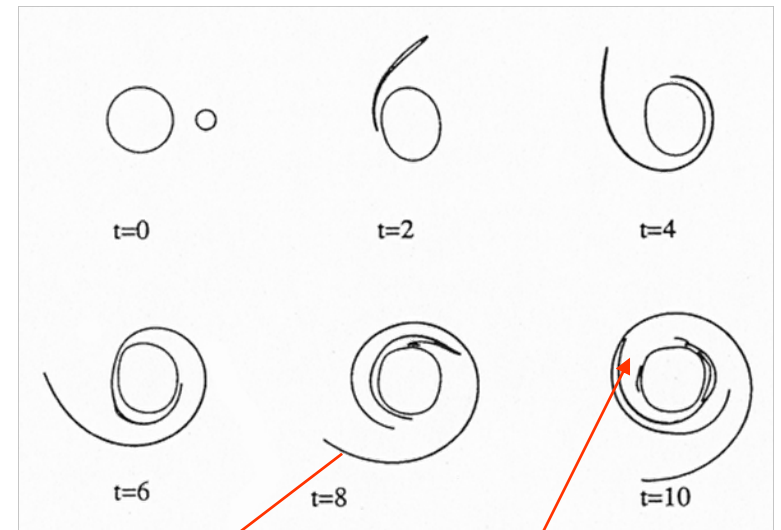
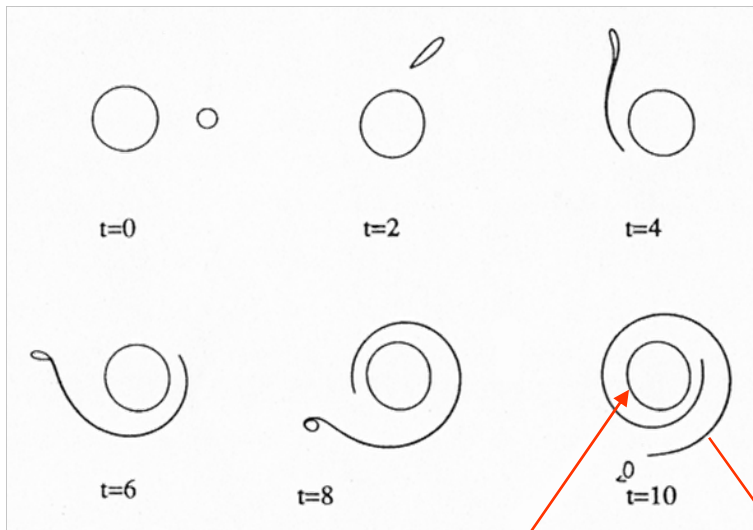


Advective rearrangement is **different from** the down-gradient diffusion

Straining out regime

partial straining - out (PSO)

complete straining - out (CSO)



Clear gap

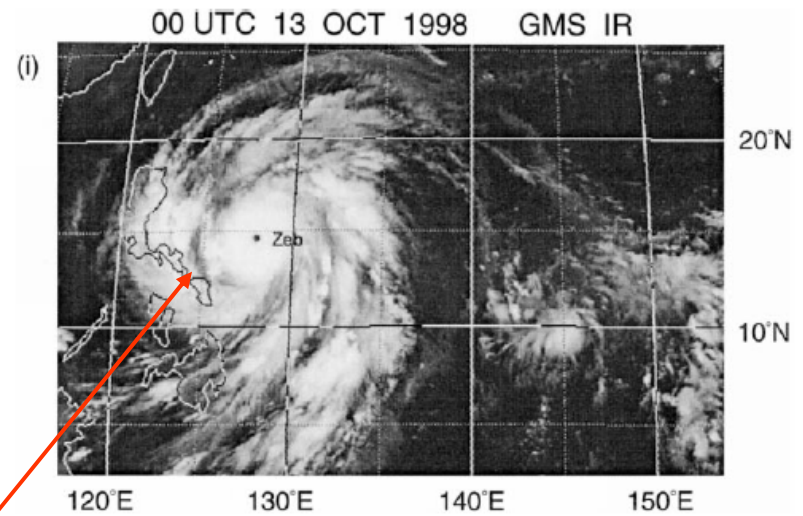
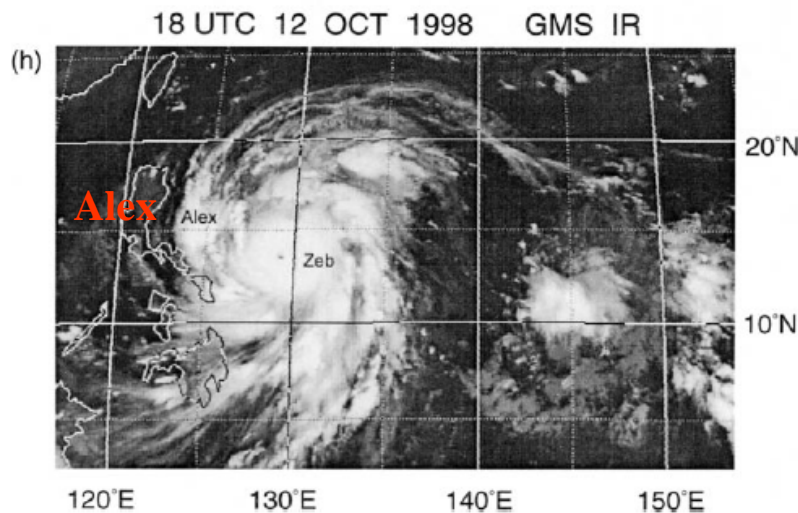
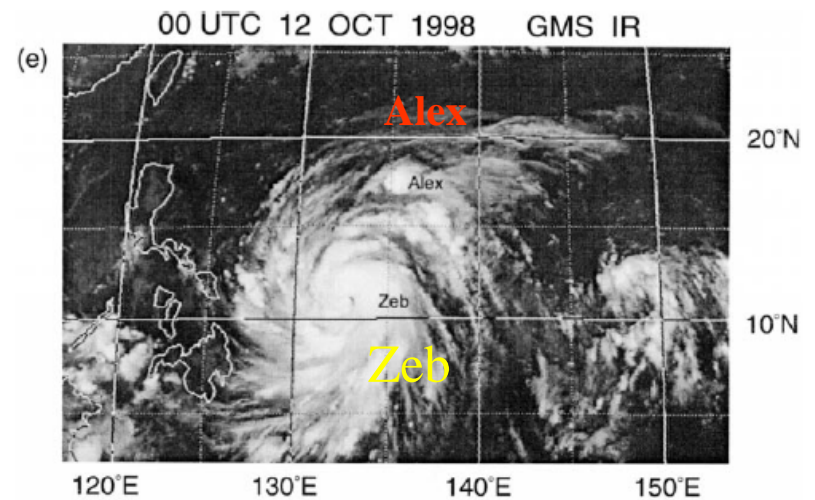
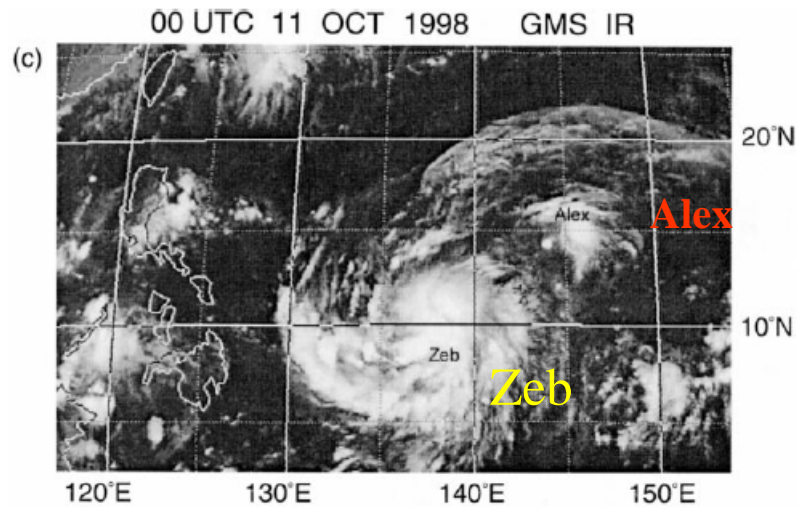
Clear gap

Adverse shear effect

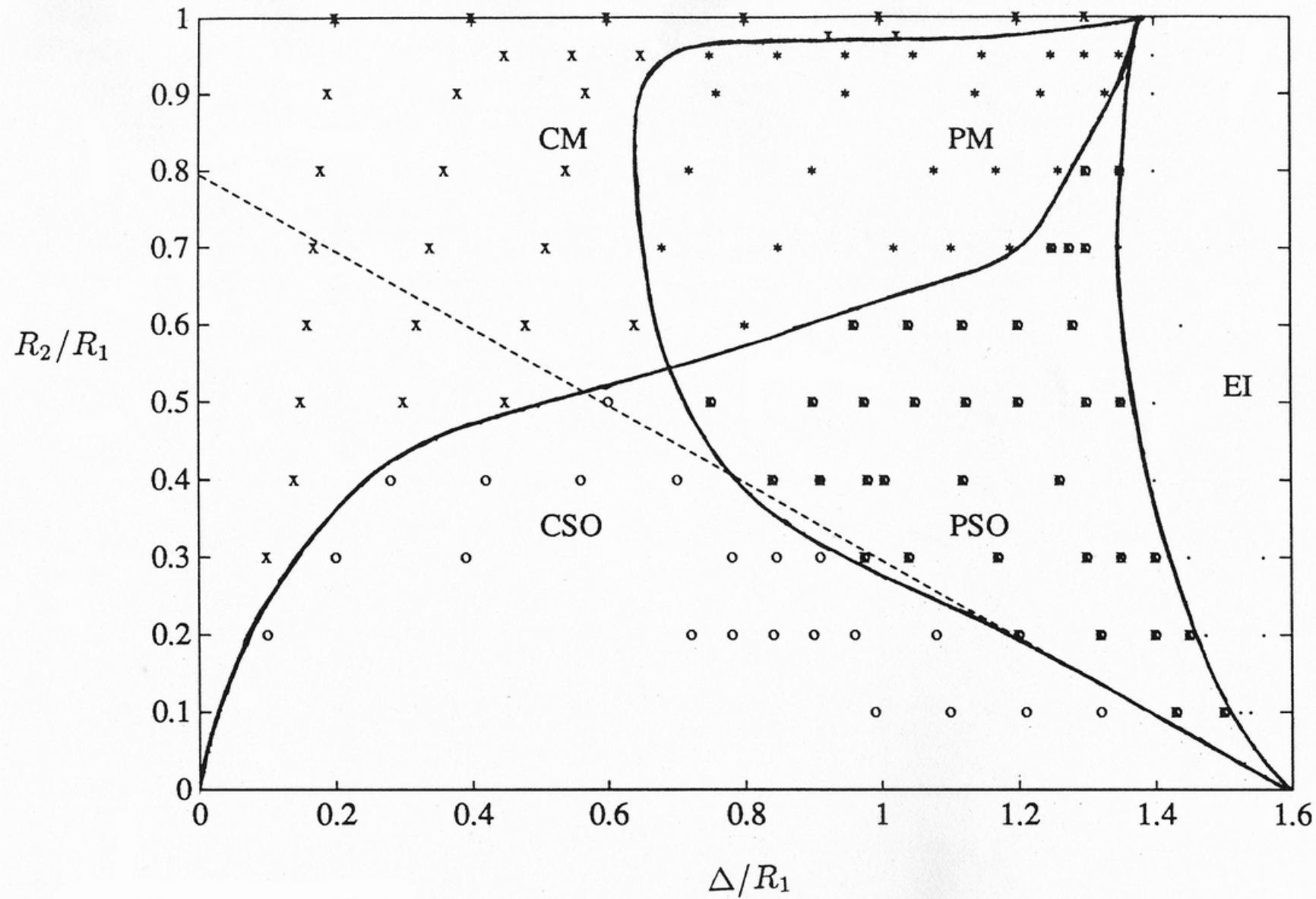
The bands are too thin to be called concentric eyewalls.

Kuo et al. (2000) MWR

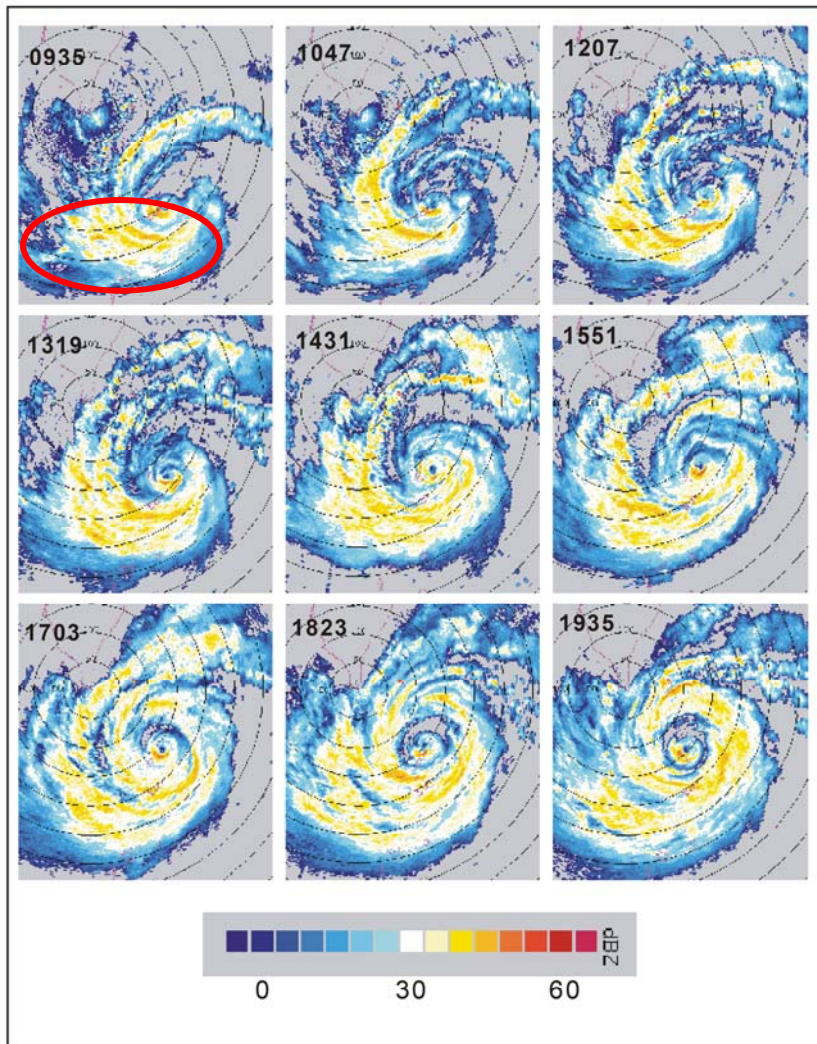
Typhoons Zeb and Alex ; straining – out regime



Clear gap between the Zeb and the remains of Alex



Dritschel and Waugh 1992



Concentric eyewalls formation in Typhoon Lekima (2001) near Taiwan

A huge area of convection outside the core vortex that wraps around the inner eyewall to form the concentric eyewalls in about 12 hours.

Figure 1a: Reflectivity at 0.5 elevation angle for Typhoon Lekima (2001) from the Central Weather Bureau WSR-88D (10 cm) radar at Kung-Ting for the period 0935 to 1935 September 25. The sequence of the images is from left to right and from top to bottom. The time interval between each image is approximately 75 min. The local time of observation is indicated on top of each image. The radial increment of the circles centered at radar station is 50 km. The nine images illustrate the formation of a

concentric eyewalls.

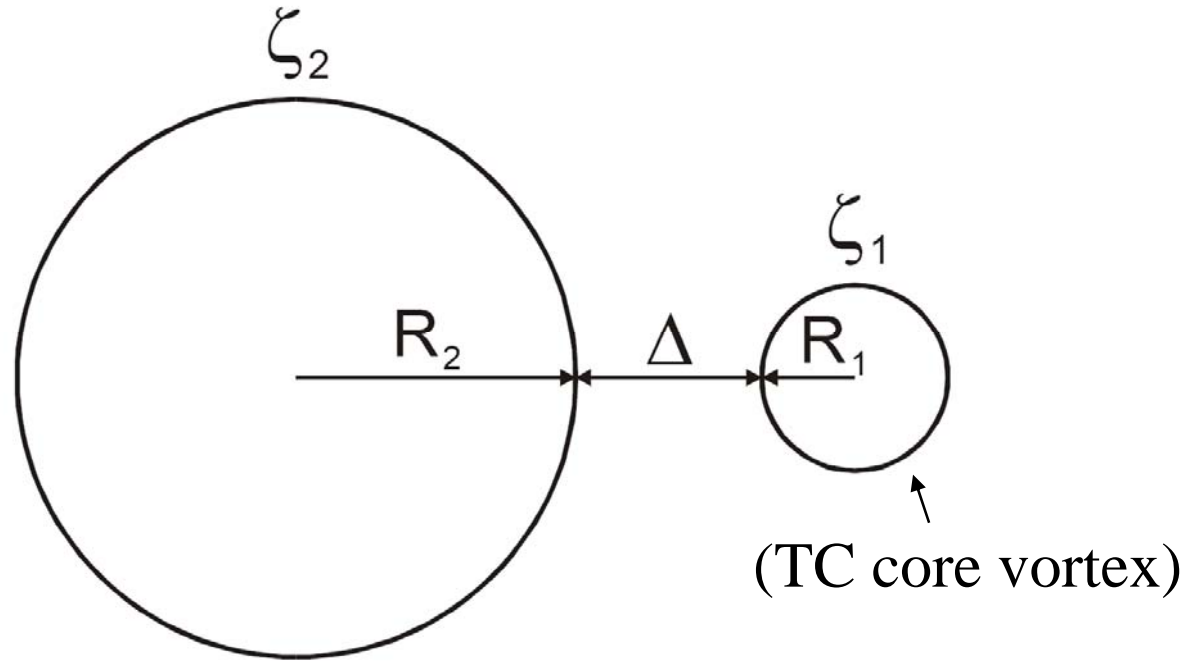
Kuo et al. (2004)

Binary vortex interaction

$$r = \frac{R_1}{R_2}$$

$$\Delta / R_1$$

$$\gamma = \frac{\zeta_1}{\zeta_2}$$



An extension of Dritschel and Waugh's (1992) work.

In addition to the radii ratio and the normalized distance between the two vortices, the **vorticity ratio** is added as a third external parameters.

Interaction of a small and strong vortex (representing TC core) with a large and weak vortex (representing vorticity induced by convection outside the TC core)

$$\gamma = \frac{\zeta_1}{\zeta_2}$$

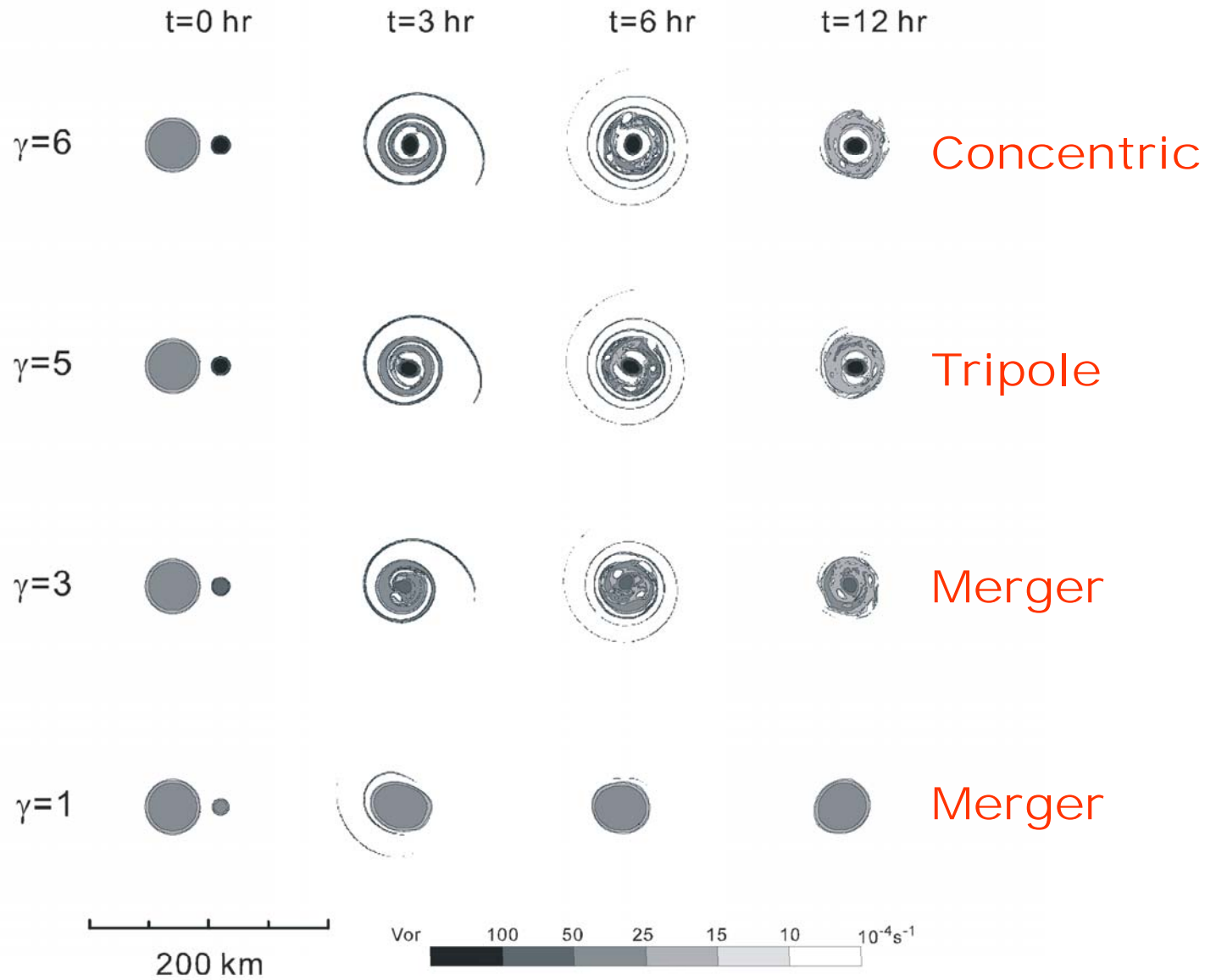


Figure 3: The sensitivity of the vorticity field in the binary vortex experiments with respect to the vorticity strength ratio (γ) at hour 0, 3, 6 and 12 with the dimensionless gap $\Delta/R_1=1$, and the vortex radius ratio $r=1/3$.

Vortex radius ratio

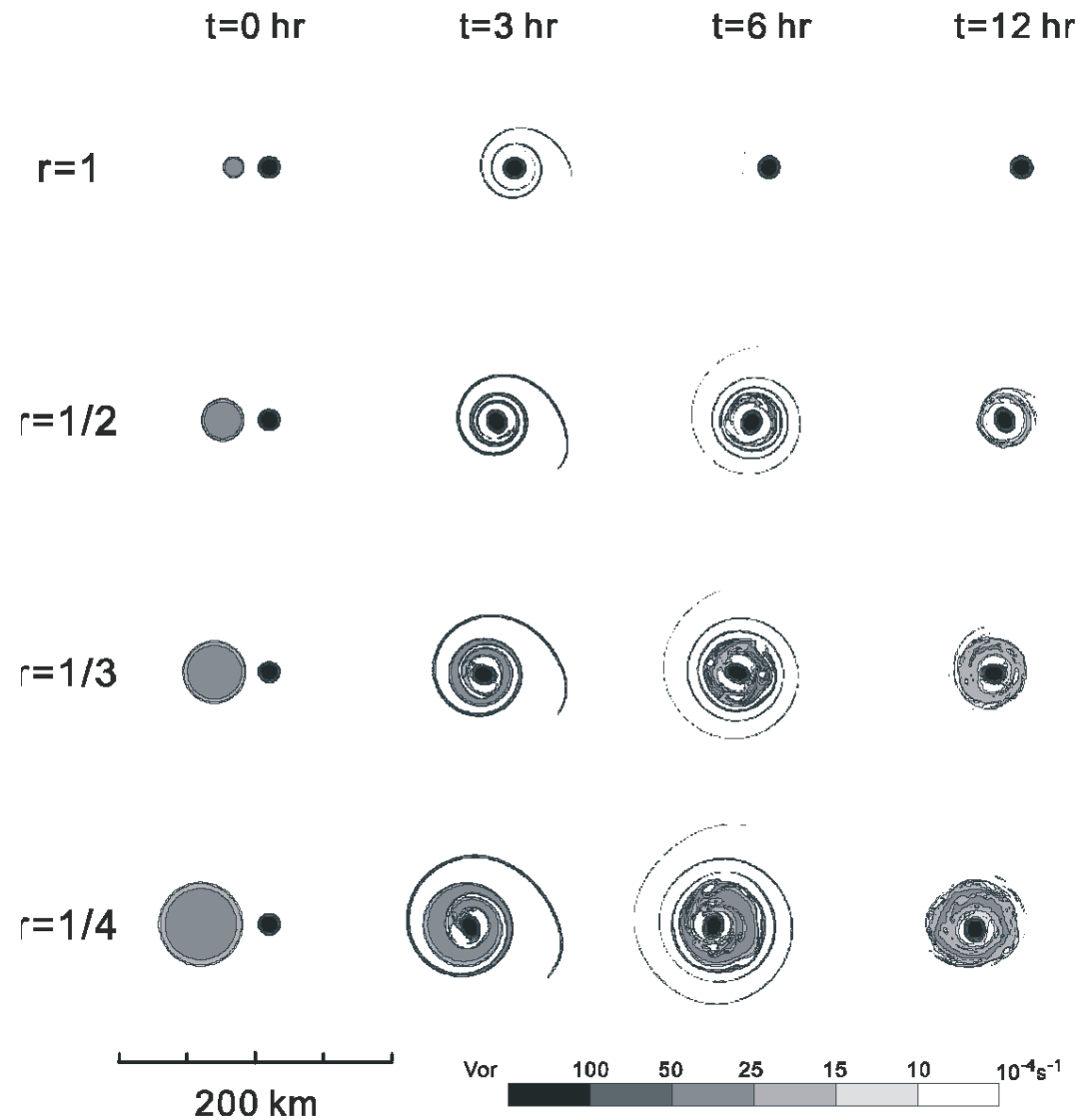


Figure 5: The sensitivity of the vorticity field in the binary vortex experiments with respect to the vortex radius ratio r at hour 0, 3, 6 and 12 with the vorticity strength ratio $\gamma=5$, and the dimensionless gap $\Delta/R_1=1$.

$$\gamma = 10$$

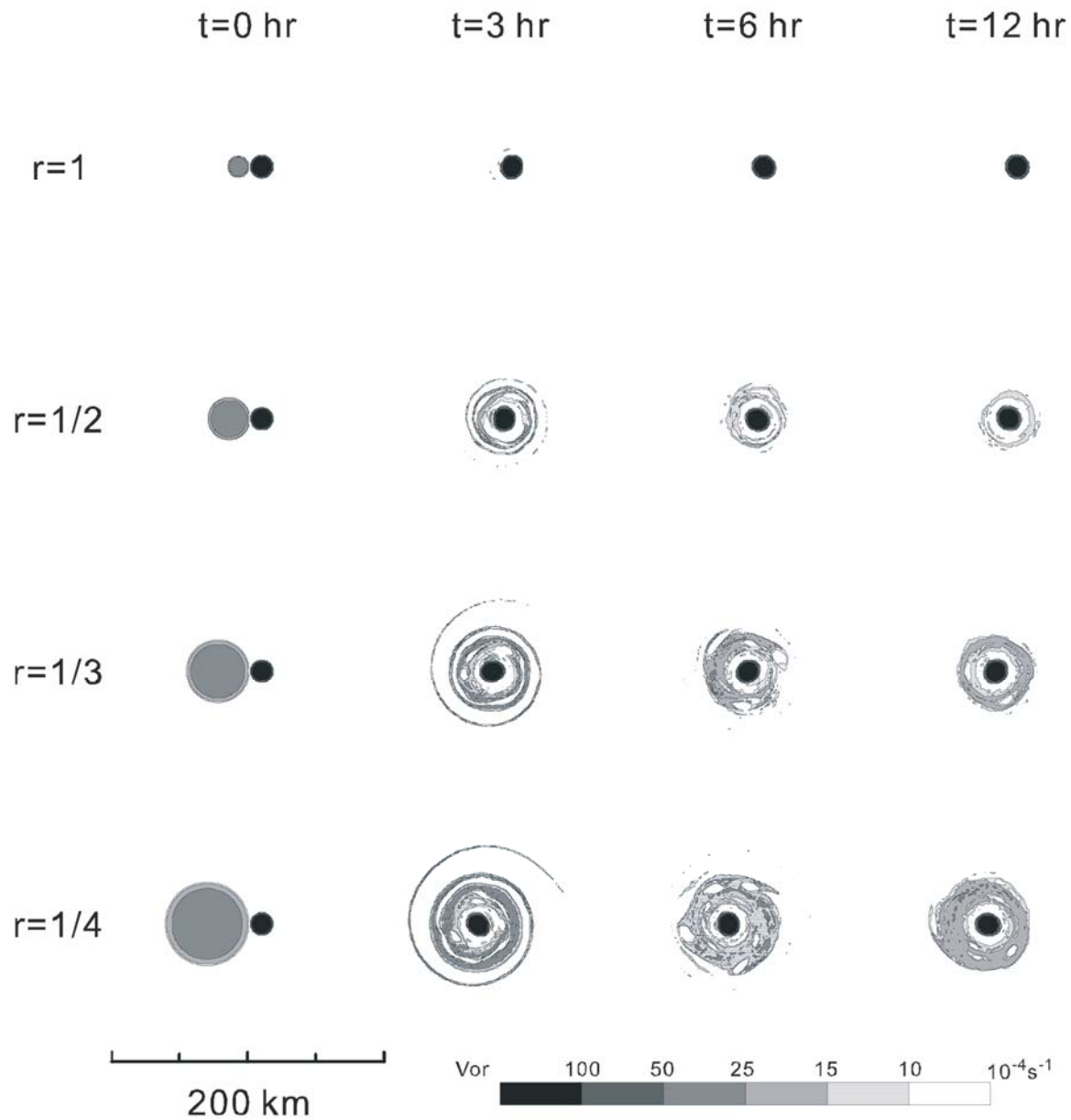


Figure 6: Similar to Figure 5 except that the dimensionless gap $\Delta/R_1=0$ and the vorticity strength ratio $\gamma=10$.

Same strain outside
the RMW but
different core vortex
strength

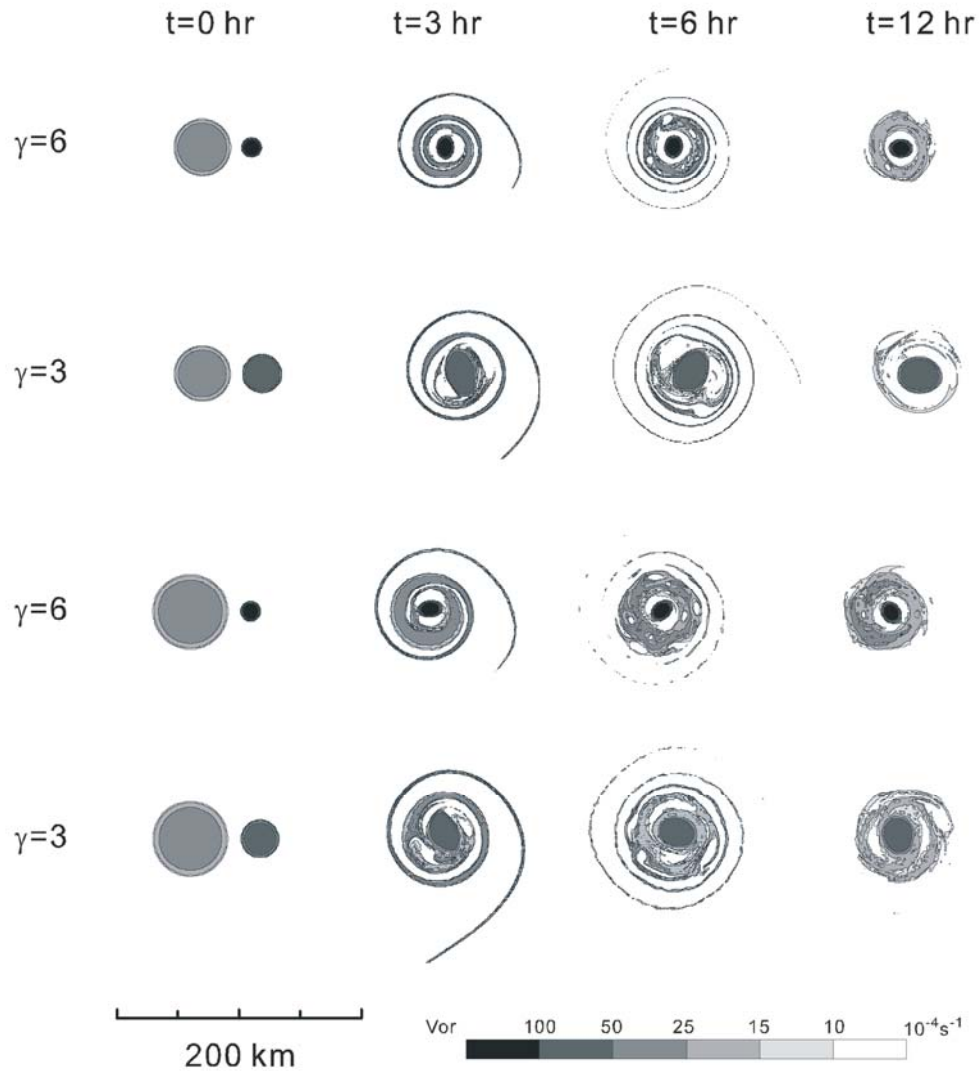


Figure 9: The sensitivity of the vorticity field in the binary vortex experiments with the core vortices process the same maximum wind but different radius of vorticity field. Two core vortices considered have the vorticity and radius of $(1.8 \times 10^{-2} \text{s}^{-1}, 10 \text{ km})$ and $(0.9 \times 10^{-2} \text{s}^{-1}, 20 \text{ km})$ respectively. The dimensionless gap is 1 in the experiments. The outer vortices considered have the radius of 30 km and 40 km respectively.

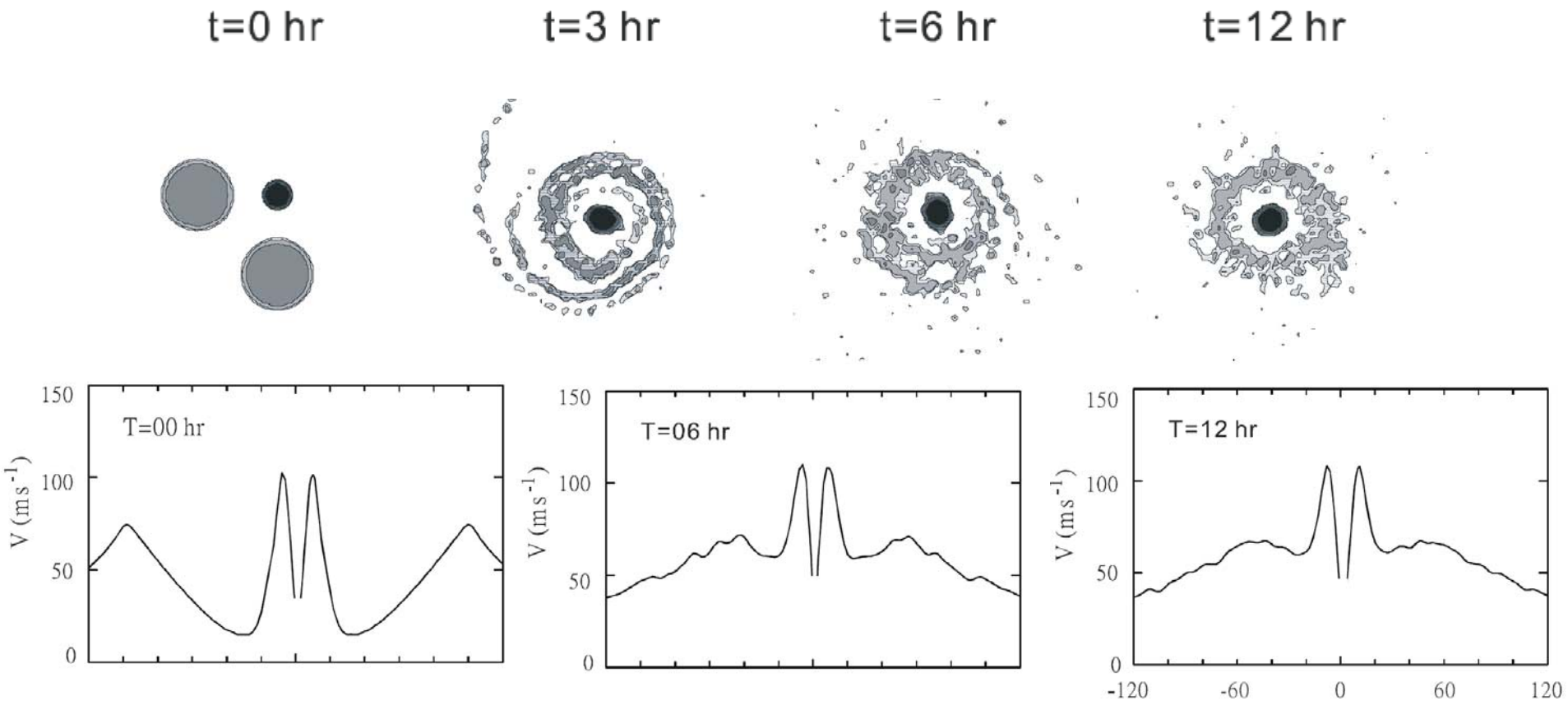


Figure 13: The tangential wind speed for radial arms toward the west (left portion) and the south (right portion) that emanate from the vortex center at various times for the experiment in the second row of Figure 11.

The contraction of the secondary wind maximum by nonlinear advective dynamics.

Kuo et al. (2004)

- A very strong core vortex at least **six times stronger** than the neighboring vorticity
- A relative **larger** neighboring vorticity area
- A separation distance within **three to four times** the core vortex radius

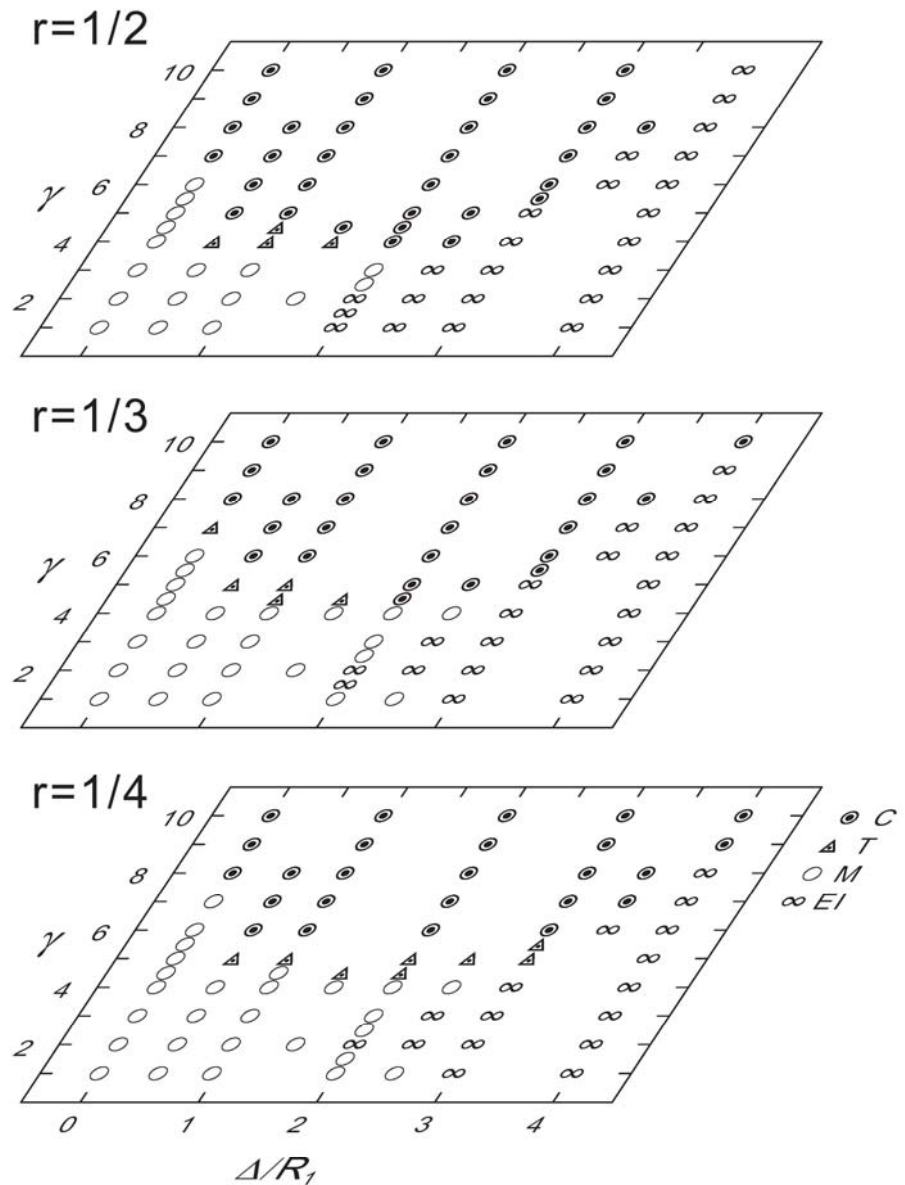
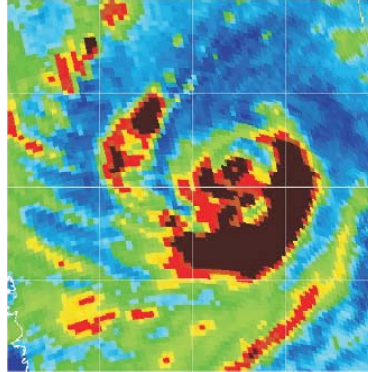


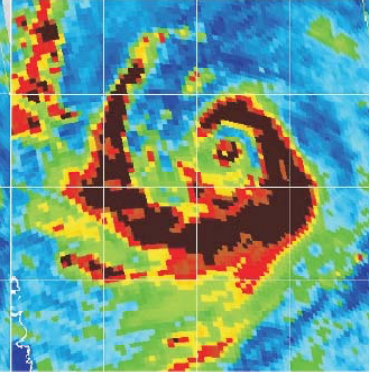
Fig 10. Summary of numerical experiments with the parameters of the vorticity strength ratio (γ), the dimensionless gap Δ/R_1 , and the vortex radius ratio r . We have classified the resulting structures into the C (concentric), T (tripole), M (complete or partial merger), and EI (elastic interaction) regimes.

2003-09W
IMBUDO

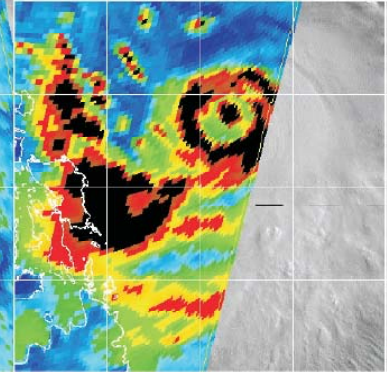
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0720-0942Z



072012Z-130kts
0720-1244Z

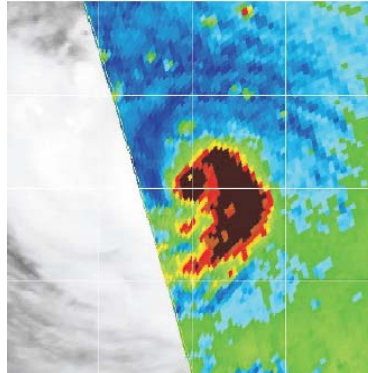


072100Z-130kts
0720-2219Z

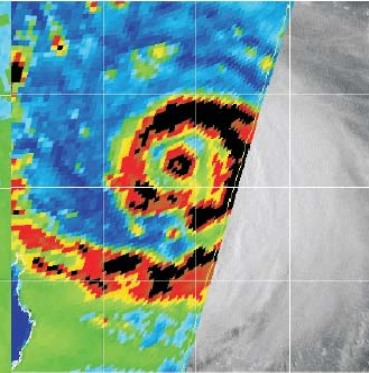


2003-14W
DUJUAN

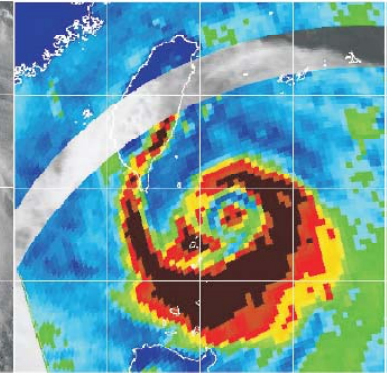
083112Z-95kts
0831-1214Z



090100Z-120kts
0831-2235Z

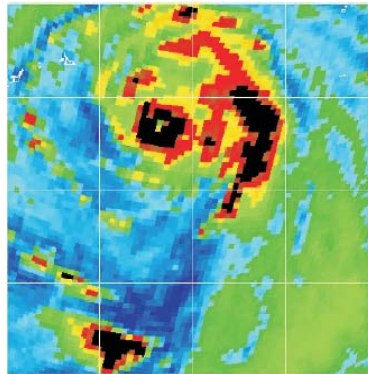


090112Z-125kts
0901-0949Z

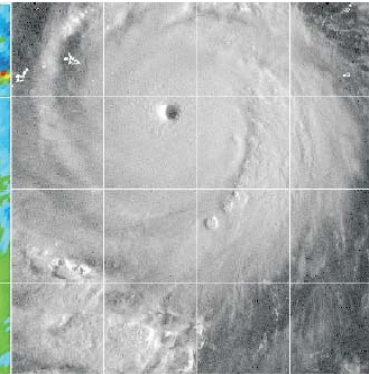


2003-15W
MAEMI

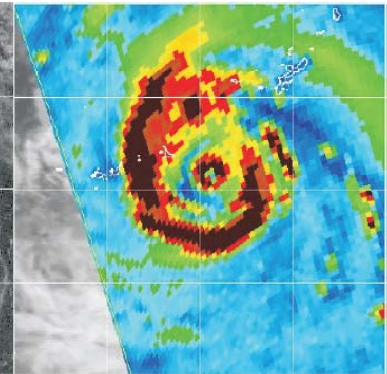
091000Z-150kts
0909-2209Z



091000Z-150kts
0909-2239Z

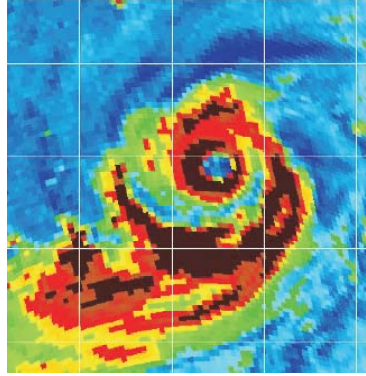


091012Z-150kts
0910-0925Z

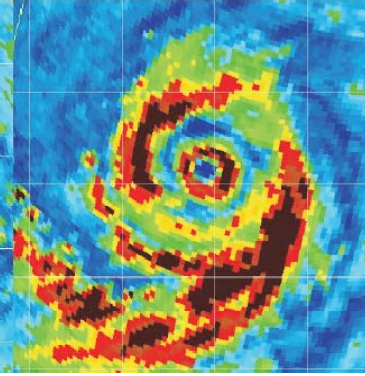


2004-09W
DIANMU

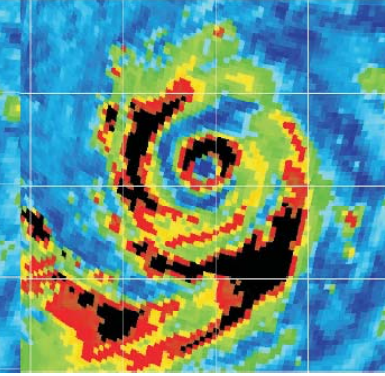
061800Z-120kts
0618-0049Z



061812Z-130kts
0618-1202Z

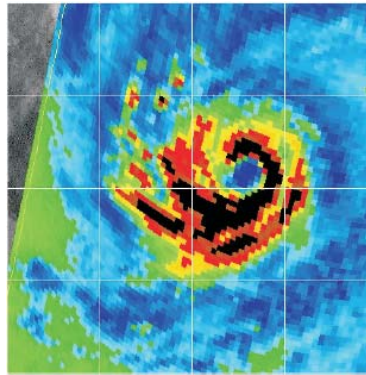


061900Z-120kts
0619-0033Z

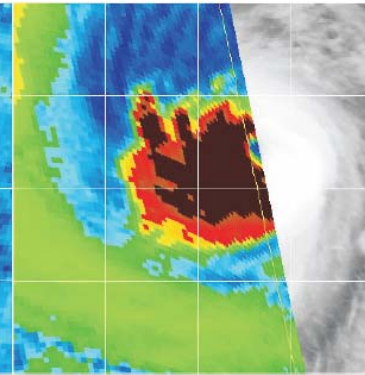


2003-20W
KETSANA

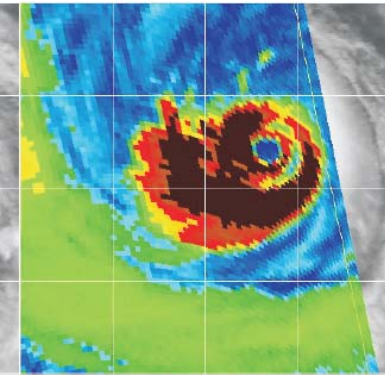
102100Z-95kts
1020-2302Z



102112Z-125kts
1021-0954Z

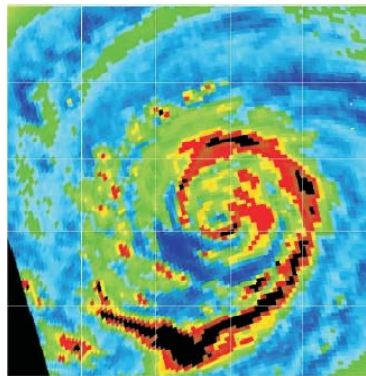


102112Z-125kts
1021-1246Z

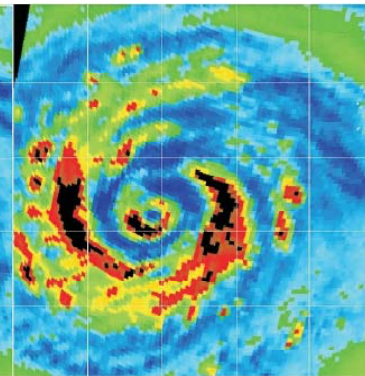


1997-14W
WINNIE

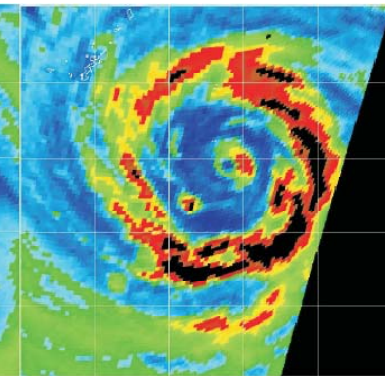
081412Z-110kts
0814-1031Z



081500Z-105kts
0814-2316Z

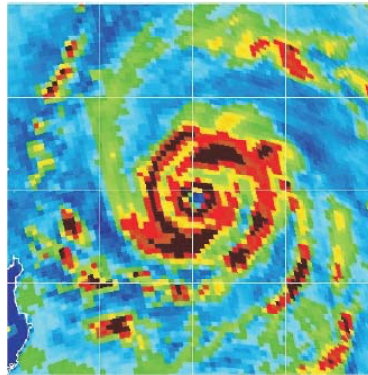


081600Z-85kts
0816-0154Z

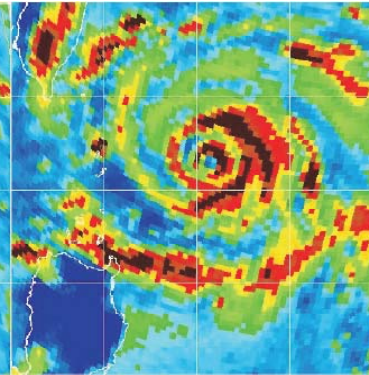


2000-18W
BILIS

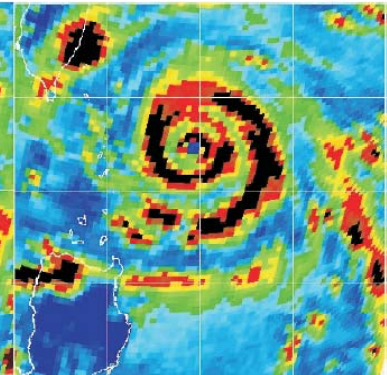
082112Z-135kts
0821-1214Z



082200Z-140kts
0821-2152Z

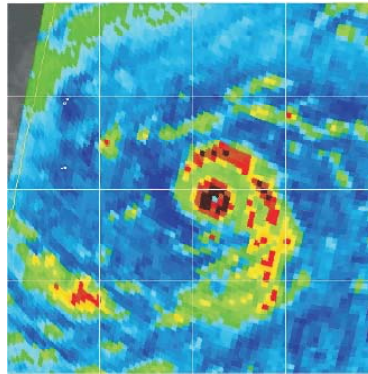


082200Z-140kts
0822-0047Z

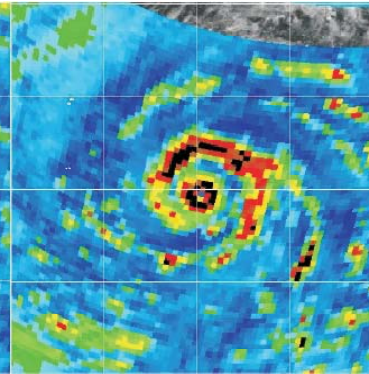


2000-22W
SAOMAI

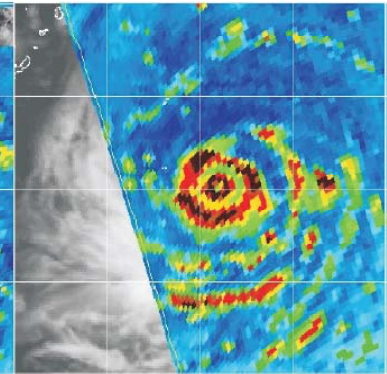
091000Z-120kts
0909-2109Z



090100Z-120kts
0831-2235Z

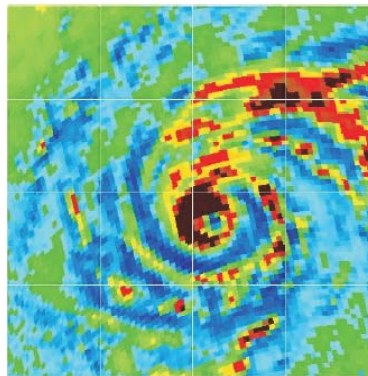


090112Z-125kts
0901-0949Z

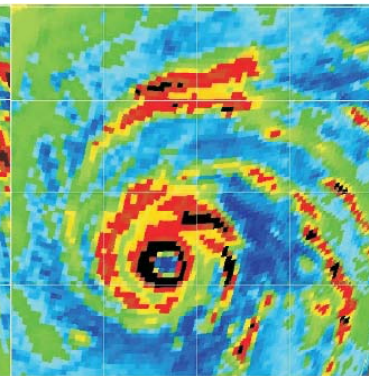


2001-26W
PODUL

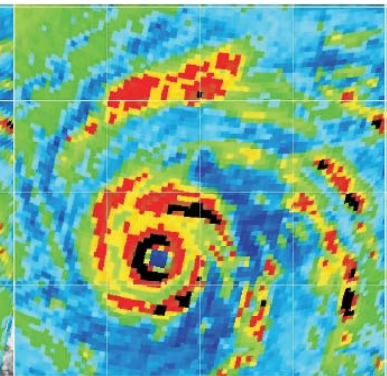
102112Z-140kts
1024-1049Z



102500Z-130kts
1024-2228Z

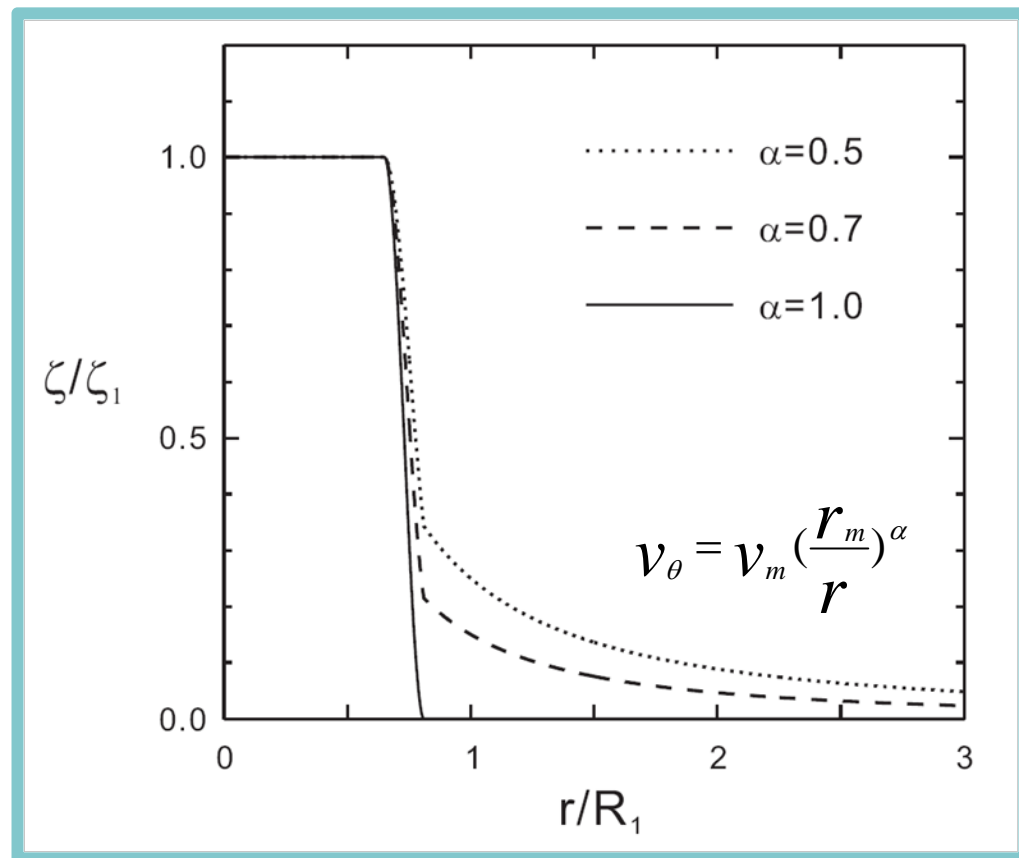


102500-130kts
1024-2323Z



Advection dynamics: Core structure important

TC structure is characterized by a relatively slow tangential wind decrease outside the radius of maximum wind (RMW) and a corresponding skirt of significant cyclonic relative vorticity. (Mallen et al. 2005)



Radial profiles of the vorticity for
the strong core vortex

$$\Delta/R_1 = 3$$

t=0hr

t=12hr

t=12hr

t=12hr

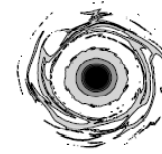
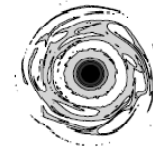
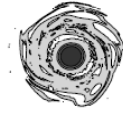
$\gamma=4$

$\gamma=7$

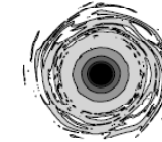
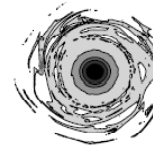
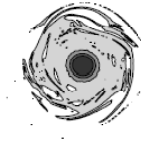
$\gamma=10$

r = 1/3

$\alpha=0.7$

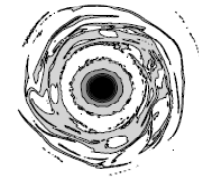
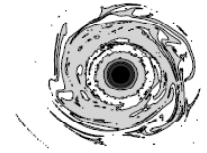
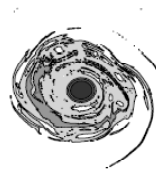


$\alpha=0.5$

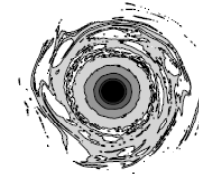
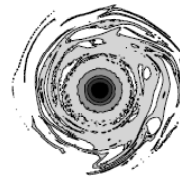
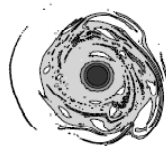
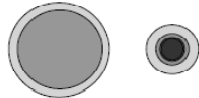


r = 1/4

$\alpha=0.7$



$\alpha=0.5$



200 km

Vor 200 100 50 25 10 10^{-4}s^{-1}

$$\Delta/R_1 = 5$$

t=0hr

t=12hr

t=12hr

t=12hr

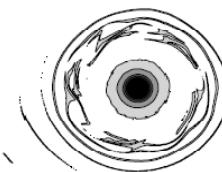
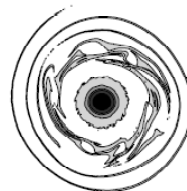
$\gamma=4$

$\gamma=7$

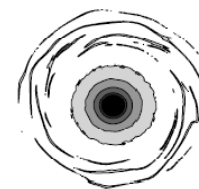
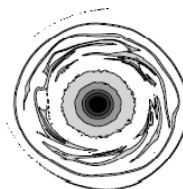
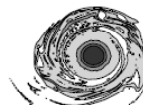
$\gamma=10$

$r = 1/3$

$\alpha=0.7$

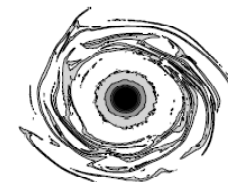
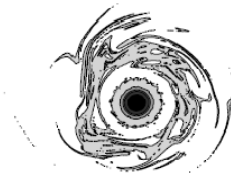
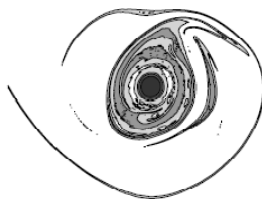


$\alpha=0.5$

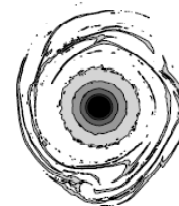
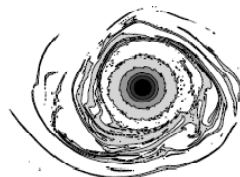
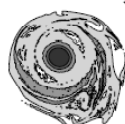


$r = 1/4$

$\alpha=0.7$



$\alpha=0.5$

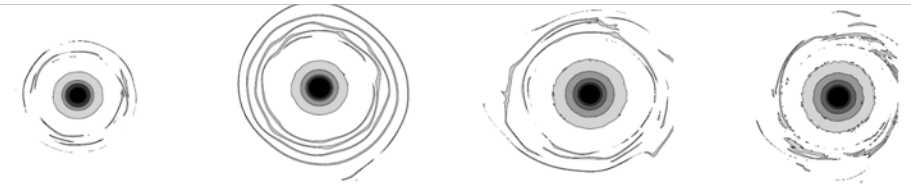


200 km

Vor 200 100 50 25 10 10^{-4}s^{-1}

Examples of the vorticity field at hour 12, showing different classifications of binary vortex interactions involving a skirted core vortex.

Straining out



Concentric



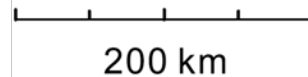
Tripole



Merger

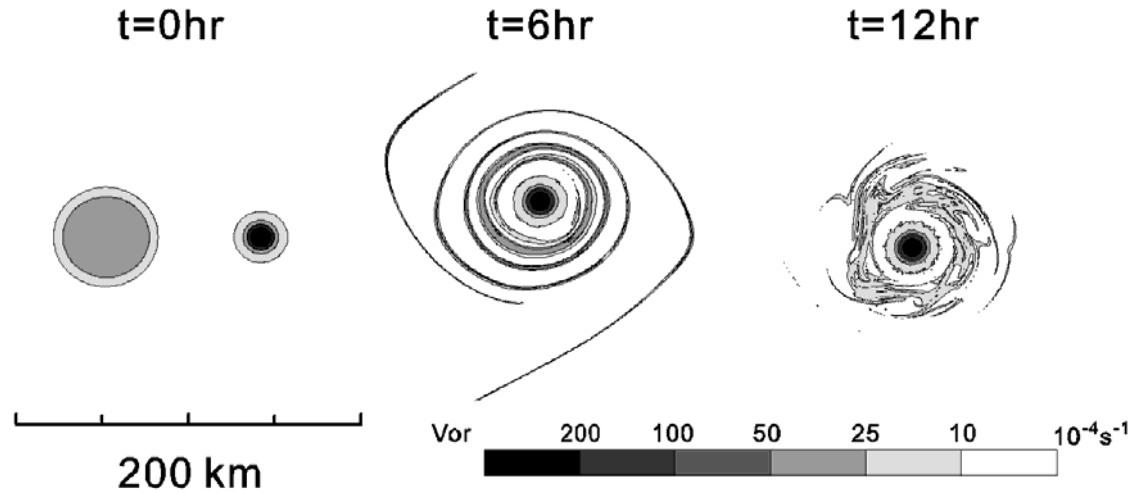


Elastic Interaction

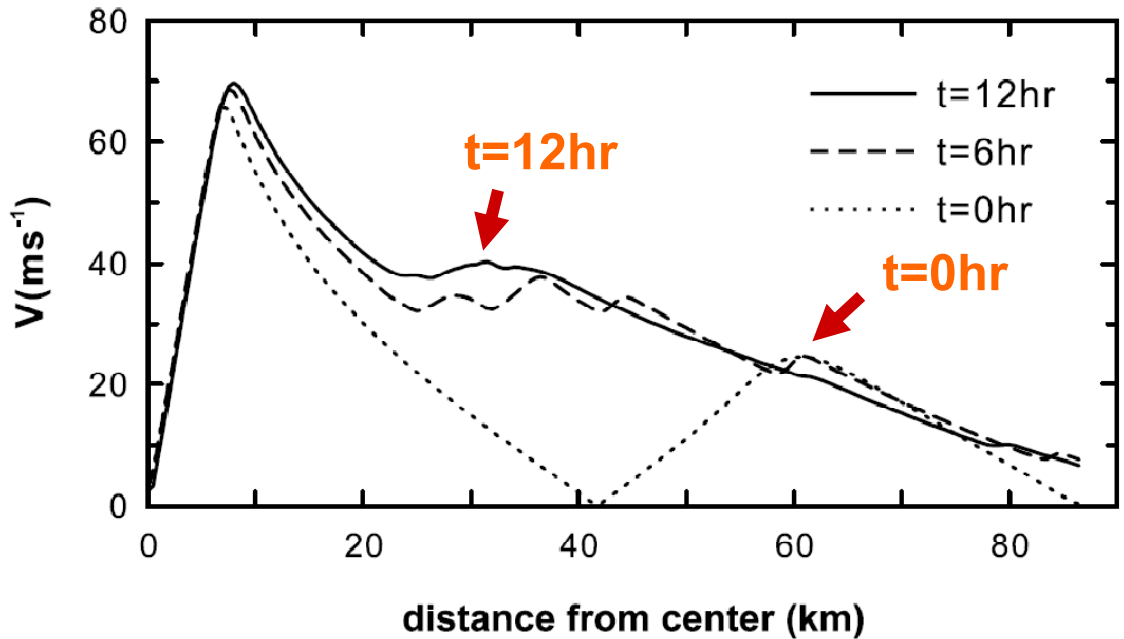


$$\alpha = 0.7 \quad \gamma = 7 \quad r = 1/4 \quad \Delta/R_1 = 5$$

The formation of concentric eyewalls



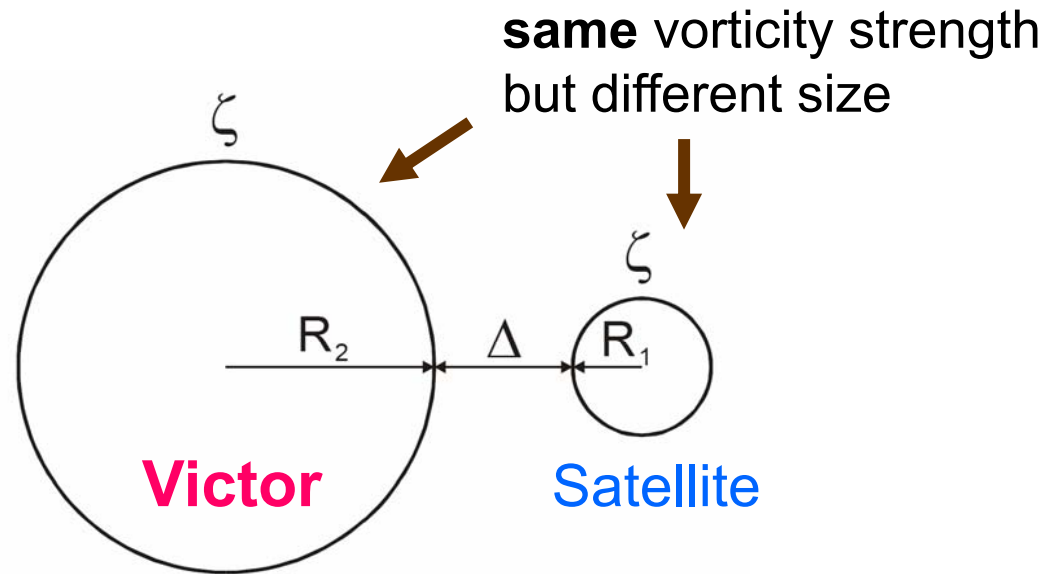
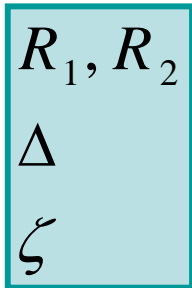
The contraction of the secondary wind maximum by nonlinear advection dynamics.



Binary vortex interaction

Dritschel and Waugh (1992)

【Variables】



【Parameters】

- Vortex radius ratio (r) = $\frac{R_1}{R_2}$
- Dimensionless gap ($\frac{\Delta}{R_1}$)

【Conclusion】

- Elastic Interaction (EI)
- Partial straining-out (PSO)
- Complete straining-out (CSO)
- Partial merger (PM)
- Complete merger (CM)

Binary vortex interaction

Kuo et al. (2004)

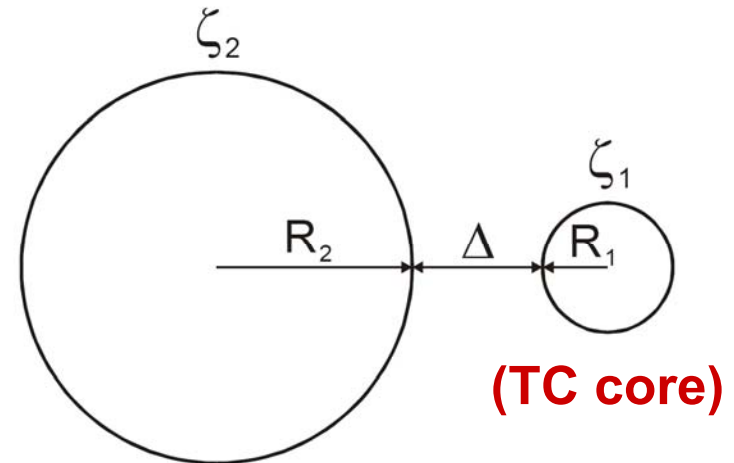
【Variables】

$$R_1, R_2; \Delta; \zeta_1, \zeta_2$$

【Parameters】

- Vortex radius ratio (r) = $\frac{R_1}{R_2}$
- Dimensionless gap ($\frac{\Delta}{R_1}$)

$$\bullet \text{ Vortex strength ratio } (\gamma) = \frac{\zeta_1}{\zeta_2}$$



- An extension of Dritschel and Waugh's (1992) work.
- In addition to **the radii ratio** and **the normalized distance** between the two vortices, **the vorticity ratio** is added as a third external parameters.

**Dritschel and
Waugh (1992)**



$$R_1, R_2, \Delta, \zeta$$
$$r = \frac{R_1}{R_2}, \Delta/R_1$$

**Kuo et al.
(2004)**



$$R_1, R_2, \Delta, \zeta_1, \zeta_2$$
$$r = \frac{R_1}{R_2}, \Delta/R_1, \gamma = \frac{\zeta_1}{\zeta_2}$$

**Kuo et al.
(2006)**

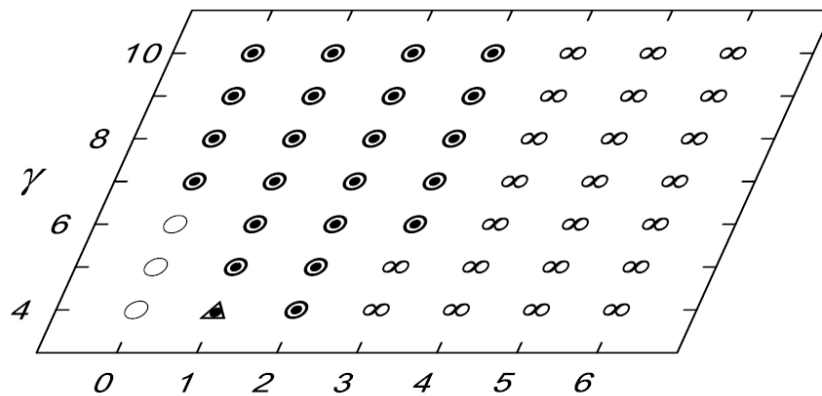


$$R_1, R_2, \Delta, \zeta_1, \zeta_2, \alpha$$
$$r = \frac{R_1}{R_2}, \Delta/R_1, \gamma = \frac{\zeta_1}{\zeta_2}, \alpha$$

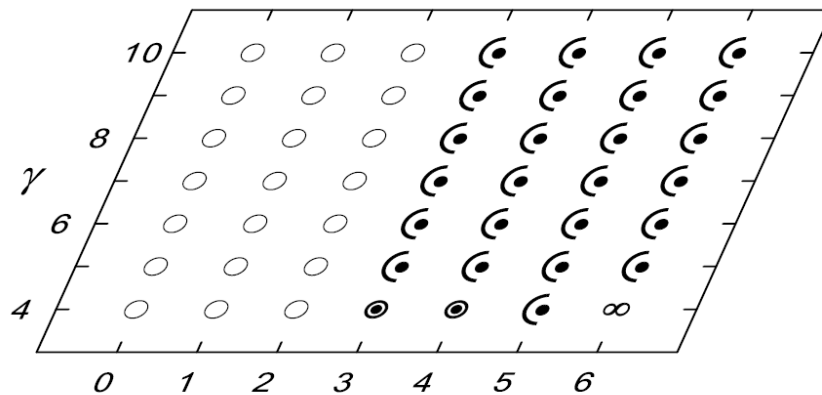
(a)

$r=1/2$

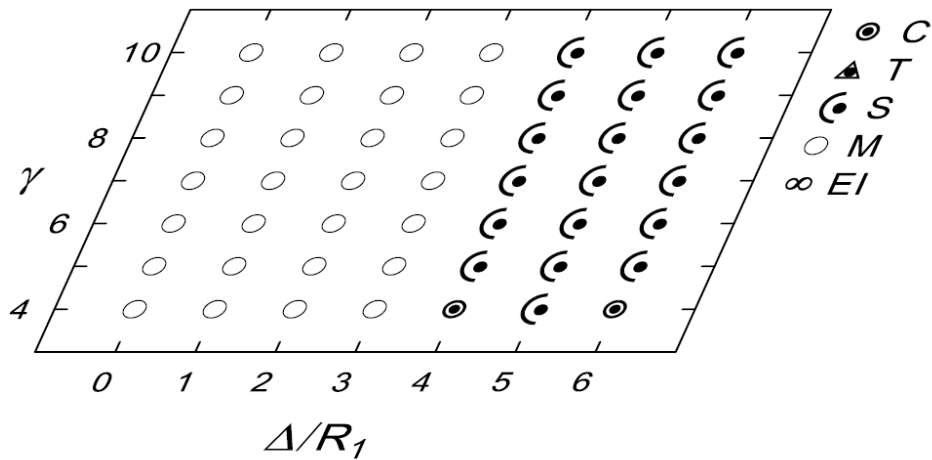
$\alpha=1.0$



$\alpha=0.7$



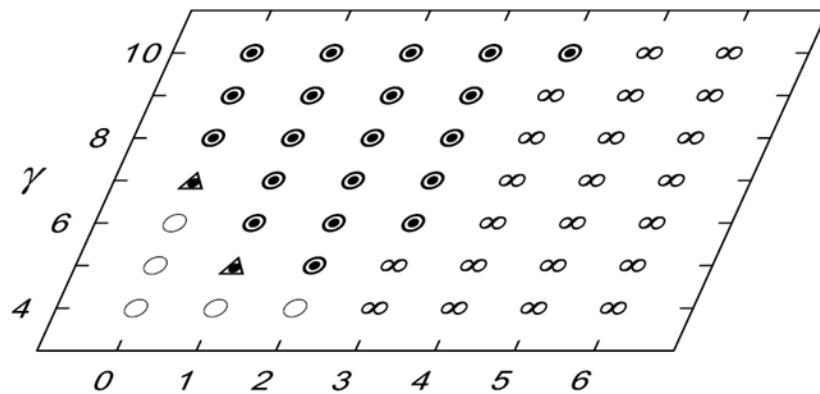
$\alpha=0.5$



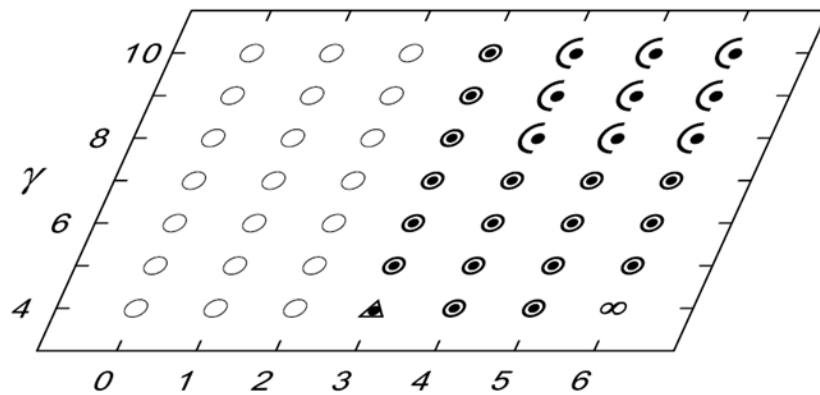
(b)

$r=1/3$

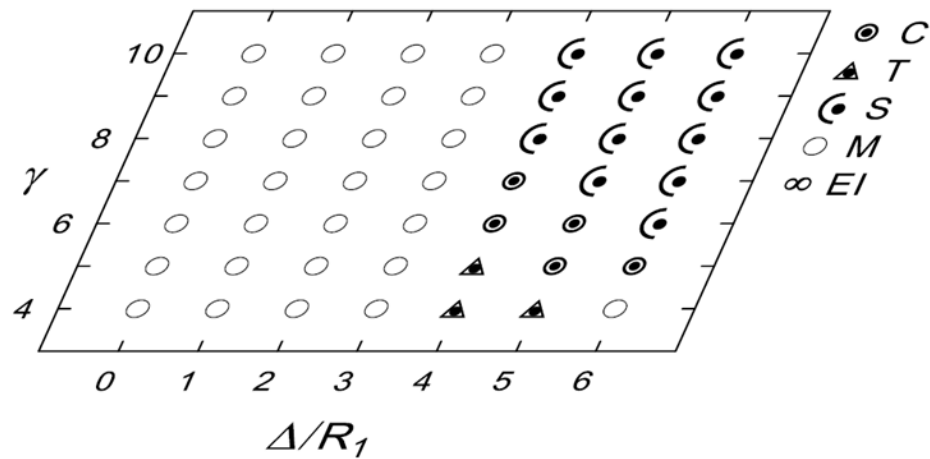
$\alpha=1.0$



$\alpha=0.7$



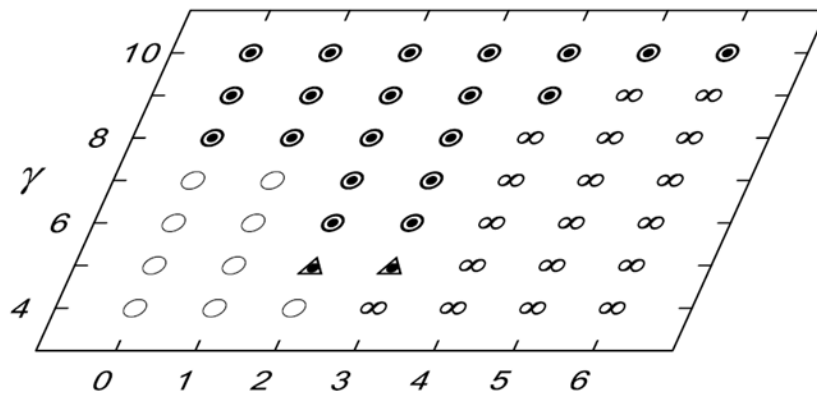
$\alpha=0.5$



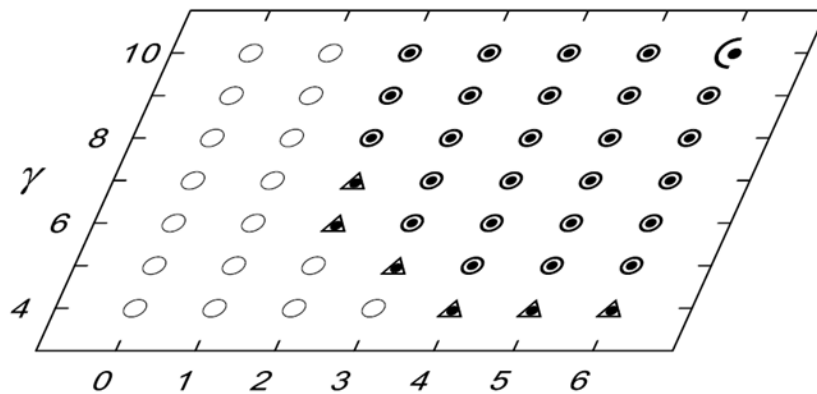
(c)

$r=1/4$

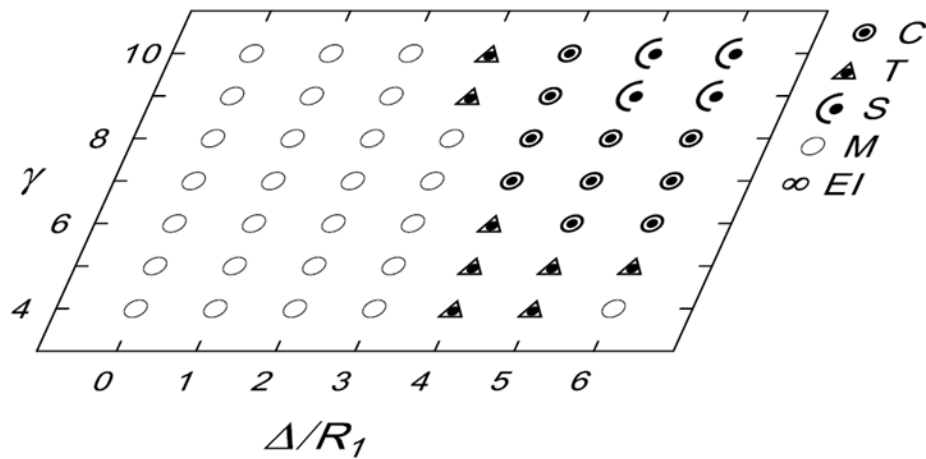
$\alpha=1.0$



$\alpha=0.7$



$\alpha=0.5$



$r^* = 3$

t = 0 hr

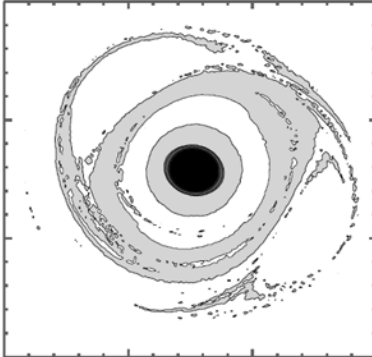
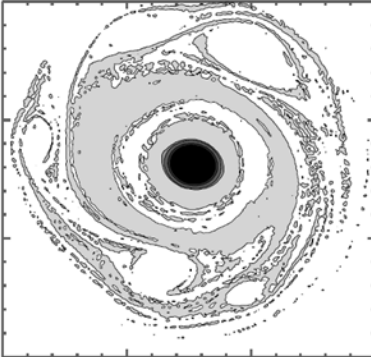
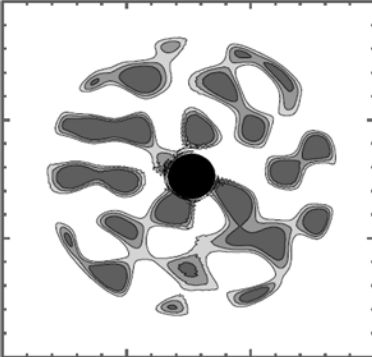
t = 24 hr

t = 24 hr

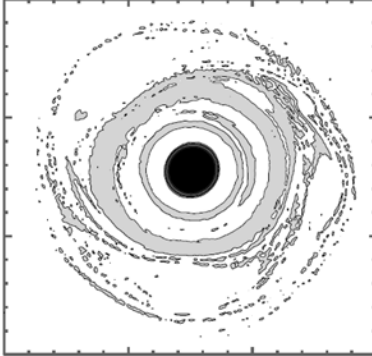
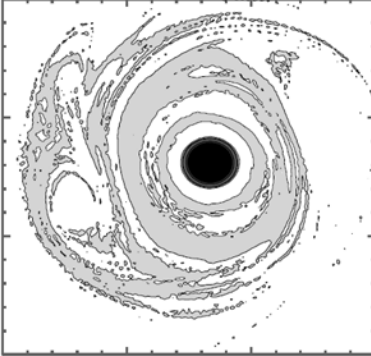
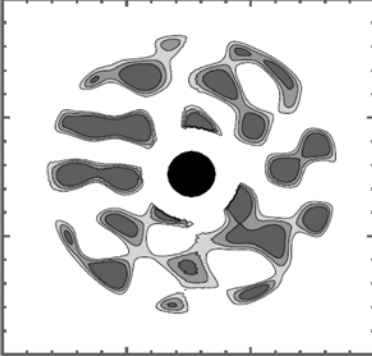
$\gamma = 6$

$\gamma = 10$

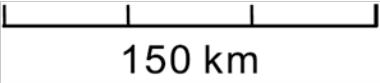
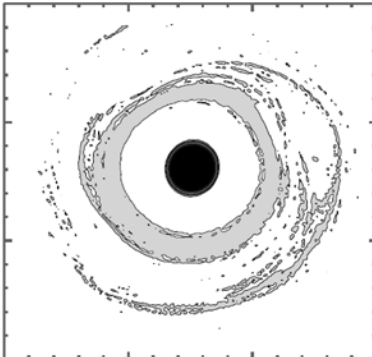
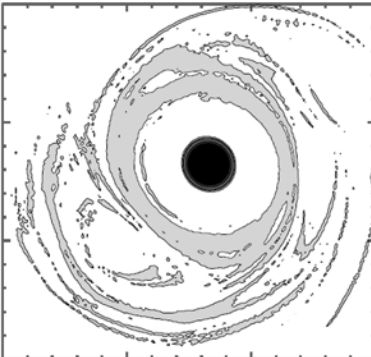
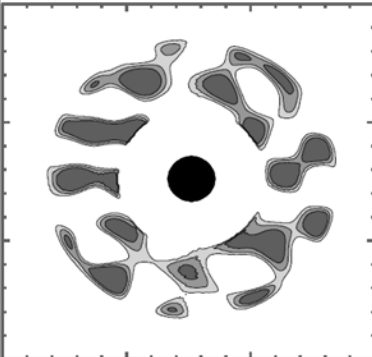
no-moat



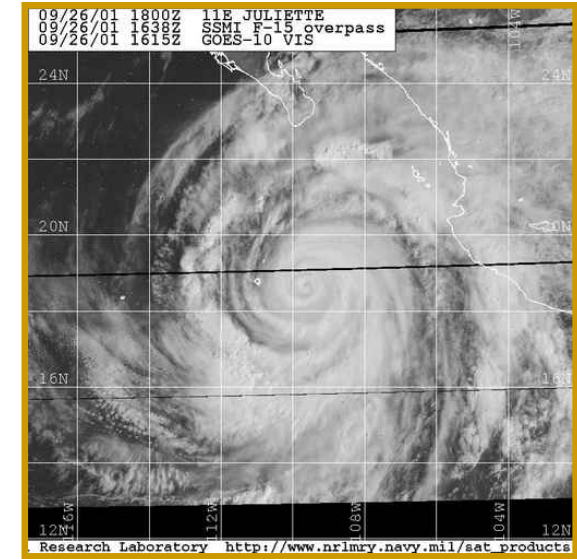
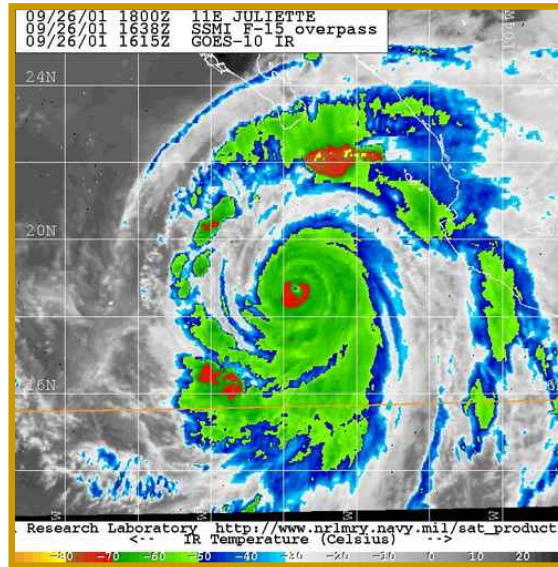
moat
= 10 km



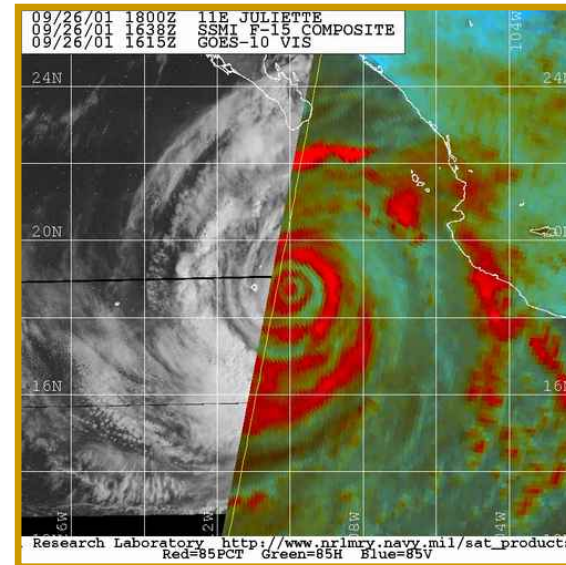
moat
= 20 km



Hurricane Juliette (2001)



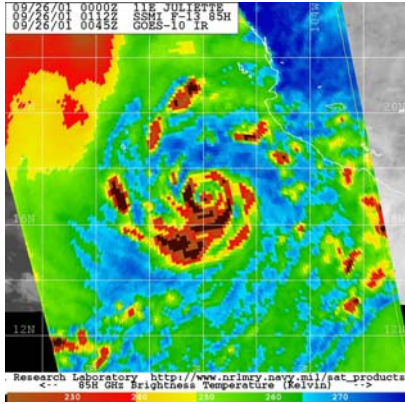
“the crew found three concentric eyewalls, defined by three peaks in tangential wind in each of the radial legs and by three complete rings of enhanced radar reflectivity with radii of 11, 56, and 90 km.”



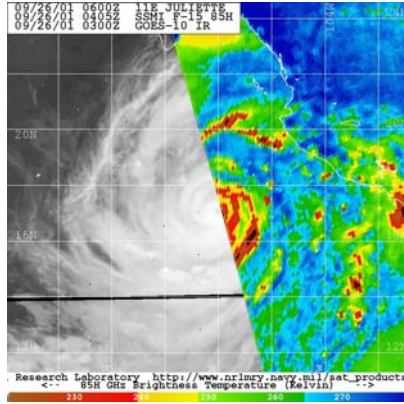
McNoldy (2004)

Observation

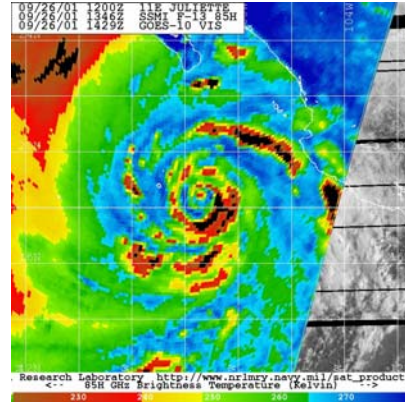
0112Z



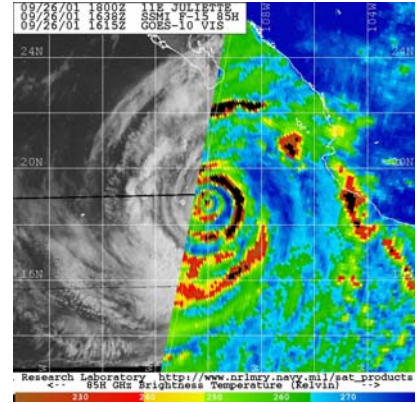
0405Z



1346Z

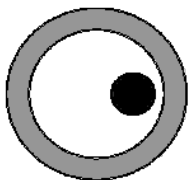


1638Z

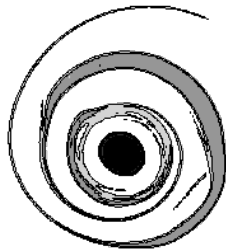
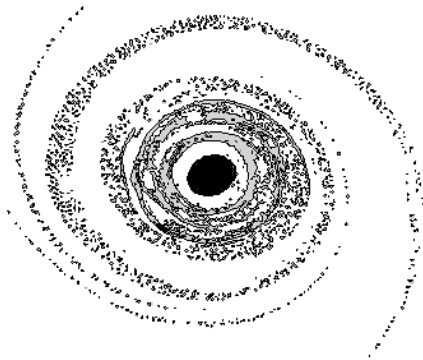


Simulation

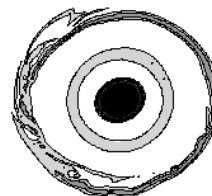
t=0hr



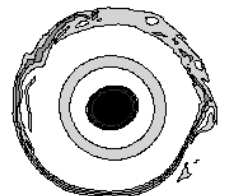
t=3hr



t=12hr



t=15hr



$r^* = 6$

$t = 0$ hr

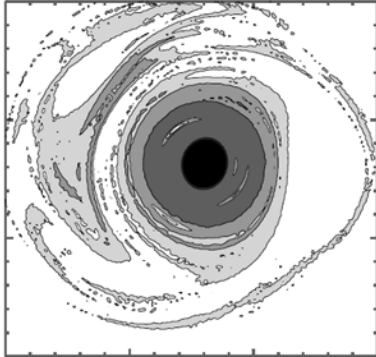
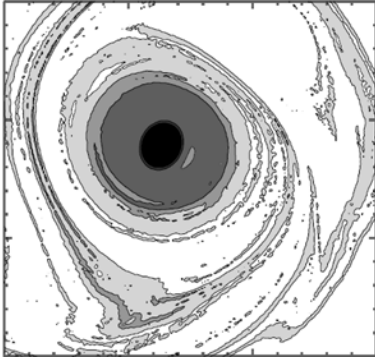
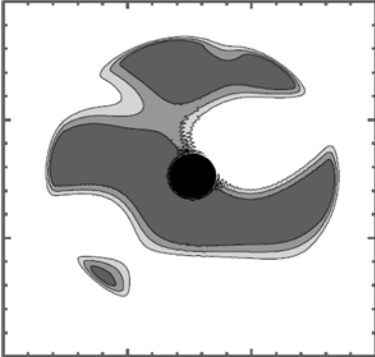
$t = 24$ hr

$t = 24$ hr

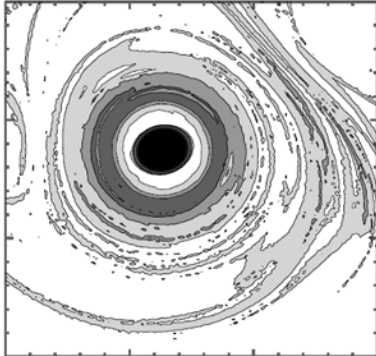
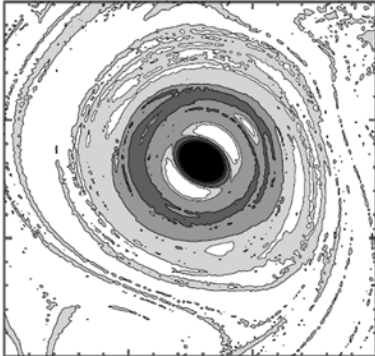
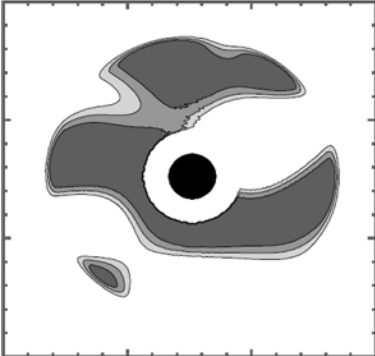
$\gamma = 6$

$\gamma = 10$

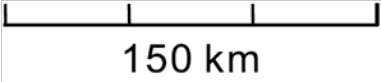
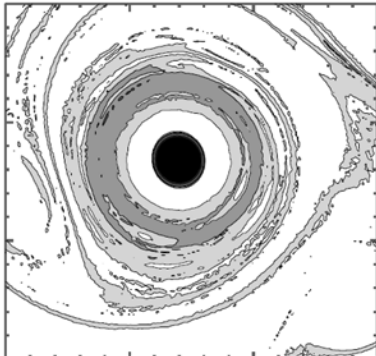
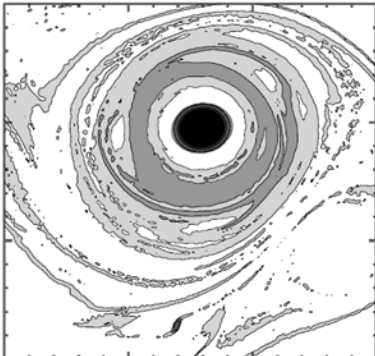
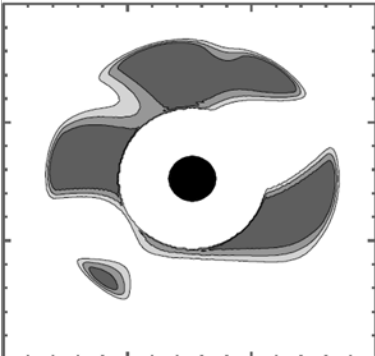
no-moat



moat
= 10 km



moat
= 20 km



$r^* = 4$

$t = 0$ hr

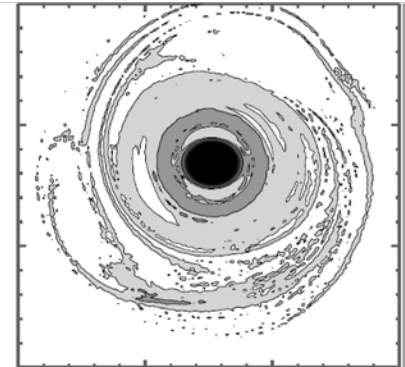
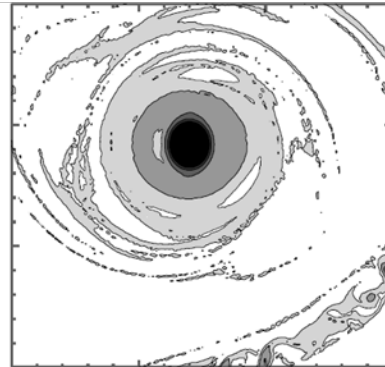
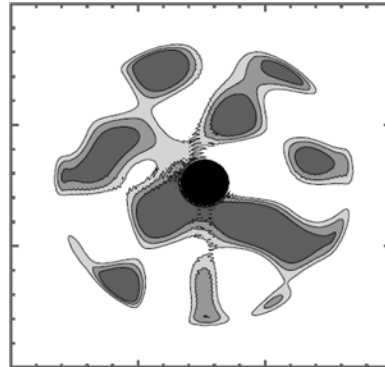
$t = 24$ hr

$t = 24$ hr

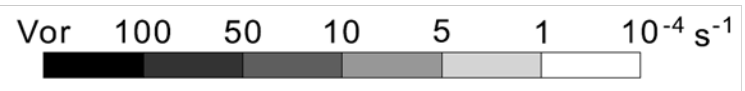
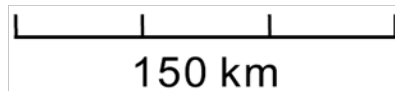
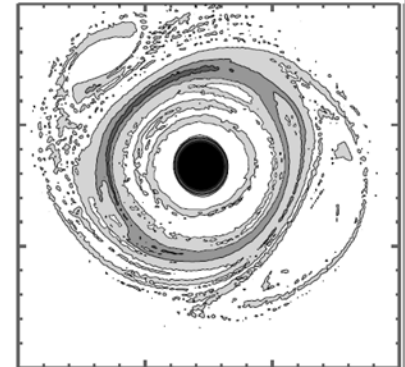
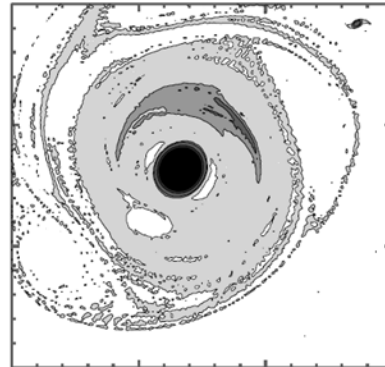
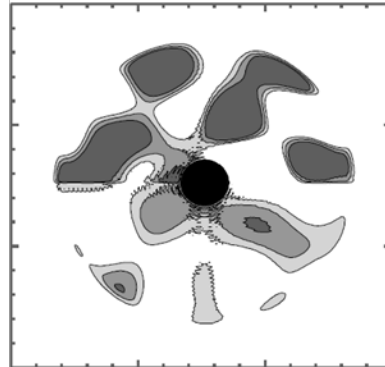
$\gamma = 6$

$\gamma = 10$

Isotropic
distribution



Anisotropic
distribution



Weiss(1981,1991), Rozoff et al. (2006)

$$\frac{D}{Dt}(\nabla \zeta) = -J(\nabla \psi, \zeta)$$

$$\rightarrow \nabla \zeta(t) \propto \exp(\lambda t) \quad \lambda = \pm \frac{1}{2} \sqrt{Q} = \pm \frac{1}{2} \sqrt{S_1^2 + S_2^2 - \zeta^2}$$

$$S_1 = \frac{\partial u}{\partial x} - \frac{\partial v}{\partial y} \quad (\text{stretch deformation})$$

$$S_2 = \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \quad (\text{shear deformation})$$

$Q > 0$ (strain dominates)

→ vorticity gradient will be stretched

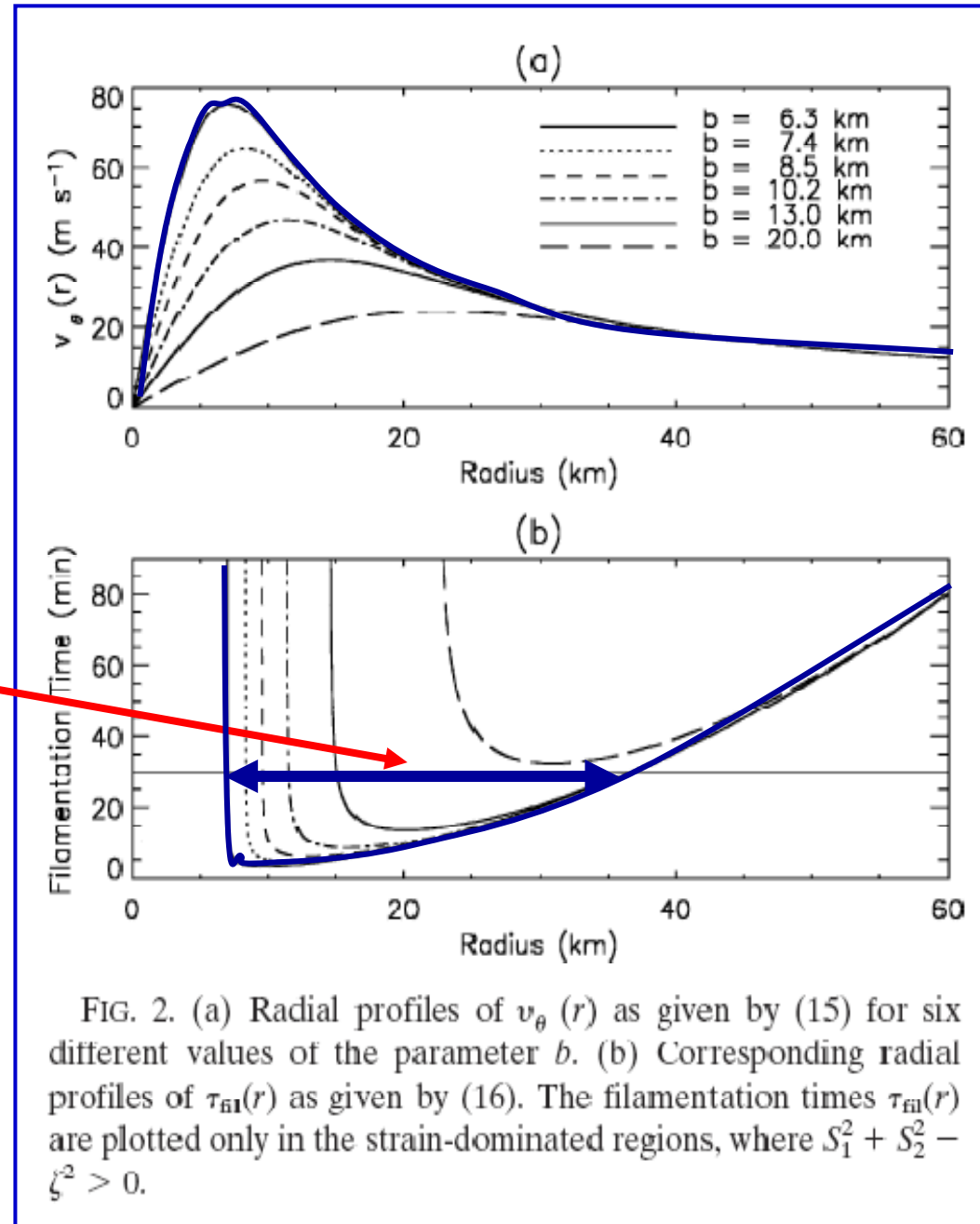
$Q < 0$ (vorticity dominates)

→ vortex is stable (survival of eyewall meso-vortices)

Rozoff et al. (2006)

The strong differential rotation outside the radius of maximum wind of the core vortex may also contribute to the formation and maintenance of the moat.

The Rapid Filamentation Zone: A zone with the filamentation time smaller than the 30 min convective turnover time.



Moat size (r_0) estimated from filamentation time

$$\tau_{fil} = 2 / \sqrt{\left(\frac{\partial v_\theta}{\partial r} - \frac{v_\theta}{r}\right)^2 - \left(\frac{\partial v_\theta}{\partial r} + \frac{v_\theta}{r}\right)^2}$$

$$\tau_{fil} = \frac{r}{v_\theta \sqrt{\alpha}} = \frac{2}{\zeta} \left(\frac{r}{r_m}\right)^{\alpha+1} \alpha^{-\frac{1}{2}} = \tau_0$$

$$\frac{r}{r_m} = \left(\frac{\zeta}{2} \tau_0 \alpha^{\frac{1}{2}}\right)^{\frac{1}{\alpha+1}}$$

$$r_0 = r - r_m = \left(\frac{\zeta}{2} \tau_0 \alpha^{\frac{1}{2}}\right)^{\frac{1}{\alpha+1}} r_m - r_m$$

v_m : best track data

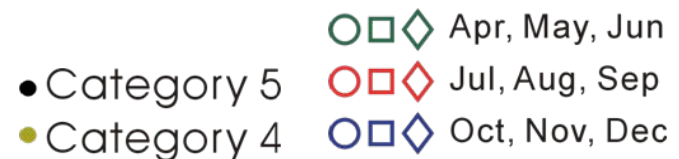
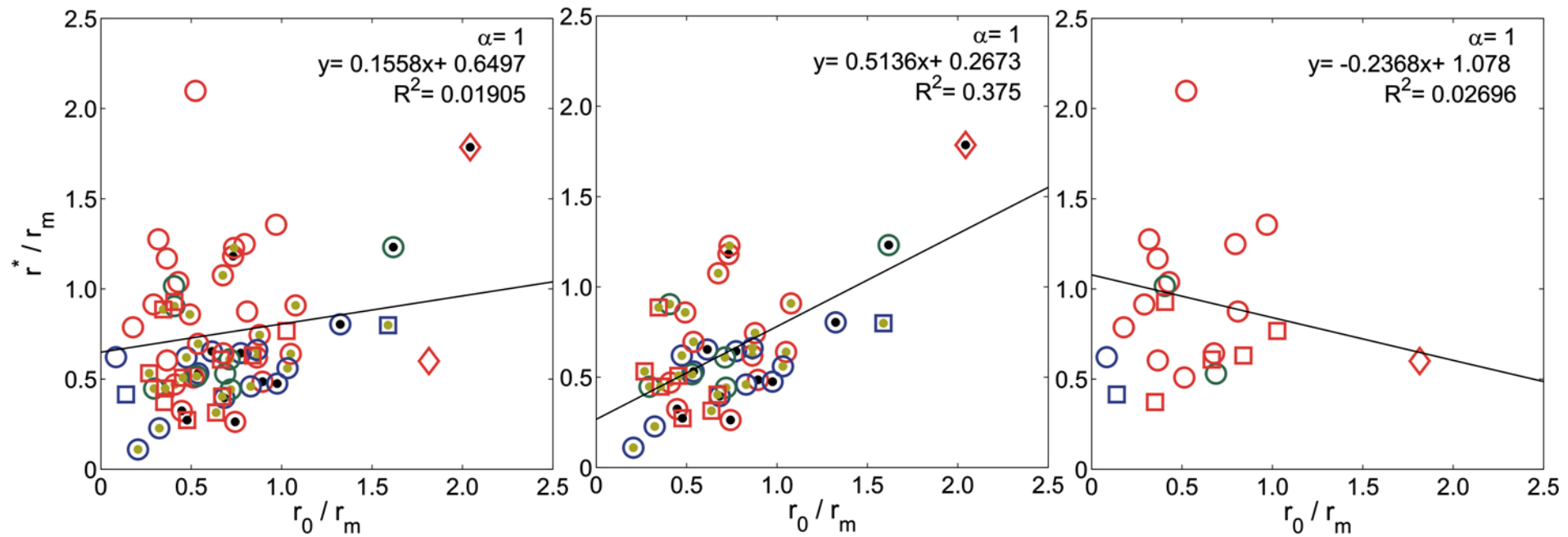
r_m : inner eyewall radius

$$v_\theta = v_m \left(\frac{r_m}{r}\right)^\alpha$$

$$\zeta = 2 \frac{v_m}{r_m}$$

$$\tau_0 = 30 \text{ min}$$

Nondimensional moat size v.s. nondimensional filamentation moat size



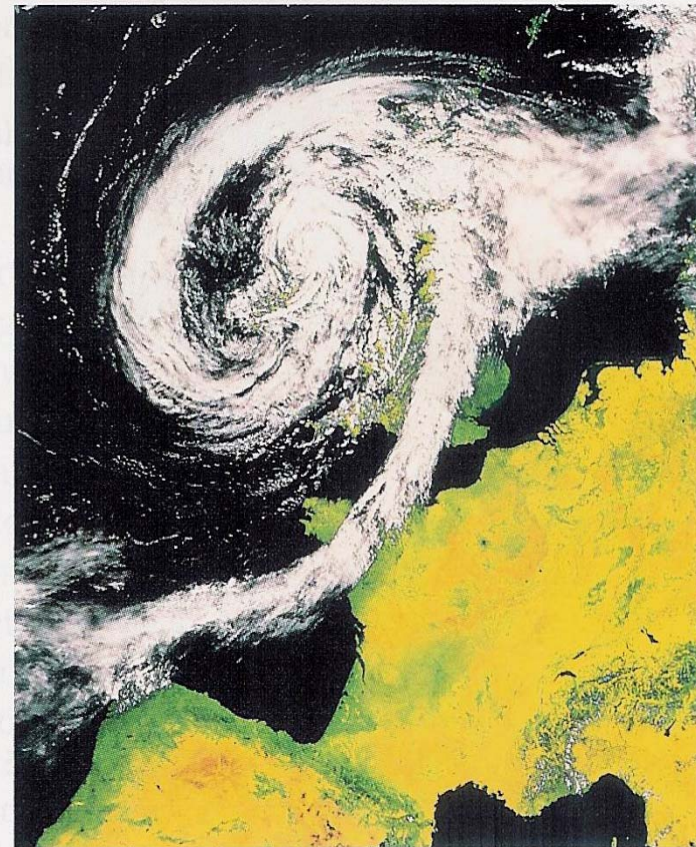
Summary

- As revealed by microwave imagery, tropical cyclones of sufficient strength (e.g. sustained wind speed ≥ 120 kts) often form concentric or double eyewalls. Area of asymmetric convection outside the core vortex that wraps around the inner eyewall to form the concentric eyewalls in about 12 hours.
- The outer eyewalls usually constricts the inflow of moisture, mass, and momentum into the inner eyewalls and the inner eyewall weakens and/or die.

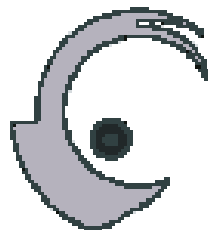
- A new dimension is added to the Dritschel-Waugh binary vortex interaction scheme that provides a proper concentric vorticity structure, the tripole vortex structure and the multiple eyewalls structure. Two important parameters are the vorticity strength ratio and the skirt parameters.
- The contraction of the secondary wind maximum and the formation of the moat are features of the vorticity dynamics. The moat formation by subsidence, rapid filamentation, and advective dynamics.

- Concentric vorticity structure formation requires a core vortex that is of 4 to 6 times stronger than the outer vortex. The outer vortex is larger than the core and with a separation distance from the core vortex that is within 6 times the core vortex radius. The vortex with (without) skirts allows the formation of outer band at large (small) radii.
- The pivotal role of the vorticity strength of the core vortex in maintaining itself, and in stretching, organizing and stabilizing the outer vorticity field, and the shielding effect of the moat to prevent further merger and enstrophy cascade processes in concentric eyewall dynamics.

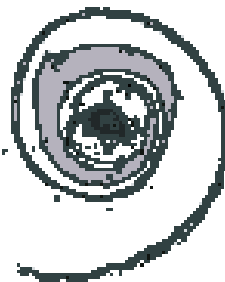
- Formation of concentric eyewalls is important for the TC intensity problem. A self-limiting process for TC intensity. A natural “STORMFURY”!!
- The organizational aspects of outer eyewall formation from an asymmetric vorticity field outside the cyclone core. The importance of core vortex structure, the moat, and the spatial characteristics of the vorticity field outside the core in the formation of double eyewalls.
- Need to understand the vorticity generating meso-scale processes in the TC environment.



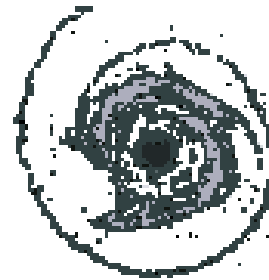
t=0 hr



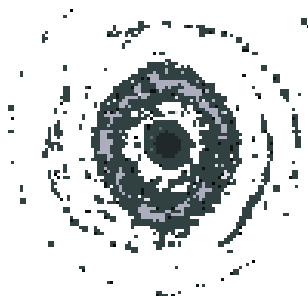
t=3 hr



t=8 hr



t=12 hr



$(\alpha = 1 \quad \gamma = 10 \quad r = 1/3)$

$r^* = 1$

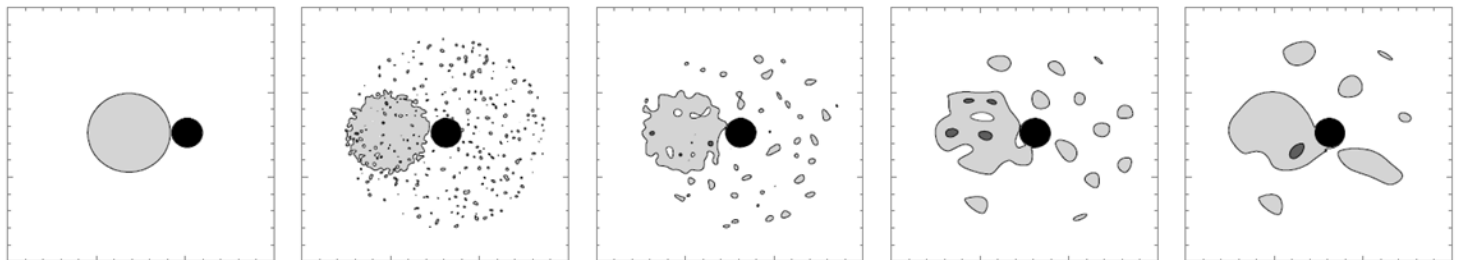
$r^* = 2$

$r^* = 3$

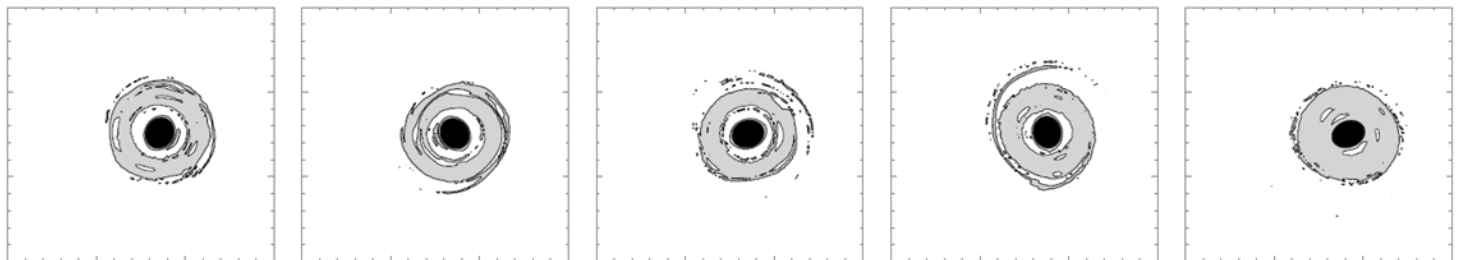
$r^* = 4$

$\Delta/R_1 = 0$

$t = 0 \text{ hr}$

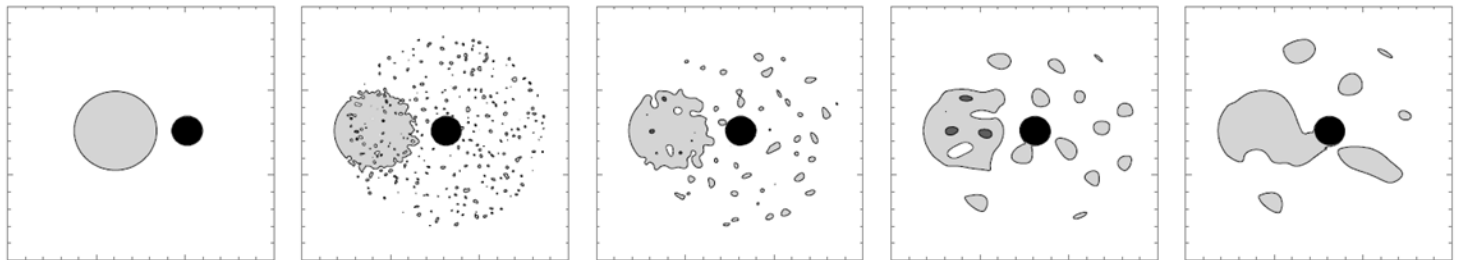


$t = 12 \text{ hr}$

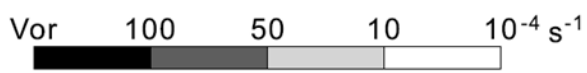
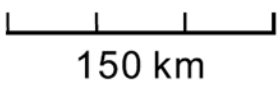
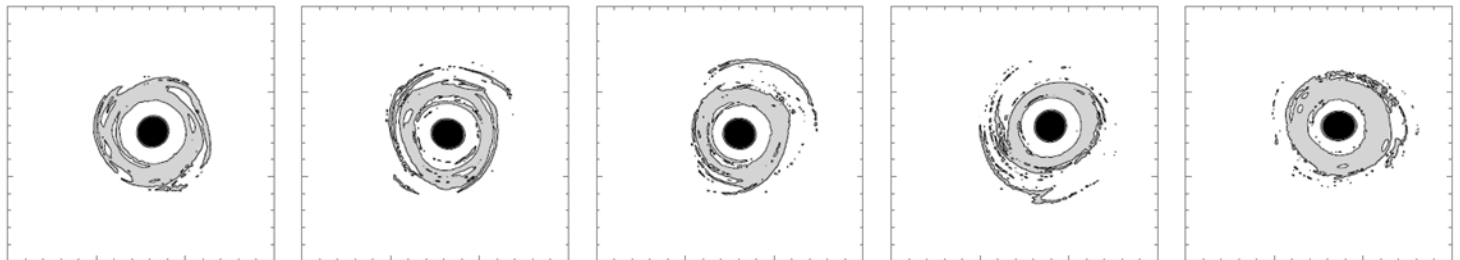


$\Delta/R_1 = 1$

$t = 0 \text{ hr}$



$t = 12 \text{ hr}$



$(\gamma = 8 \quad r = 1/4 \quad \Delta/R_1 = 5)$

$r^* = 1$

$r^* = 2$

$r^* = 3$

$r^* = 4$

$\alpha = 0.7$

$t = 0$ hr

$t = 12$ hr

$\alpha = 0.5$

$t = 0$ hr

$t = 12$ hr

